

# On Dualities in QFTs and String theory

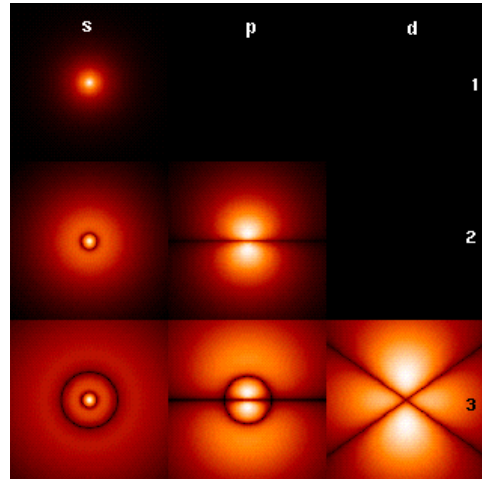
Hee-Cheol Kim

*POSTECH*

May 22, 2019

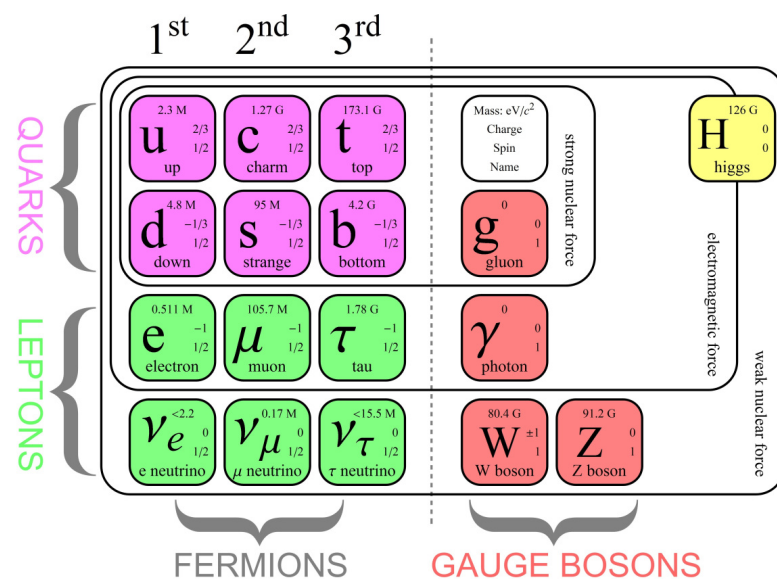
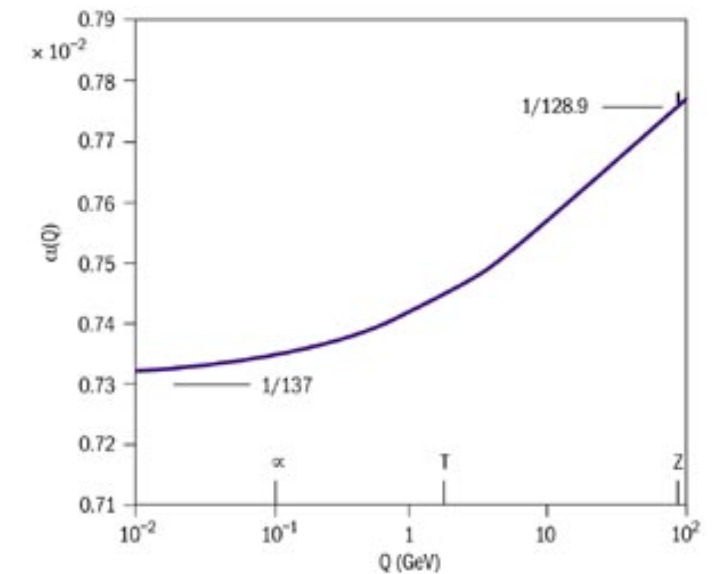
*Colloquium at SNU*

# Success of Quantum Frameworks



Electron density in Hydrogen Atom

## Renormalization of fine structure constant in QED

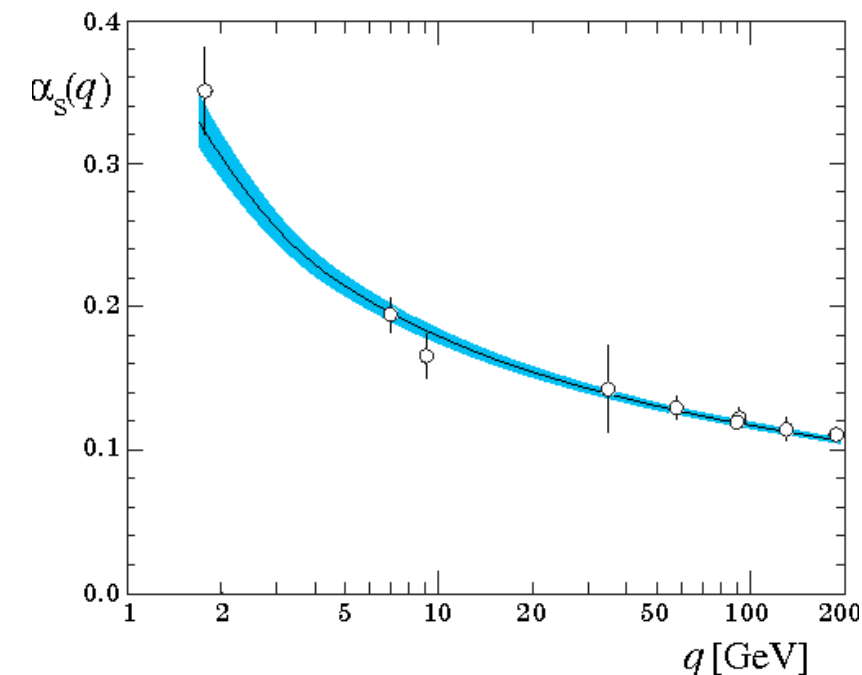


Standard Model

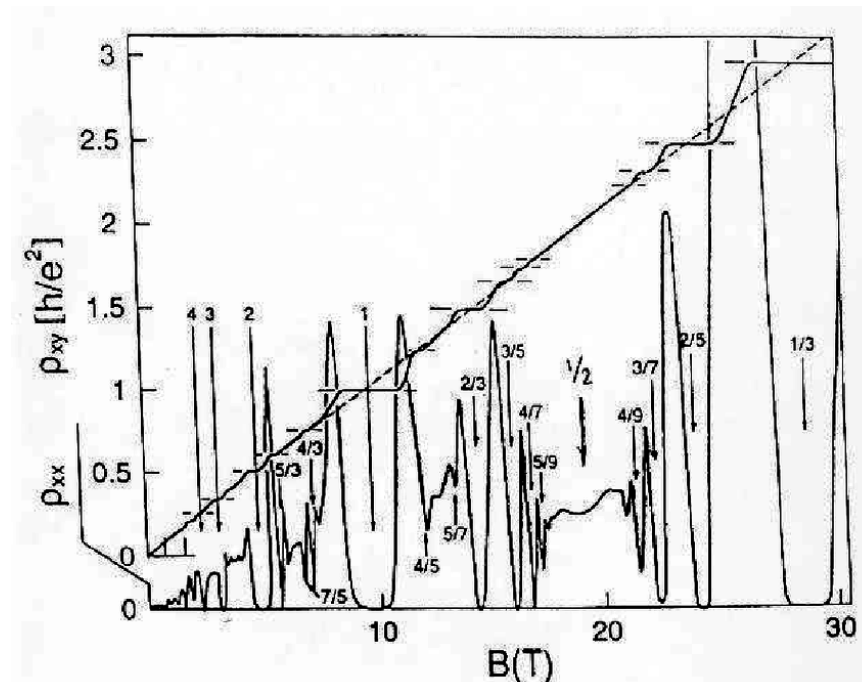
# Challenges in Modern Physics

However, “strong dynamics” in Quantum systems (QM or QFT) cannot in general be understood by using perturbative approaches.

