

CURRENT STATUS AND FUTURE PROSPECTS OF BRIGHT X-RAY SOURCES



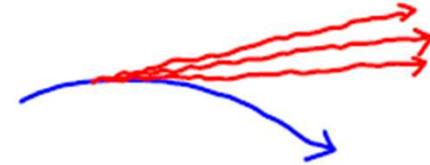
KWANG-JE KIM

University of Chicago
Argonne National Lab

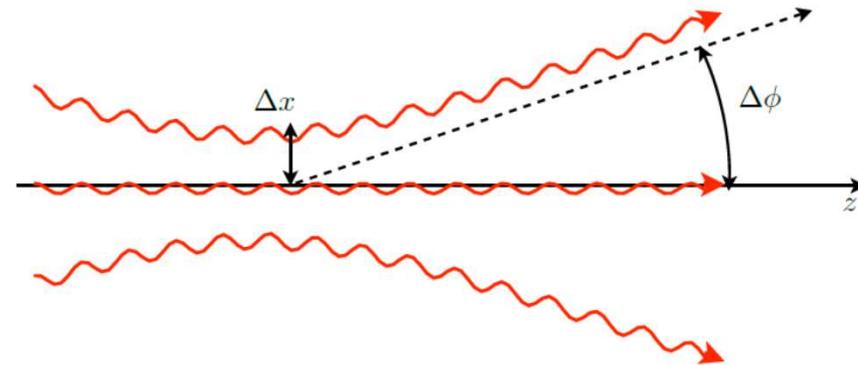
November 21, 2018
Physics Department, Seoul National University

RELATIVISTIC ELECTRONS AND SYNCHROTRON RADIATION

- If an electron with velocity $v \rightarrow c$, $\gamma = \frac{1}{\sqrt{1-(v/c)^2}} \rightarrow \infty$ makes a curved motion under magnetic field, short wavelength radiation is radiated: $\lambda \sim \rho/\gamma^3$
- # of photons in $\frac{1}{\gamma}$ angle $\sim \alpha \simeq 1/137$



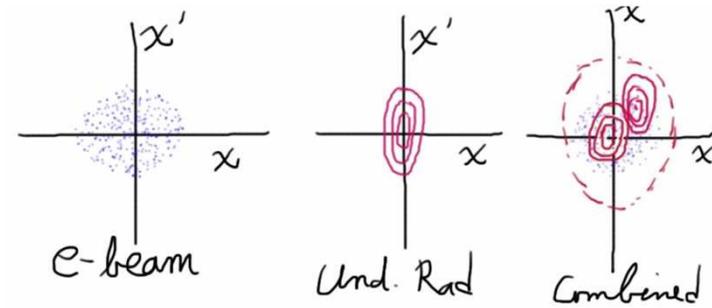
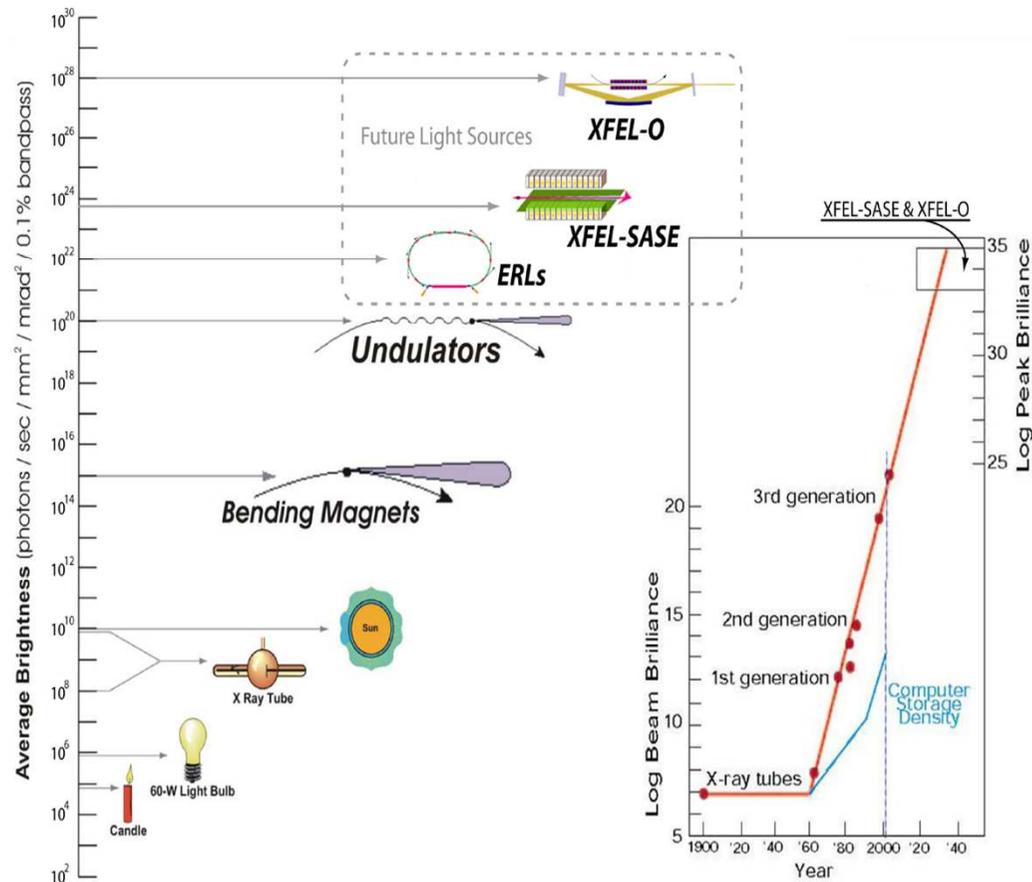
RADIATION (PHOTON) IS USEFUL (BRIGHT) IF THE TRANSVERSE EXTENT IS **SMALL & DIRECTIONAL**



- $\Delta x \cdot \Delta \phi =$ phase space area $= \Delta x \cdot \Delta p_x / m\gamma$: Phase space area
- Liouville's theorem in Hamiltonian mechanics: Phase space area is conserved
- Beam physics terminology: $\Delta x \cdot \Delta \theta =$ "emittance" $= \varepsilon_x$
- Brightness = # of photons / $\varepsilon_x \varepsilon_y$. *Real figure of merit since it is conserved*

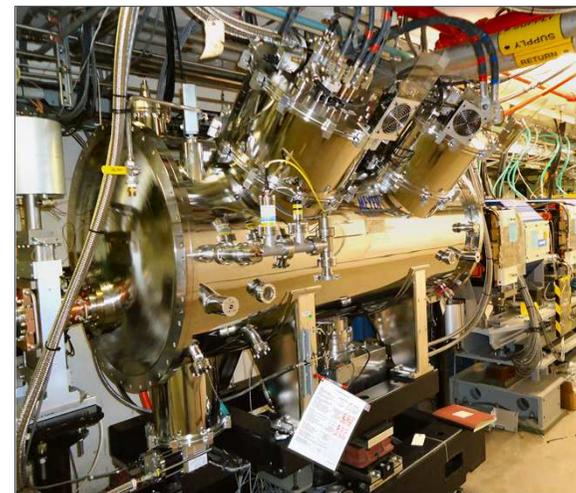
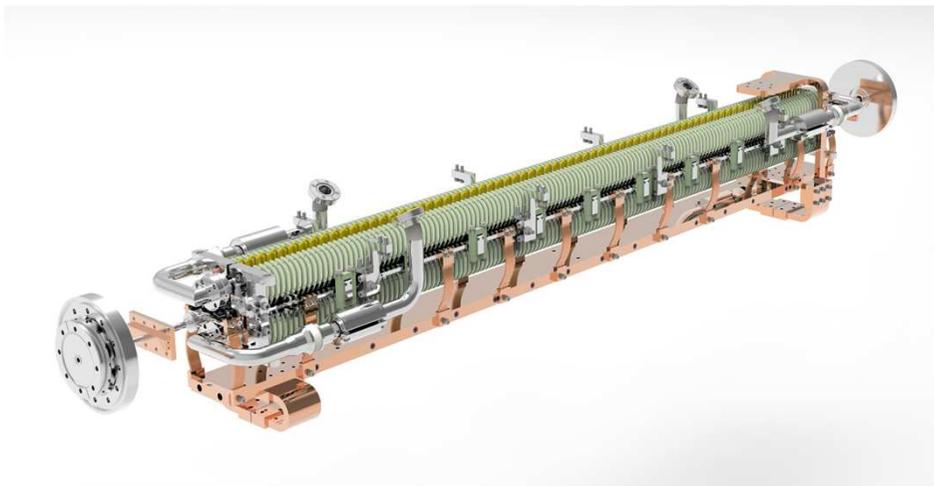
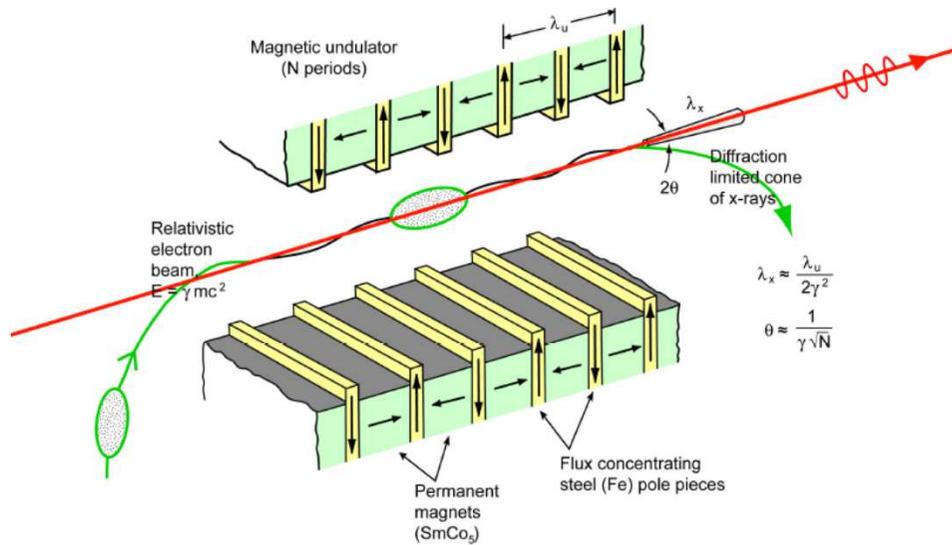
X-RAY BEAM PHASE SPACE IS A CONVOLUTION OF THE RADIATION AND E-BEAM PARTS

Advance of X-ray beam brightness follows the advance of the electron storage ring performance



PRODUCING DIRECTED RADIATION (I)

UNDULATORS: EM → PM → SC



SCU18-1 in Sector 1 of the APS ring.

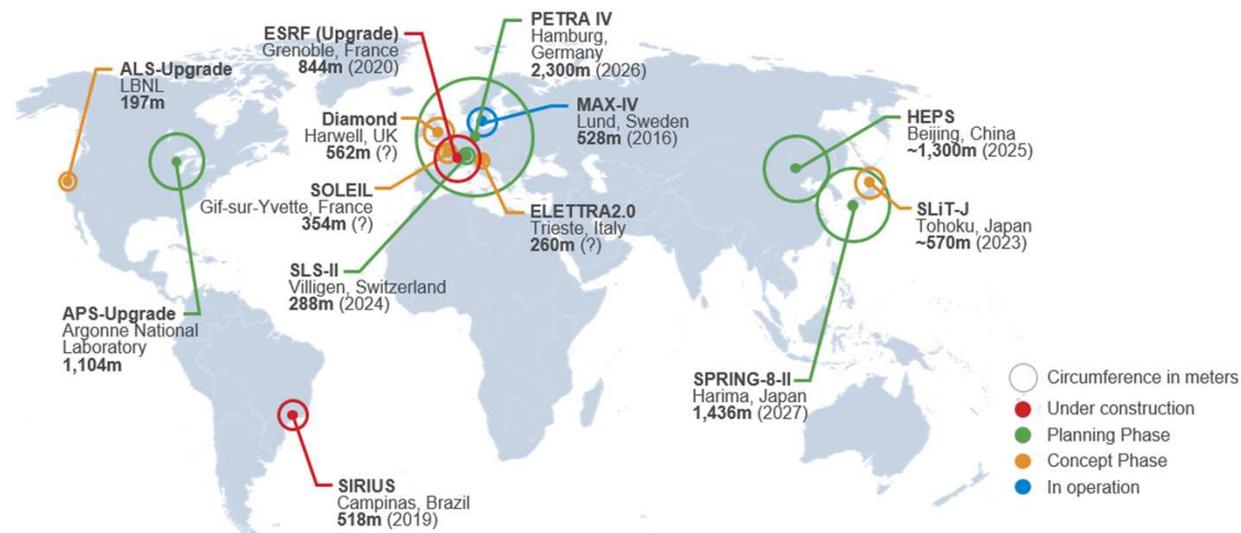
PRODUCING DIRECTED RADIATION (II): REDUCING THE ELECTRON BEAM EMITTANCE

- E-beam emittance is determined from the balance of two effects: damping from the classical nature and diffusion due to the quantum nature (discrete photon effect) of the synchrotron radiation in bending magnets