- Quark confinement in QCD
- High  $T_c$  superconductor
- Fractional quantum Hall effect
- Black hole thermodynamics



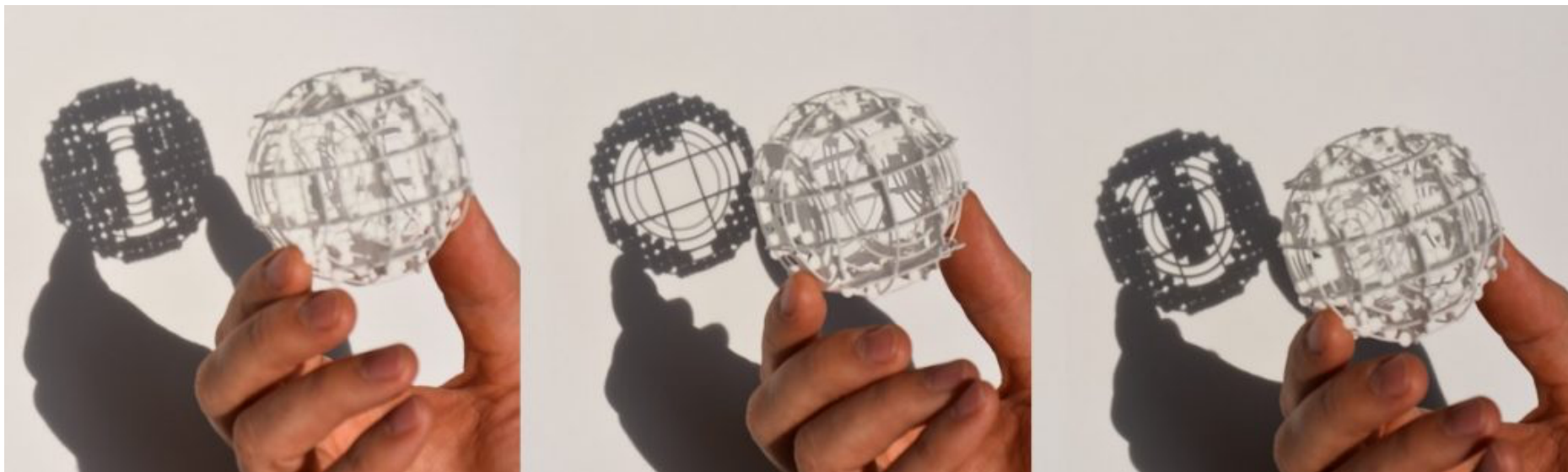
QCD running coupling constant



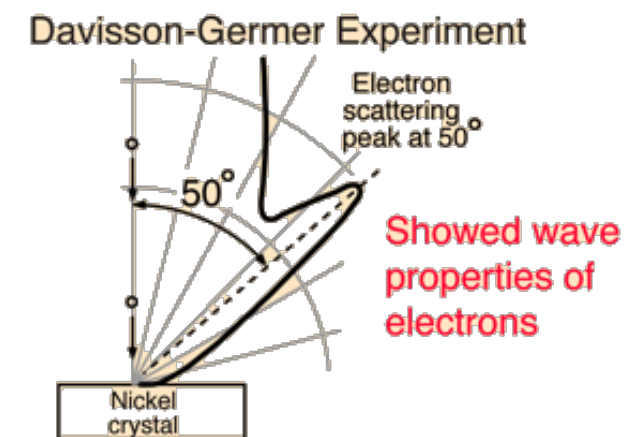
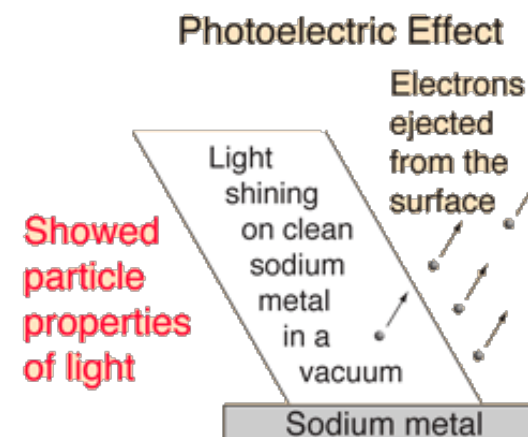
Fractional QHE

# Duality

- Strong dynamics in quantum systems can be studied by using “duality”!
- Two (or more) different descriptions of same physics  $\longrightarrow$  **Duality**



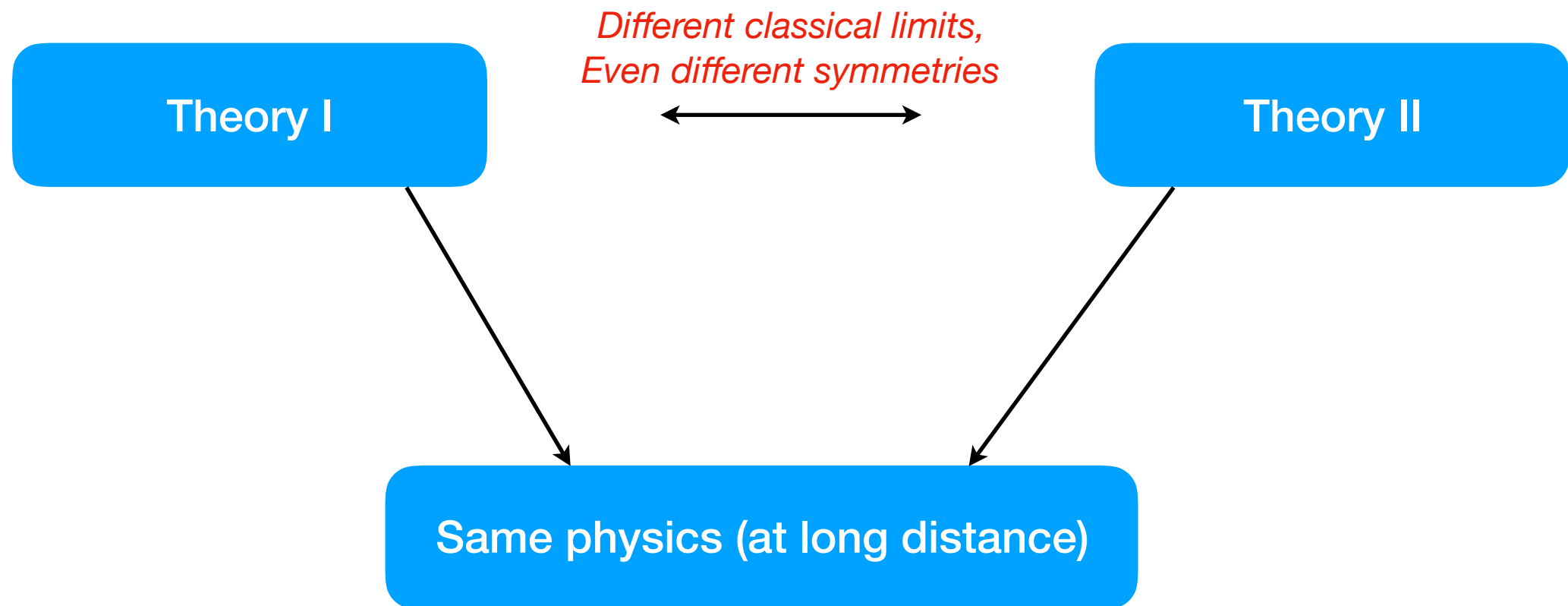
- Particle - Wave duality in QM





# Duality

- Two (or more) different descriptions of same physics



- Different descriptions provide different insights in physics.
- Strong-Weak duality : **String coupling**  $\longleftrightarrow$  **Weak coupling**