→ Use a short achromatic bending units

– $(N_D = \# \text{ of dipoles/sector})$

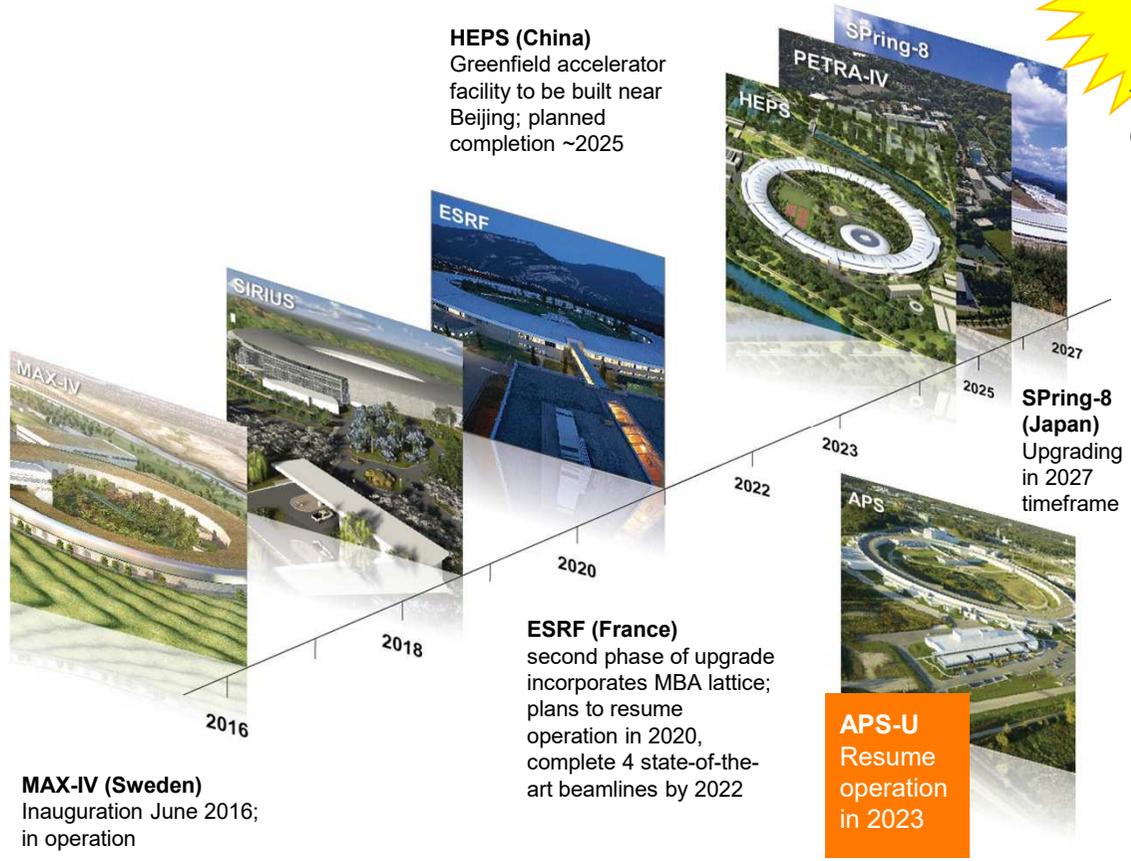
- 3rd generation source: $N_D = 2,3$
- “Diffraction limited storage ring (DLSR)” sources for X-rays: $N_D = 7$
- Many DLSR projects, new or retrofitting 3rd gen facilities around the world



INTERNATIONAL HIGH ENERGY “DIFFRACTION LIMITED (DL) SR DEVELOPMENT

$$\frac{\lambda}{4\pi}$$

Diffraction limited transverse emittance

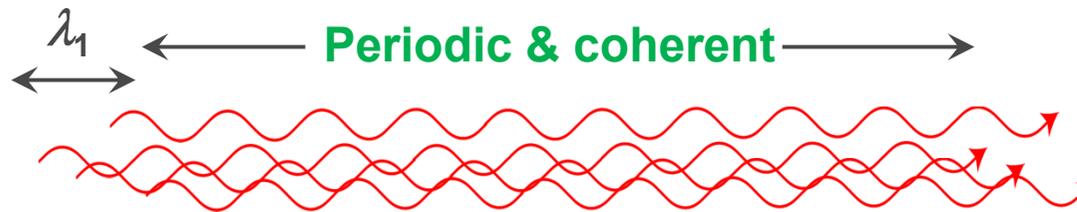


TEMPORAL COHERENCE

- Monochromator improves temporal coherence $:\Delta\omega \rightarrow \Delta\omega_M \ll \Delta\omega$,



- If e-beam bunch length \gg wavelength \rightarrow the intensity of wave-trains adds incoherently $\sim N_e$

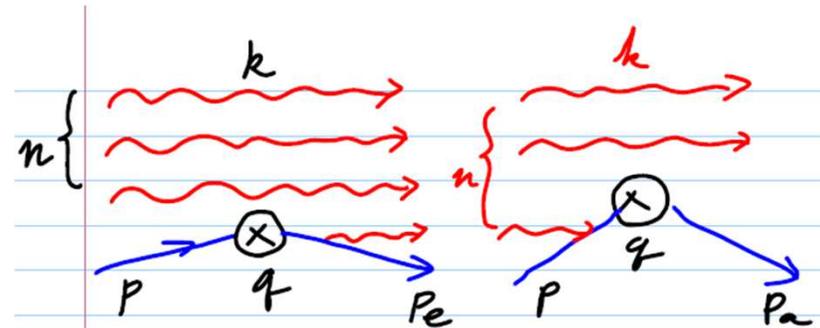


- The amplitudes add in phase if $\Delta z_{el} \ll \lambda_1$ or if electrons are concentrated at positions $z = n\lambda_1, n=1,2,..$ **Intensity $\sim N_e^2$**



- This happens in free-electron laser

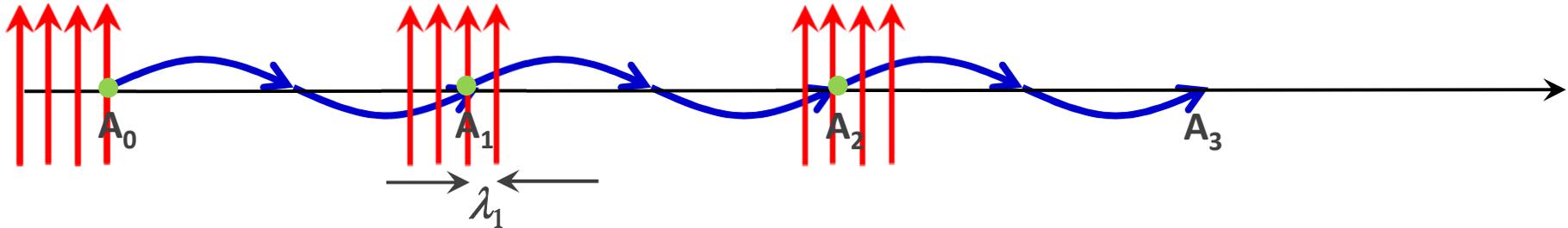
FREE ELECTRON LASER IS A LASER



- For every spontaneous process, there exists stimulated emission
- In the presence of n coherent modes, the emission probability $\propto n + 1$ while absorption probability $\propto n \rightarrow$ light amplification
 - J. Weber (1951), Maiman, Townes, Shawlow, Basov, Porkhorov,..
- Free electron laser—the emission process involves “free electrons” in an undulator magnetic field rather than bound electrons in atoms or molecules
 - John Madey (1971)
- In the limit of low photon energy, drops out \rightarrow classical interpretation

Classical View of a Free Electron laser

- When the EM wavelength satisfies the undulator condition, an electron sees the same EM field in the successive period \rightarrow sustained energy exchange

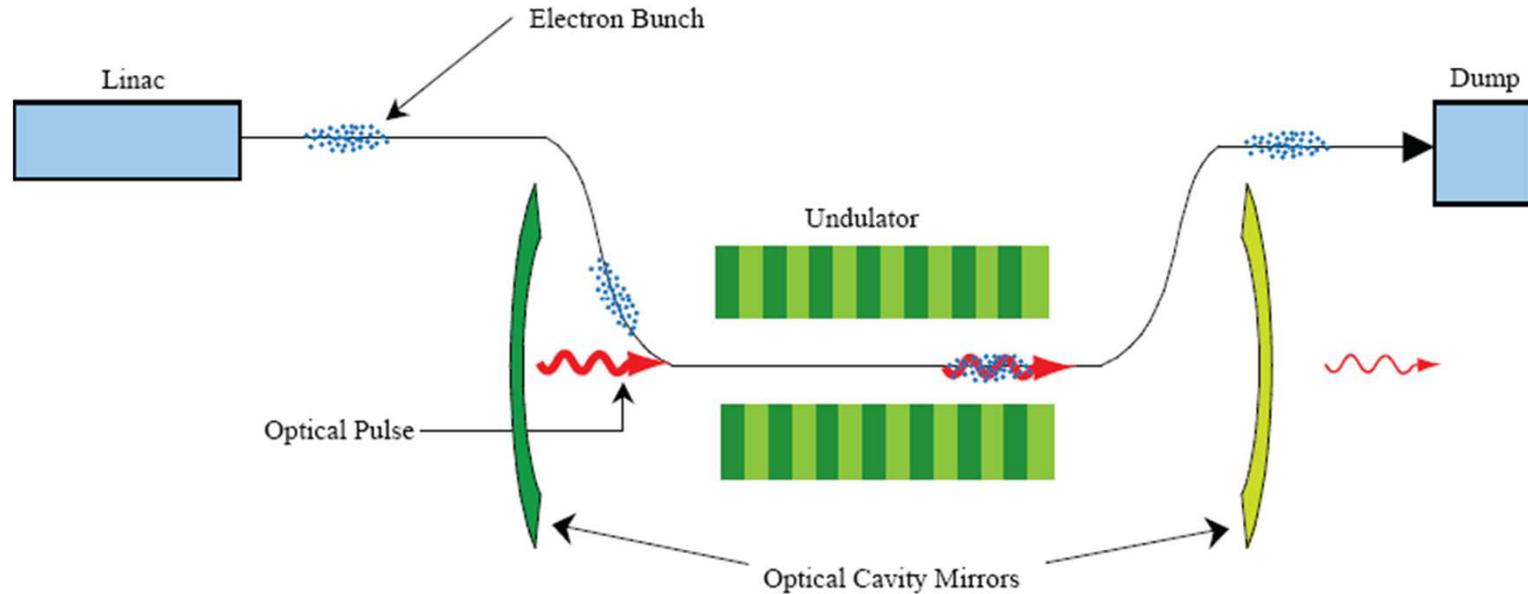


- An e^- arriving at A_0 loses energy to the field ($e\mathbf{v} \cdot \mathbf{E} < 0$). Similarly the e^- at distance $n\lambda_1$, $n=1,2,\dots$ also loses energy. However, those at $\lambda_1(1/2 + n)$ away gain energy.
- The electron beam develops energy modulation (period length λ_1).



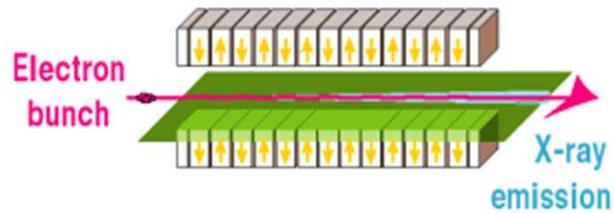
- Higher energy electrons are faster \rightarrow density modulation develops
- 
- A second green wavy line representing density modulation, showing a sinusoidal wave with a period corresponding to λ_1 .
- Coherent EM of wavelength λ_1 is generated \rightarrow “Free electron laser”

CHALLENGES FOR X-RAY FELS



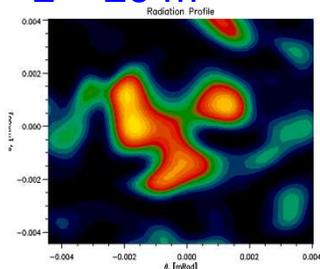
- **X-rays are difficult to reflect** → **Low-gain oscillators are difficult**
 - Make a single pass, high-gain ($\sim 10^6$) device for **self-amplified spontaneous emission (SASE)** for intense, **quasi-coherent** X-rays
 - The required e-beam can be produced from an RF photo-cathode gun, compression, and acceleration
- **Use Bragg reflector (diamond)** → **XFELO** producing **fully coherent** X-rays

X-ray FEL solution I :Initial noise is amplified to intense, quasi-coherent radiation (SASE)