# Contents

- Introduction
- Dualities in  $d + 1$  ( $d \leq 2$ ) dimensions : Simple duality examples
- Dualities in  $d + 1$  ( $d \geq 3$ ) dimensions : More dualities
- Gauge/Gravity duality
- Conclusion

# ***Duality***

*that I found in “Sanya Nanshan Temple” in January*



*“Door of non-duality”*



*“Three sided statue of Guanyin Buddha” -> **Triality ?***

# Dualities in lower dimensions

# Duality of free Harmonic Oscillator

Free Harmonic Oscillator in quantum mechanics

$$\hat{H} = \frac{\hat{P}^2}{2m} + \frac{1}{2}m\omega^2 \hat{X}^2 \quad \text{with } \hat{P} = -\hbar \frac{\partial}{\partial X}$$

- Duality exchanges  $\hat{X} \leftrightarrow \hat{P}$  given by Fourier transformation

$$(\hat{X}, \hat{P}) \rightarrow \left( \frac{\hat{P}'}{m\omega}, -m\omega \hat{X}' \right)$$

- Duality maps  $\hat{H}$  to itself !

# Duality in 1+1d free bosons

Action for a compact free boson  $\phi(x) \sim \phi(x) + 2\pi$

$$S = \int d^2x \frac{\beta^2}{2} (\partial_\mu \phi)^2$$

Equation of motion :  $\partial^\mu \partial_\mu \phi(x) = 0$

Conserved currents ( $\partial_\mu j^\mu = 0$ ) :  $J^\mu = \beta^2 \partial^\mu \phi$  for  $\phi(x) \rightarrow \phi(x) + c$

$$J_W^\mu = \frac{1}{2\pi} \epsilon^{\mu\nu} \partial_\nu \phi \quad \text{from Bianchi id. } \epsilon^{\mu\nu} \partial_\mu \partial_\nu \phi = 0$$

- Duality by  $\partial^\mu \phi \leftrightarrow \epsilon^{\mu\nu} \partial_\nu \phi'$  leads to same action with  $\beta \rightarrow \frac{1}{\beta}$ .

$$\begin{array}{ccc} \beta \updownarrow \bigcirc & \xleftrightarrow{\text{T-duality}} & \bigcirc \updownarrow \beta' \sim \beta^{-1} \\ (J, J_W) & & (J'_W, J') \end{array}$$

- Equation of motion becomes trivial in dual theory of  $\phi'$ .



# Bosonization of 1+1d QFT

- Chiral fermions  $\psi_+, \psi_-$  can be written as functions of chiral bosons  $\phi_+, \phi_-$

$$\psi_+ \sim e^{-i\phi_+}, \quad \psi_- \sim e^{i\phi_-}$$

- Ex) Massless Thirring model is dual to a free boson [Thirring 1958], [Luscher, Peschel 1974], [Coleman 1975]

$$S = \int d^2x \, i\bar{\psi}\gamma^\mu \partial_\mu \psi - g(\bar{\psi}\gamma^\mu \psi)^2 \quad \longleftrightarrow \quad S = \int d^2x \, \frac{\beta^2}{2} (\partial_\mu \phi)^2$$

with  $\beta^2 = \frac{1}{4\pi} + \frac{g}{2\pi^2}$

- “Boson” and “Fermion” are not fundamental concepts in duality.

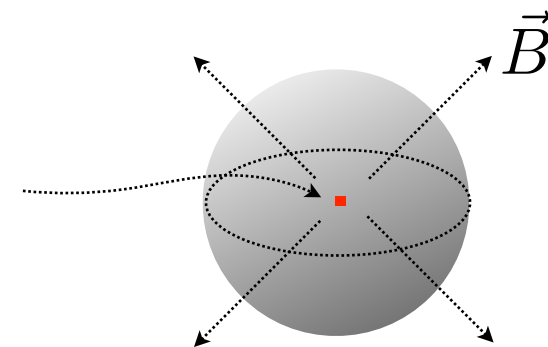
Free boson  $\longleftrightarrow$  Interacting fermion

# Simplest duality in 2+1d Maxwell theory

$$S = -\frac{1}{4e^2} \int d^3x F_{\mu\nu} F^{\mu\nu} \quad \text{with} \quad \begin{aligned} F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu \\ A_\mu &: \text{U(1) gauge field} \\ A_\mu(x) &\rightarrow A_\mu(x) + \partial_\mu \Lambda(x) \end{aligned}$$

Topological  $U(1)_{\text{top}}$  current  $J_{\text{top}}^\mu = \frac{1}{4\pi} \epsilon^{\mu\nu\rho} F_{\nu\rho}$  from Bianchi id.  $\epsilon^{\mu\nu\rho} \partial_\mu F_{\nu\rho} = 0$

Charged object is magnetic monopole OP  $\mathcal{M}(x)$



- Dual to a free theory of a compact scalar field  $\sigma$  by  $F^{\mu\nu} \sim \epsilon^{\mu\nu\rho} \partial_\rho \sigma$ .
- Monopole is implemented explicitly in dual theory as  $\mathcal{M}(x) \sim e^{i\sigma(x)}$ .
- Gauge symmetry can disappear or emerge, so it's not fundamental.

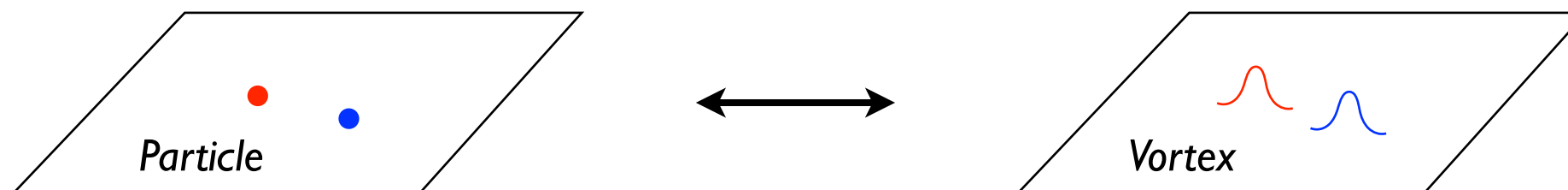
# Particle-Vortex duality in 2+1d

	Particle theory	Vortex theory
Lagrangian	$ D_A \phi ^2 -  \phi ^4$	$ D_a \Phi ^2 -  \Phi ^4 + \frac{1}{2\pi} a dA$
$U(1)$ current	$j_\phi$	$\frac{1}{2\pi} \epsilon_{\mu\nu\lambda} \partial^\nu a^\lambda$
Operator	$\phi$	$\mathcal{M}_a$
Mass deformation	$-m \phi ^2$	$m \Phi ^2$

$A$  : background gauge field

$a$  : dynamical gauge field

[Peskin 1978], [Dasgupta, Halperin 1981]



- Non-perturbative “soliton” is mapped to elementary particle.
- Two distinguished phases  $\begin{cases} m > 0 & : \text{Mass gap} \\ m < 0 & : \text{Superfluid} \end{cases}$
- Emergent gauge symmetry of “ $a$ ” in “vortex” theory.