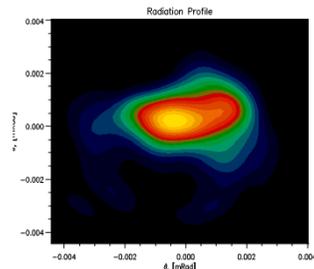


Transverse mode

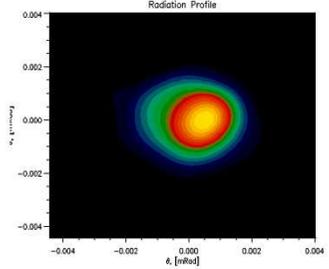
$z = 25 \text{ m}$



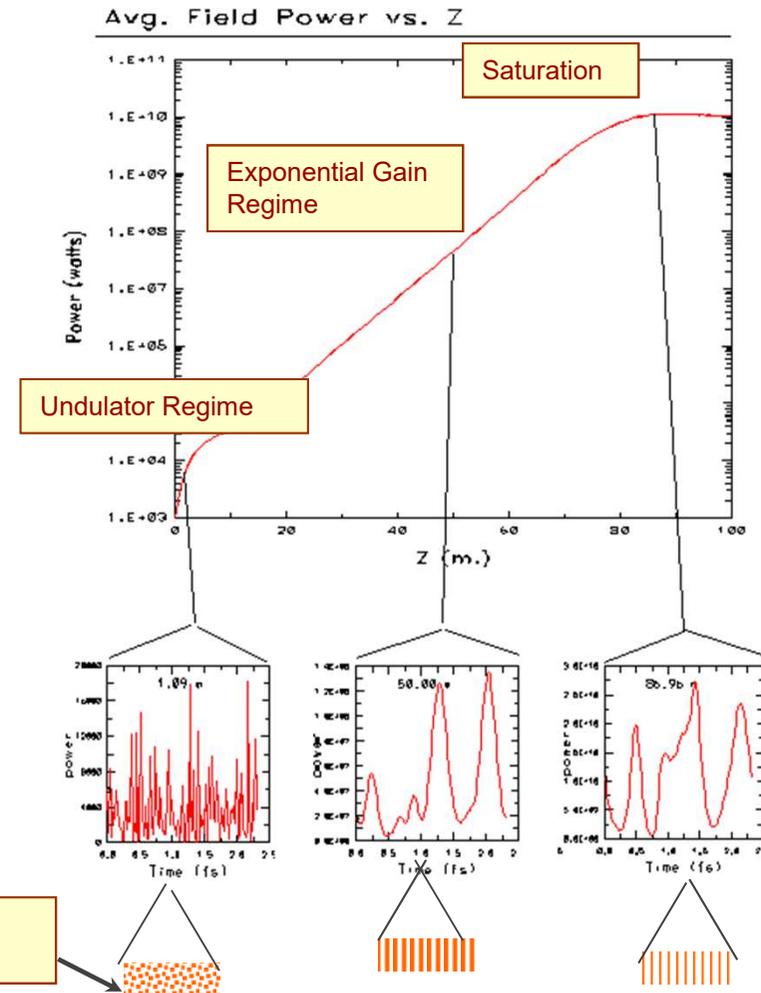
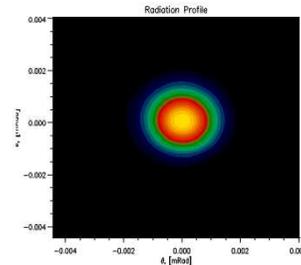
$z = 37.5 \text{ m}$