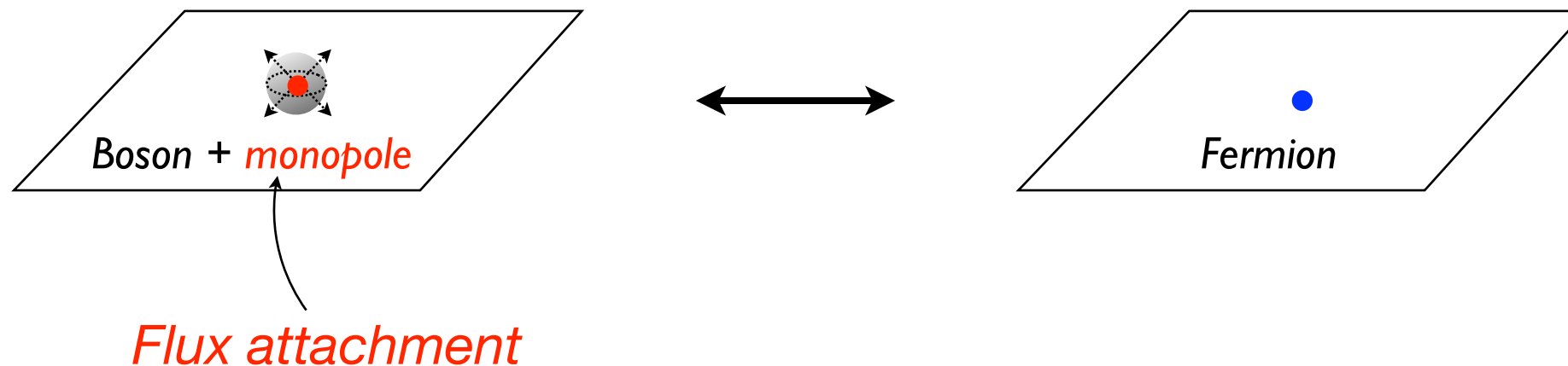
# Boson-Fermion duality in 2+1d

Scalar field  $\phi$  coupled to U(1) gauge field with Chern-Simons coupling is dual to a free Fermion  $\psi$ !

<i>Lagrangian</i>	$ D_a \phi ^2 -  \phi ^4 + \frac{1}{4\pi} a da$	$i\bar{\psi}\gamma^\mu \partial_\mu \psi$
U(1) <i>current</i>	$\frac{1}{2\pi} \epsilon_{\mu\nu\rho} \partial^\mu a^\nu$	$j_\psi$
<i>Operator</i>	$\mathcal{M}_a \phi$	$\psi$

**“Bosonization”**

[Wilczek, Zee 1983], [Polyakov 1988]



- “Spin” is not fundamental concept under duality!

# Dualities in higher dimensions

# Duality in 3+1d Maxwell theory

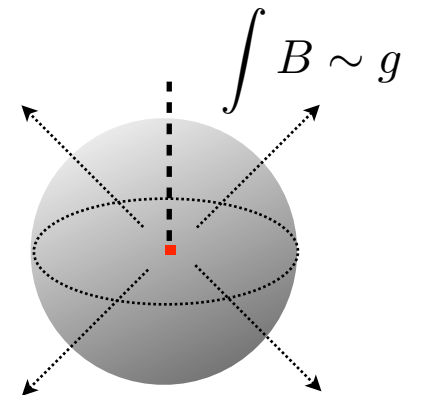
Maxwell equation without charge source

$$\begin{aligned}\nabla \cdot E &= 0 & \nabla \cdot B &= 0 \\ \nabla \times E &= -\frac{\partial B}{\partial t} & \nabla \times B &= \frac{\partial E}{\partial t}\end{aligned}$$

- Electric-magnetic duality :  $(E, B) \rightarrow (B', -E')$  or  $F^{\mu\nu} \rightarrow \epsilon^{\mu\nu\lambda\rho} F'_{\lambda\rho}$

- Charge sources under duality

Electron with charge  $e$   $\longleftrightarrow$  Monopole with charge  $g$



Dirac monopole

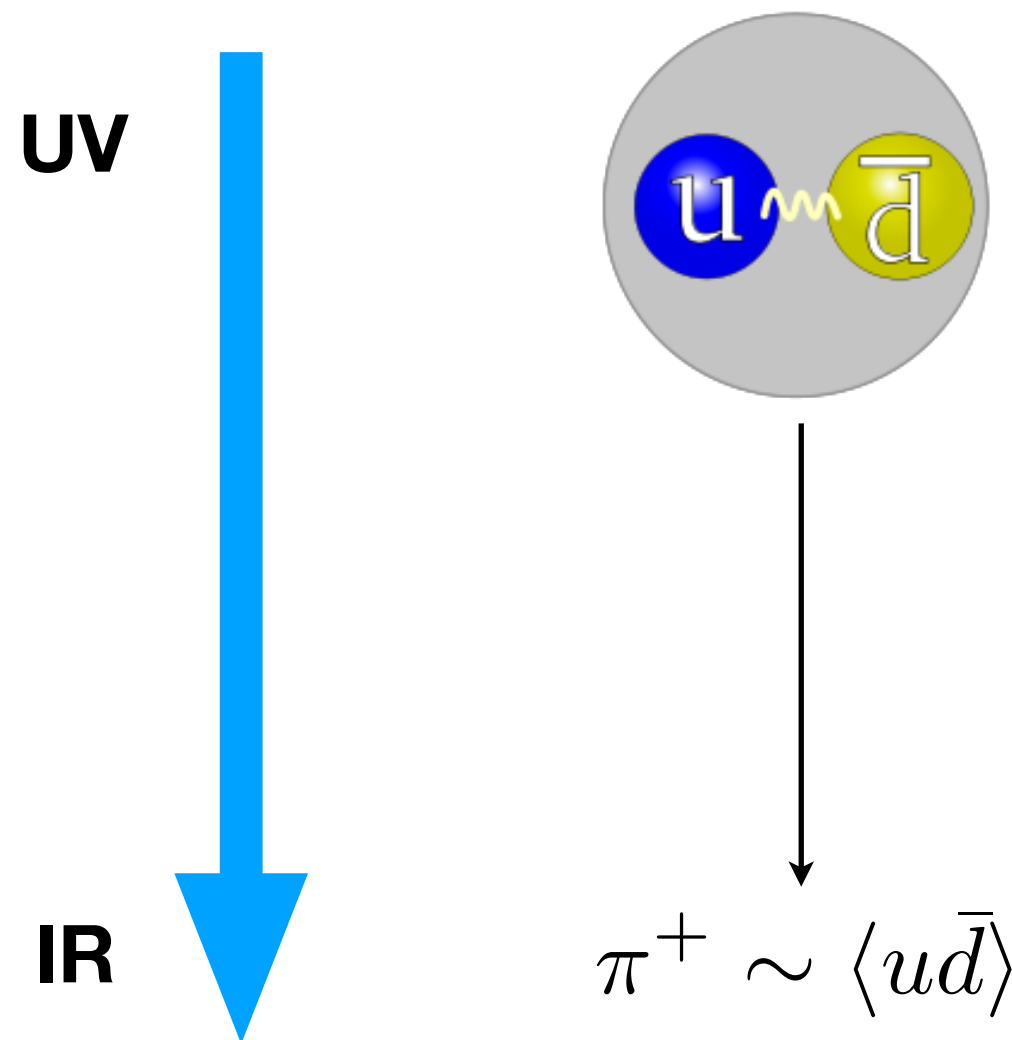
- Strong/Weak duality

$A_\mu$  with  $e$   $\longleftrightarrow$   $A'_\mu$  with  $g \sim e^{-1}$  due to charge quantization  $e \cdot g = 2\pi\hbar$



# Quarks vs Mesons

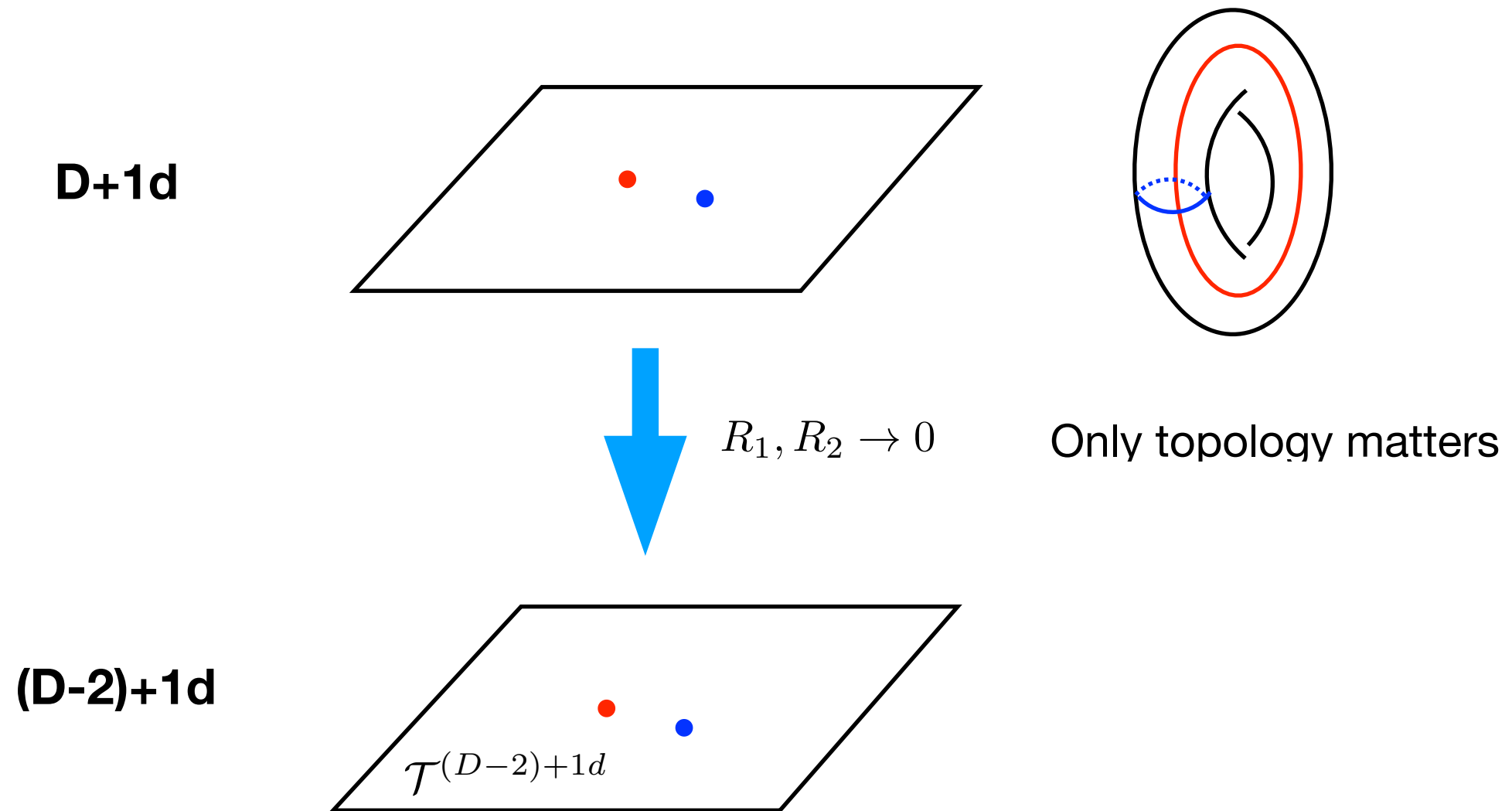
- **QCD** describes strong forces between quarks and gluons **at short distance**.
- On the other hand, **physics at long distance** is described by **Chiral perturbation theory** of mesons such as pions  $\pi^\pm, \pi^0$ .



- Degrees of freedom at short distance  $\neq$  those at long distance.

# Idea of Compactification

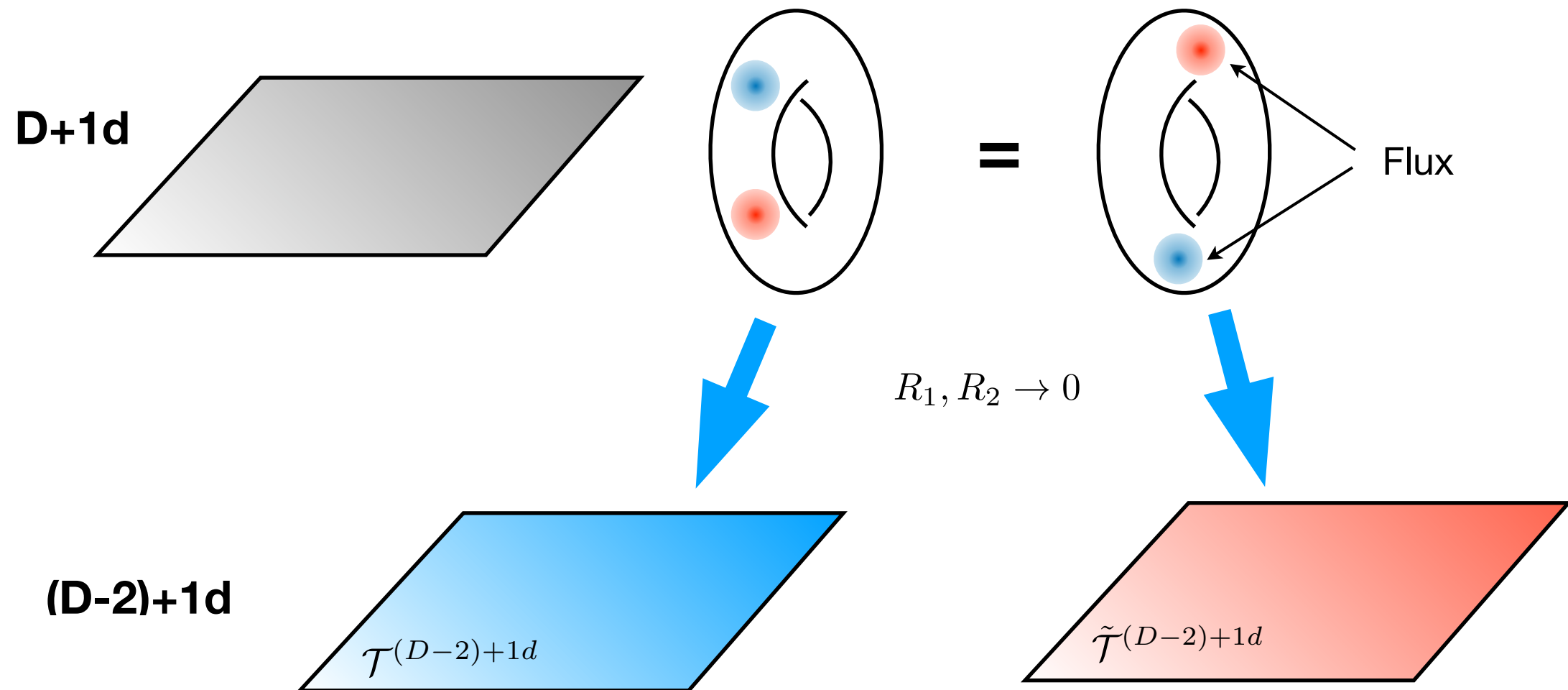
- $D+1d$  QFT compactified on a torus leads to  $(D-2)+1d$  QFT.