$z = 50 \text{ m}$

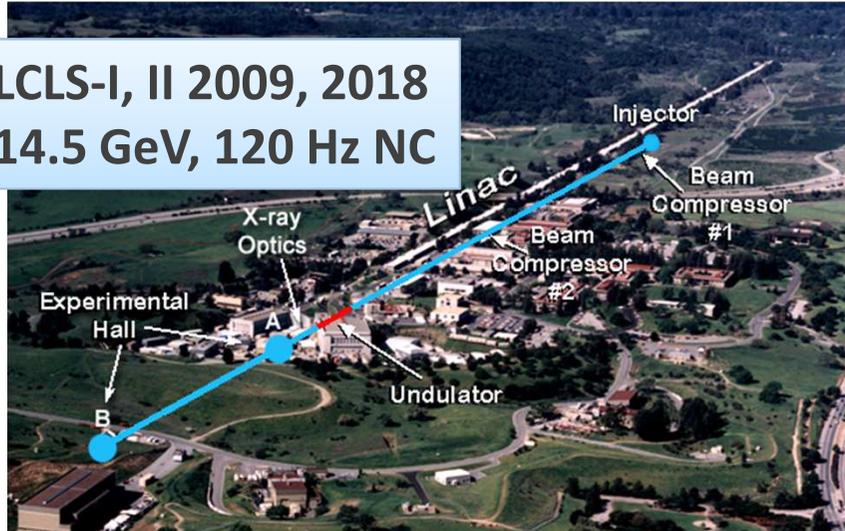


$z = 90 \text{ m}$



Hard X-Ray FELs Operating in 2017

LCLS-I, II 2009, 2018
14.5 GeV, 120 Hz NC



SACLA 2011
8.5 GeV, 60 Hz NC



XFEL 2017
17.5 GeV, 3000 x 10 Hz SC



PAL XFEL 2017
10 GeV, 60 Hz NC

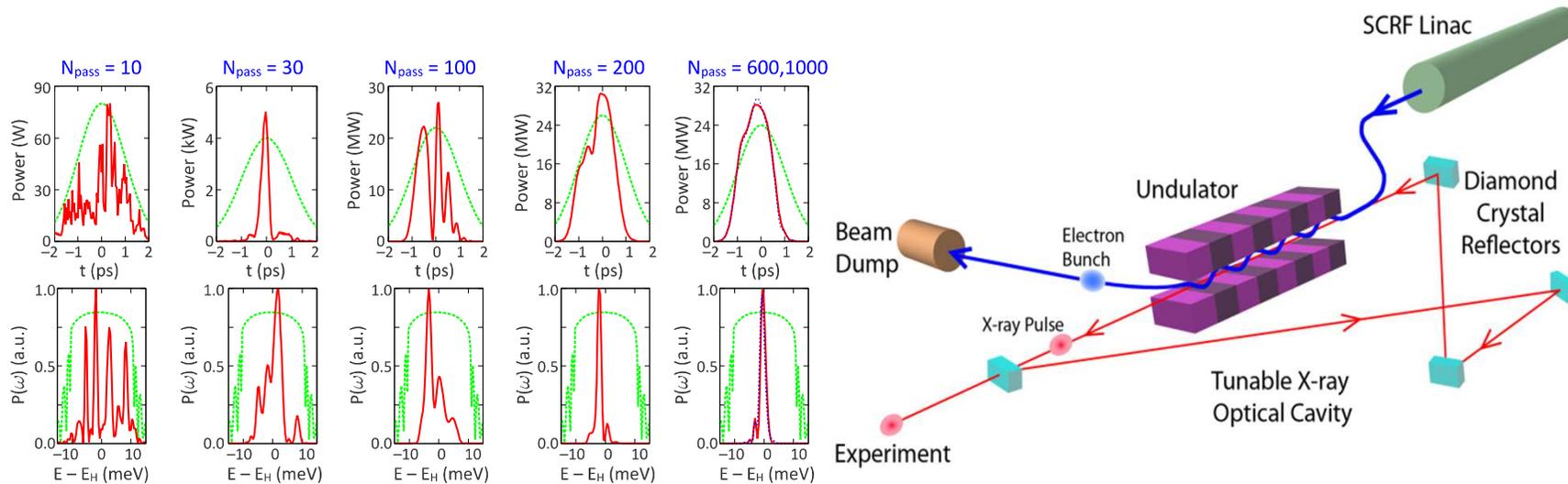


$\lambda_U=15\text{mm}$
 $K=1.2$
 $G=4.5\text{ mm}$

SWISS FEL 2017
5.8 GeV, 100 Hz NC



X-RAY FEL SOLUTION II: FEL OSCILLATOR BY USING BRAGG REFLECTORS



- First proposed by R. Collela and A. Luccio (1983)
- Revived by KJK, Y. Shvyd'ko, S. Reiche (2008) spurred by the ERL development
- Will provide ultra-narrow BW (a few meV) and **full coherence**
- *Real X-ray laser!!*



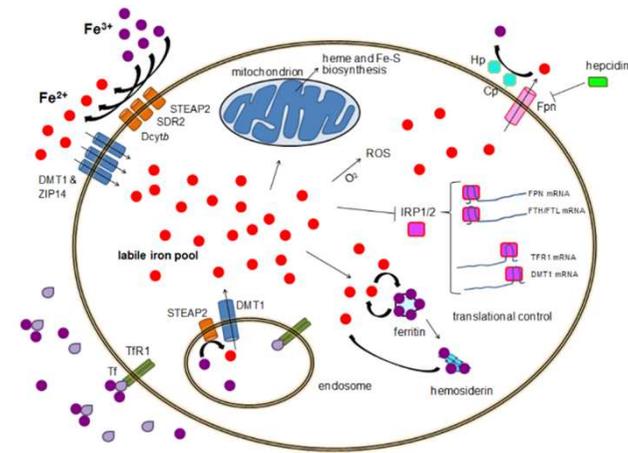
A SHORT LESSON IN KOREAN AND CHINESE

- Je (제, 齊)– ordered, coherent
- Kwang (광, 光)--light
- Kwang-Je → *Coherent Radiation*



According to Yuan T. Lee

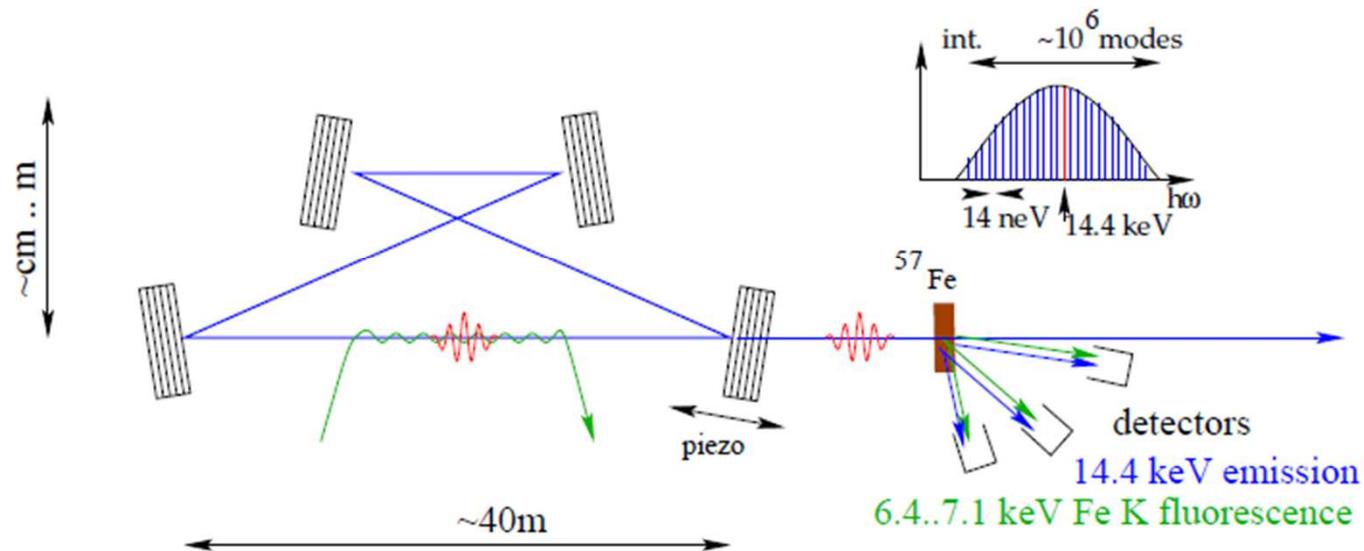
SCIENCES WITH XFELO



- **Enhanced application of techniques developed at 3GSR sources**
 - IXS, XPCS, NRS → higher resolution, smaller volume,..
 - Deep-earth core material, strongly correlated system,..
- **Techniques in infancy at current sources**
 - Coherent control with X-rays, X-ray NLO,..
 - Practical applications of NRS, such as study of red cells without enriching the excited states of Fe
- **Emergence of new areas**
 - Nuclear quantum optics, entangled state
 - X-ray spectral comb → 10 orders of magnitude improvement in precision atomic spectroscopy (QED test), quantum optics of nuclear states

Nuclear-resonance-stabilized XFEL-O (B.W. Adams and K.-J. Kim, 2015, PRSTAB)

- The XFEL-O output pulses are copies of the same circulating intra-cavity pulse → **By stabilizing cavity length to a fraction of λ , the spectrum of XFEL-O output becomes a comb**
- The extreme-stabilized XFEL-O will
 - establish an x-ray-based length standard
 - Revolutionize the quantum optics of nuclear states
 - have applications in fundamental physics such as x-ray Ramsey inter