- Duality exchanging  $\bullet \longleftrightarrow \bullet$  comes from  $R_1 \longleftrightarrow R_2$  on torus.

# Idea of Compactification

- $D+1d$  QFT on a torus with fluxes leads to a family of  $(D-2)+1d$  QFTs.

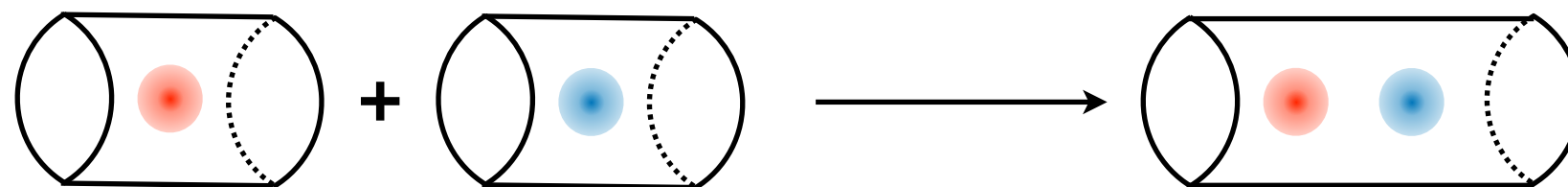


- Duality exchanging  $\mathcal{T} \longleftrightarrow \tilde{\mathcal{T}}$  comes from  $\text{blue circle} \longleftrightarrow \text{red circle}$  on torus.

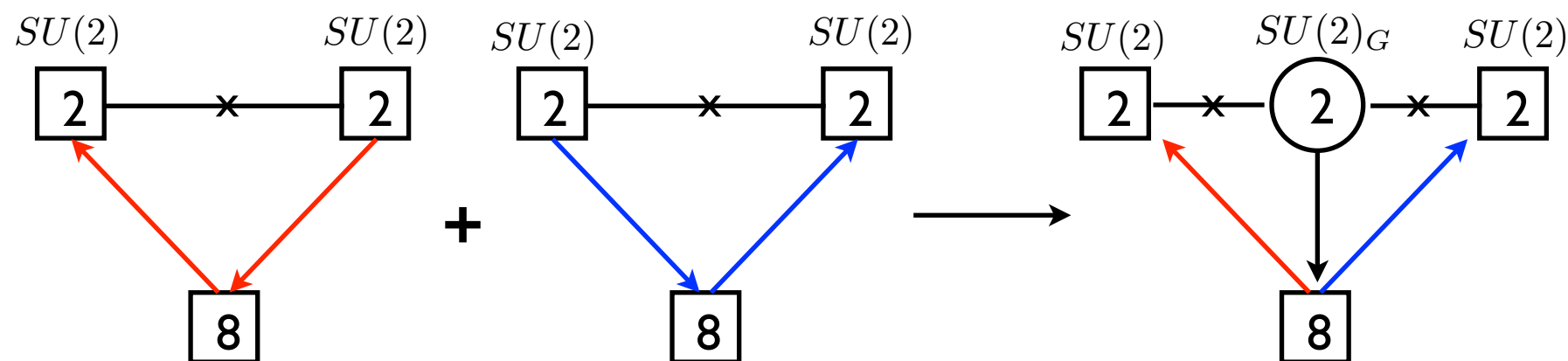
# 5+1d E-string theory and compactification

- 5+1d E-string theory [Witten 1995], [Ganor, Hanany 1996],...
  - Minimal supersymmetric conformal theory with  $E_8$  global symmetry
- LEGO-like construction** of Lagrangians for such 3+1d QFT from 5+1d E-string theory compactified on a torus with  $E_8$  fluxes. [H-C Kim, Razamat, Vafa, Zafrir 2017]

LEGO blocks for fluxes



3+1d QFT operators



- Generic flux on torus  gives rise to infinite number of 3+1d QFTs

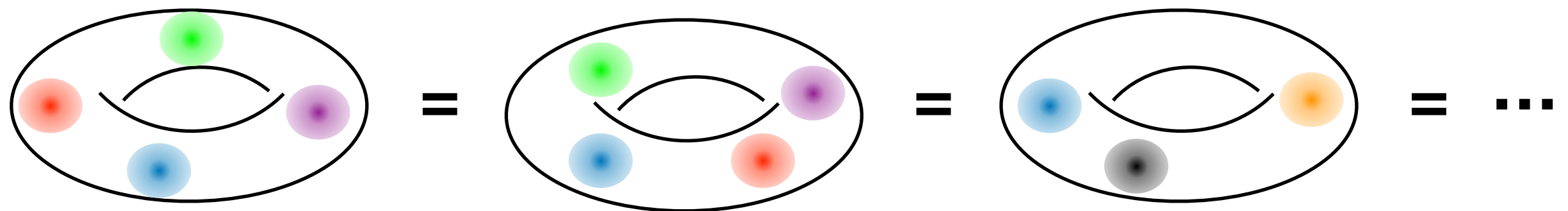
# 3+1d Dualities from 5+1d E-string theory

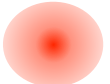
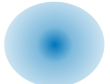
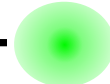
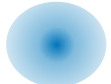
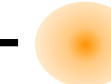
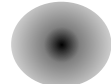
Generic flux on torus  gives rise to infinite number of 3+1d QFTs

3+1d QFTs only depend on topology and fluxes.

[[H-C. Kim, Razamat, Vafa, Zafrir 2017](#)]

→ Infinitely many dualities !