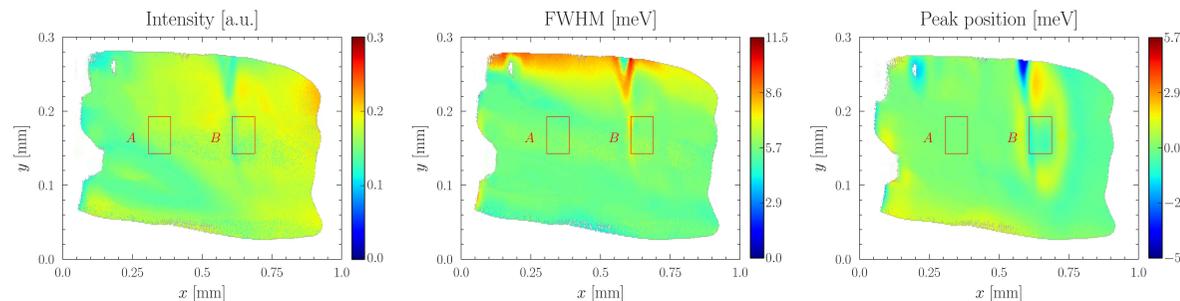
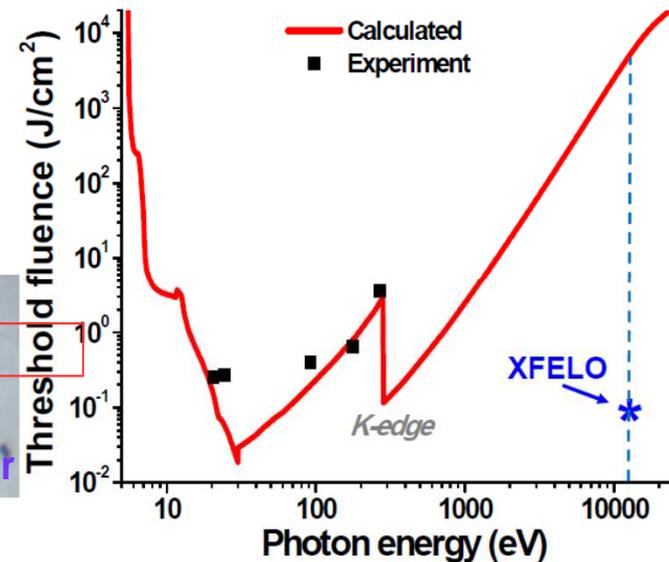
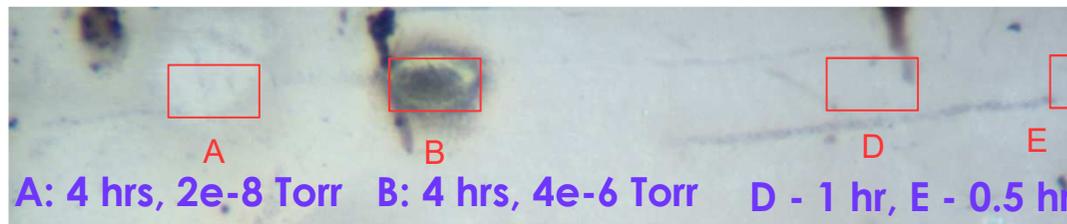
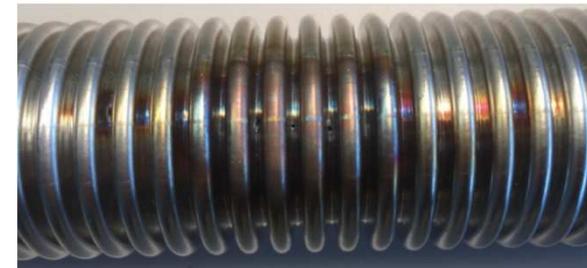


TEST FOR DIAMOND ENDURANCE @ X-RAY POWER DENSITY 10-20 KW/MM²

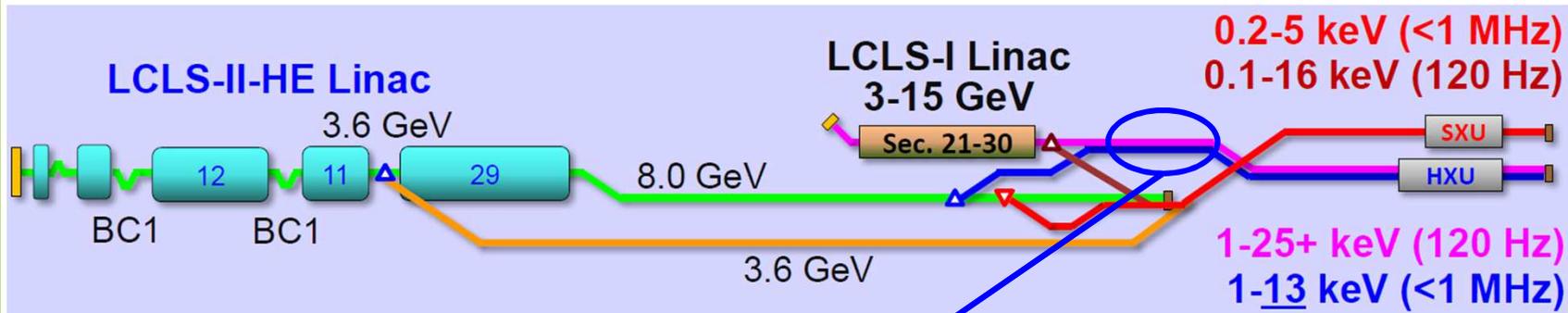
T.

KOŁODZIEJ, ET AL (JSR 2018)

- Steel will melt in < milli-seconds
- Irradiation up to 4 hours at APS
 - 9 kW/mm² in 30x120 μm² spots (K-B mirror focusing) under medium vacuum
 - 12.5 kW/mm² in 30x40 μm² spots (Be-CRL focusing) under UHV (~10⁻⁸)