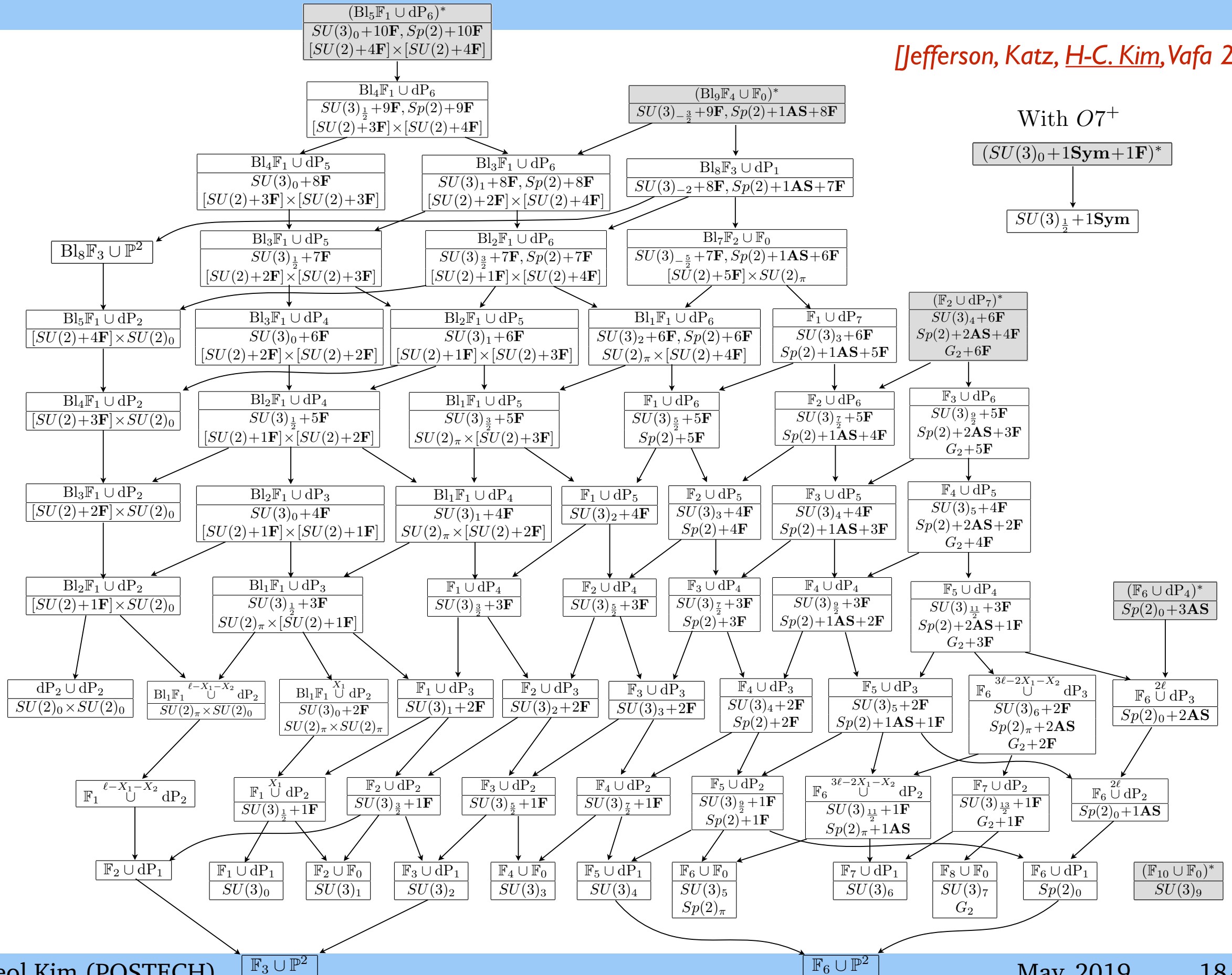


when  +  +  +  =  +  +  = ...

- Higher dimensional QFTs can provide easy interpretations and also new predictions of lower dimensional physics such as dualities, enhanced symmetries.

# Classification of 4+1d rank 2 QFTs and Dualities

[Jefferson, Katz, H-C. Kim, Vafa 2018]



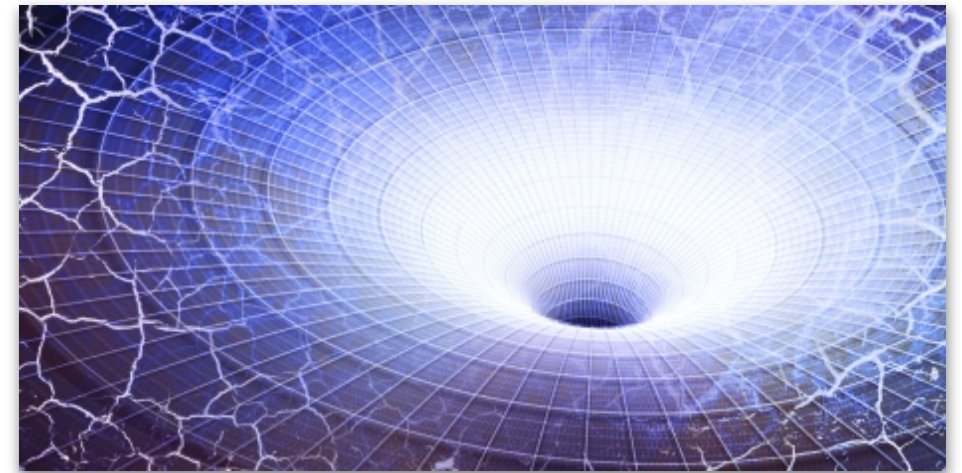


# Gauge/Gravity Duality

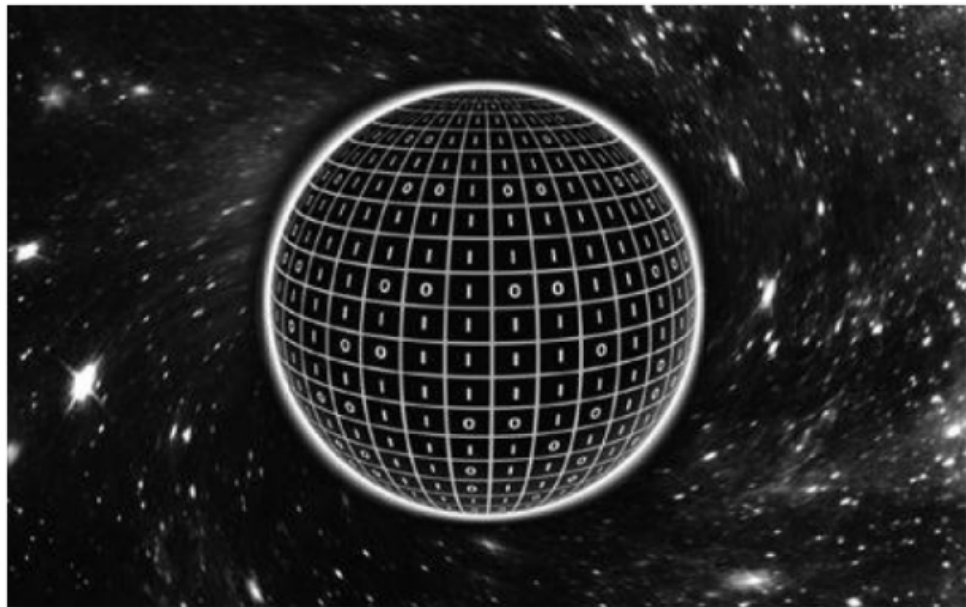
## Spacetime is dynamical [Einstein 1915]

- In quantum world, spacetime geometry arises from **quantization** of graviton field (or **metric field**  $g_{\mu\nu}$ ).

*However, Gravity is non-renormalizable.*



<https://www.nottingham.ac.uk/mathematics/research/mathematical-physics>



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- Thermodynamical entropy of black holes (BH)

$$S_{\text{BH}} = \frac{k A_{\text{BH}}}{4\hbar G_N}$$

$A_{\text{BH}}$  : BH surface area

$G_N$  : Newton constant

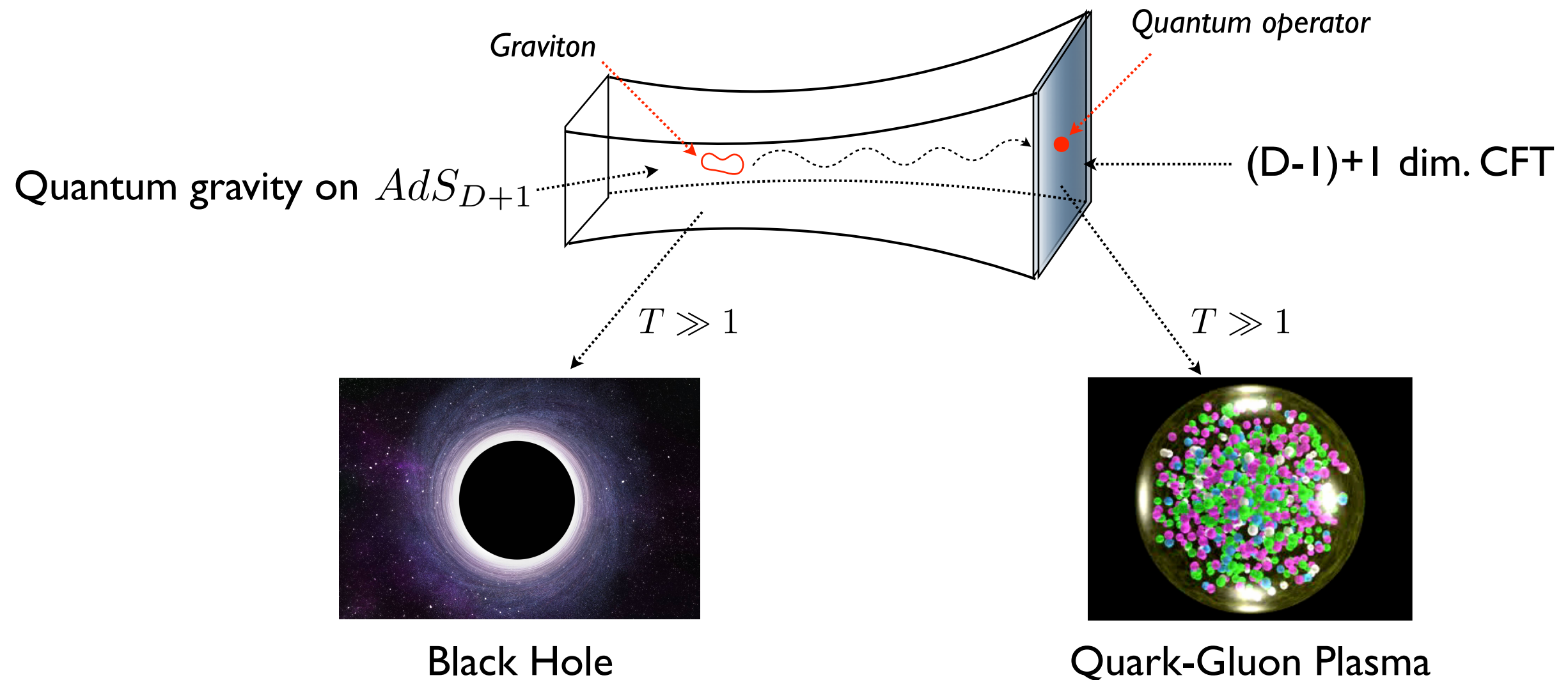
[Bekenstein 1973], [Hawking 1974]

Quantum gravity can be formulated in terms of degrees of freedom living on the boundary of the spacetime.  $\longrightarrow$  **Holography**

[’t Hooft, Susskind]

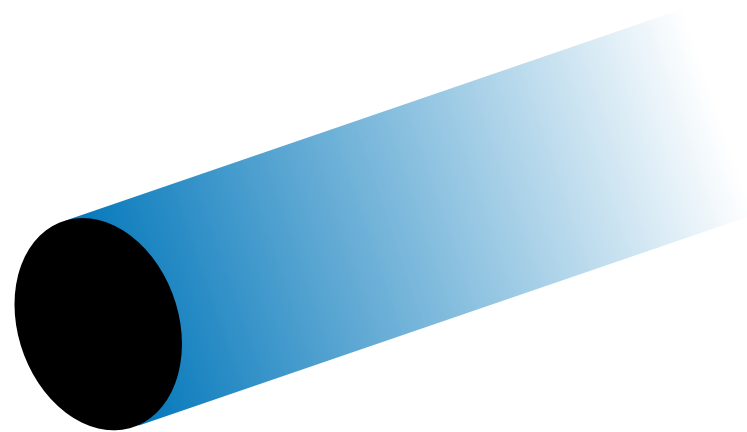
# Gauge/Gravity Duality (Holography)

- Quantum gravity in a  $D+1$  dimensional curved “anti-de Sitter” spacetime (AdS space) is holographically described by  $(D-1)+1$  dimensional conformal field theory (CFT) living on the boundary of AdS space. [Maldacena 1997]

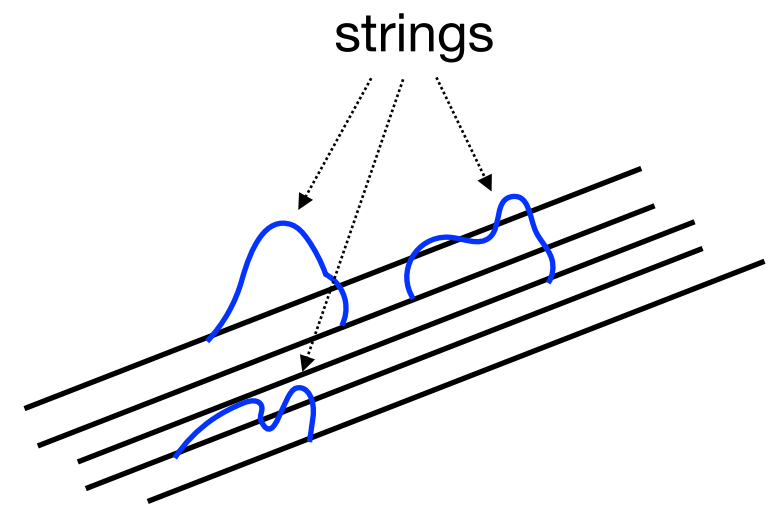


# Black Hole Entropy Counting

- Strominger and Vafa black hole counting (1996) :  
We can count **black hole microstate** and find black hole entropy  $S_{BH}$  by using dual theory living on D-branes.



**Black brane**



**D-branes**

# Summary

- Duality is a fact that two or more seemingly different quantum systems describe the same physics.
- Strong-Weak duality can be used to understand strong dynamics in quantum systems.
- Many examples : Bosonization, Particle/Vortex duality, QCD vs Mesons, Gauge/Gravity duality, ....