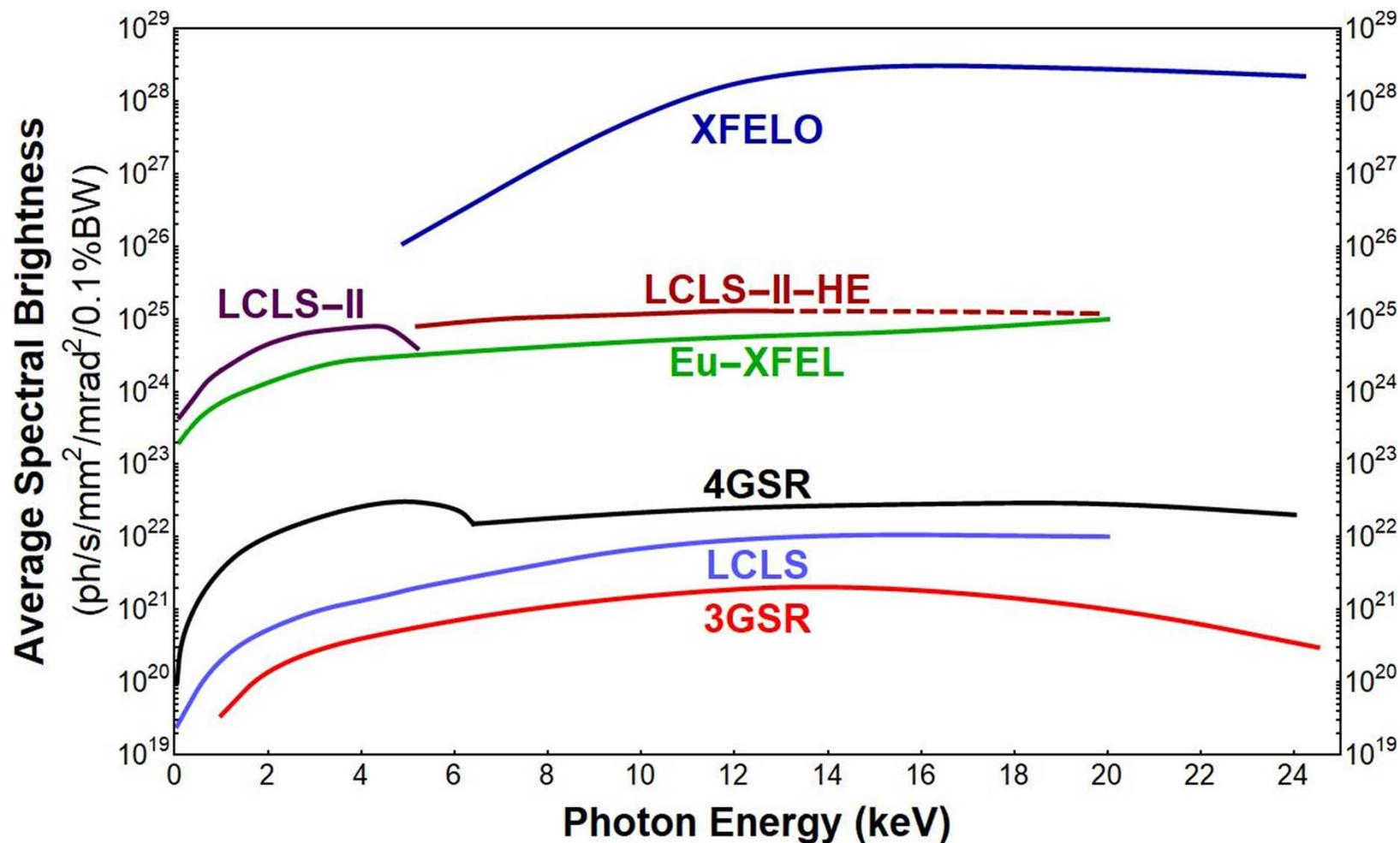


XFELO TEST EXPERIMENT AT LTU, LCLS BEING PLANNED FOR 2019-21



- Use four LCLS undulator units (4x3.4 m)
- Begin with the 120 Hz LCLS linac, but produce a pair of bunches separated by the X-ray cavity round trip time
- The first bunch produce X-ray pulse, and the second pulse amplify
- When the LCLS II SCRF linac available, attempt a steady state operation

AVERAGE BRIGHTNESS OF VARIOUS SOURCES



SUMMARY

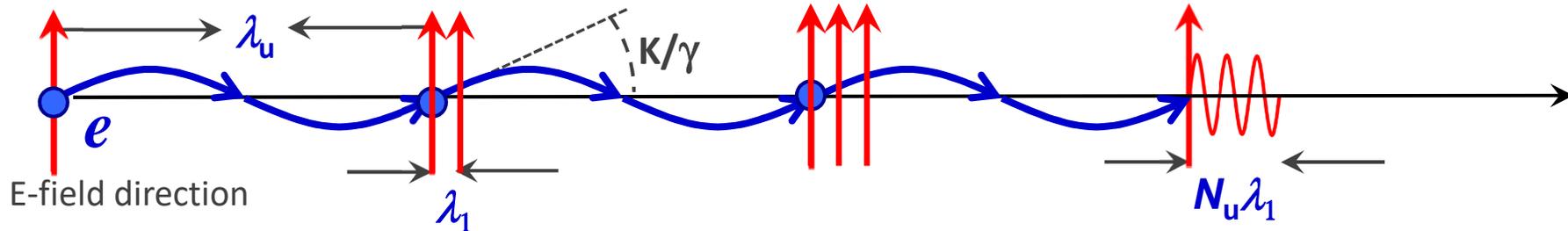
- **We have reviewed the development of high-brightness X-ray generation techniques:**
- **Storage ring-based:**
 - **2nd and 3rd generation SR sources**
 - **DLSR in near future**
- **Linac-based high-gain FELs producing quasi-coherent, femto-second SASE pulses**
 - **Mostly low rep rate XFELs using pulsed linac in the past and current**
 - **High rep rate, even CW MHz, current and future**
 - **Femto-/ atto- second dynamics**
- **X-ray FEL oscillator producing fully coherent X-rays in the future**
 - **Ultra-narrow BW applications**



BACK UP SLIDES

Radiation by one electron in N_u period undulator

- The e^- emits EM wave in the forward direction due to its x-acceleration. Consider the wave fronts from successive undulator periods:



- The e^- is slower since (1) $c > v = c(1-1/2\gamma^2)$, and (2) its trajectory is curved. Thus, the EM wave slips ahead of the e^- in one undulator period by a distance λ_1 =wavelength:

$$\lambda_1 = \lambda_u(1+K^2/2)/2\gamma^2 \quad , \quad \varepsilon_1[\text{keV}] = 12.4/\lambda_1[\text{\AA}]$$
- N_u periods of the undulator $\rightarrow N_u$ cycle wave-train

MINIMUM RADIATION EMITTANCE

$$\mathcal{E}_\omega(\phi; z) = \frac{1}{\lambda^2} \int d\mathbf{x} e^{-ik\phi \cdot \mathbf{x}} E_\omega(\mathbf{x}; z)$$

$$E_\omega(\mathbf{x}; z) = \int d\phi e^{ik\phi \cdot \mathbf{x}} \mathcal{E}_\omega(\phi; z).$$

$$E(x; 0) = E_0 \exp\left(-\frac{x^2}{4\sigma_r^2}\right) \quad \mathcal{E}(\phi; 0) = \mathcal{E}_0 \exp\left(-\frac{\phi^2}{4\sigma_{r'}^2}\right)$$

$$\sigma_r^2(z) = \langle x^2 \rangle = \frac{\int d\mathbf{x} x^2 |E_\omega(\mathbf{x}; z)|^2}{\int d\mathbf{x} |E_\omega(\mathbf{x}; z)|^2} \quad \sigma_{r'}^2(z) = \langle \phi^2 \rangle = \frac{\int d\phi \phi^2 |\mathcal{E}_\omega(\phi; z)|^2}{\int d\phi |\mathcal{E}_\omega(\phi; z)|^2}$$

Minimum emittance Radiation emittance for
Gaussian amplitude profile

$$\sigma_r \sigma_{r'} = \frac{\lambda}{4\pi} \equiv \varepsilon_r.$$

