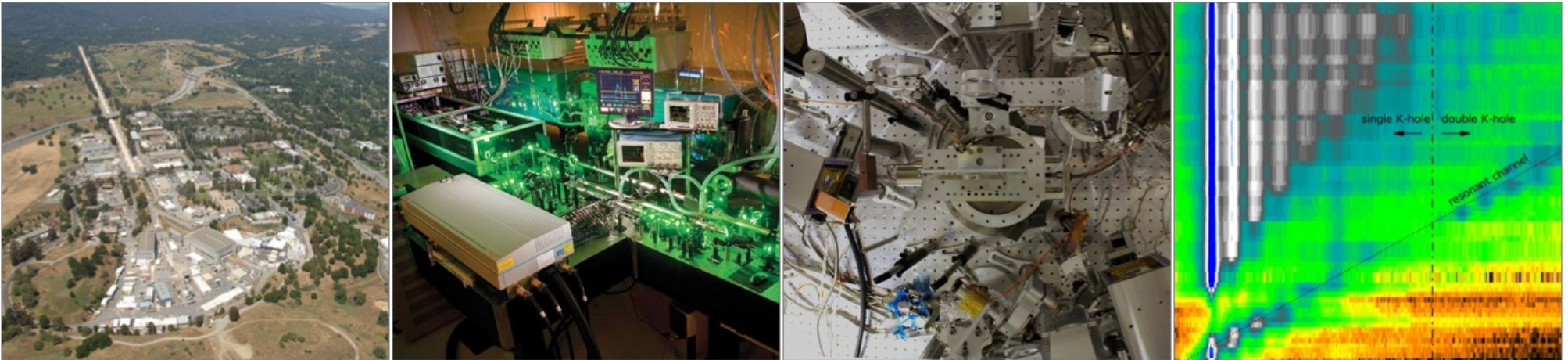


High Energy Density Physics with Extreme Light



Byoung-ick Cho

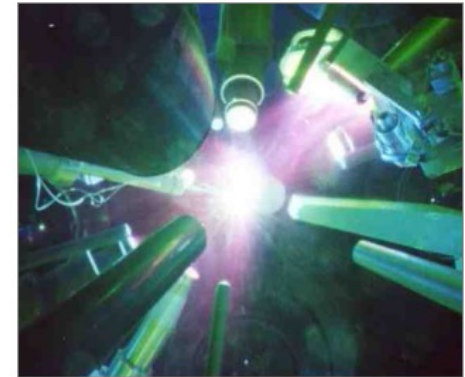
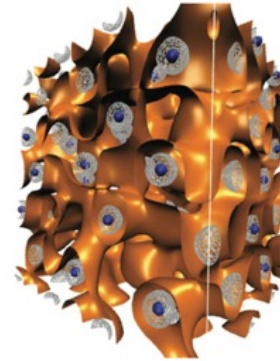
Dept. Physics & Photon Science, GIST

Acknowledgement

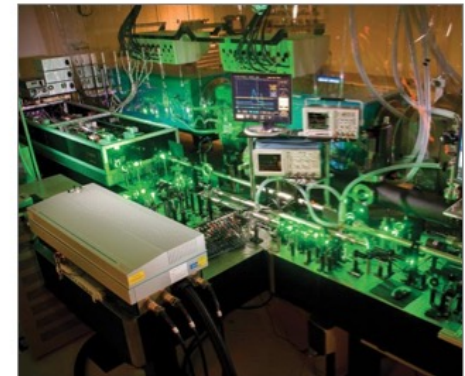


- › GIST: Leejin Bae, Jongwon Lee, Minju Kim, Young Hoon Kim, Minsang Cho, Jaehyung Chung
- › CoReLS - IBS : Chang Hee Nam and CoReLS members
- › ALS / UC Berkeley : Roger Falcone, Richard Lee
- › LCLS : Philip Heimann, Hee Ja Lee, Bob Nagler
- › LLNL : Tadashi Ogitsu, Alfredo Correa, Yuan Ping
- › Oxford: Justin Wark, Sam Vinko, Orlando Ciricosta
- › Univ. Bordeaux/CEA: Jerome Gaudin, Fabien Dorchies

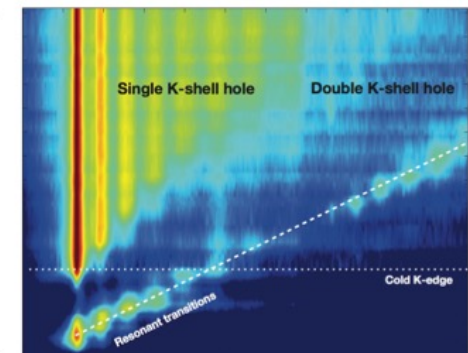
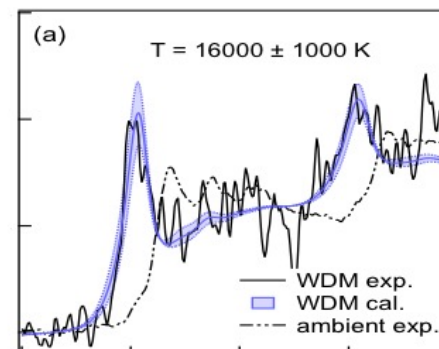
- **What is HEDP?**



- **Extreme Light:**
X-ray Free electron laser & High power laser



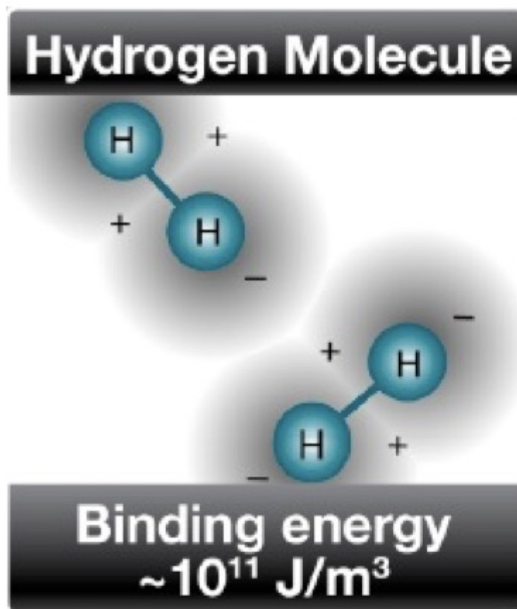
- **Ultrafast X-ray spectroscopies**
for IPD, x-ray transmission,
el-ph coupling of HED matters, etc



High Energy Density Physics (HEDP) can be defined as study of matter at energy density $\sim 10^{11} \text{ J/m}^3$



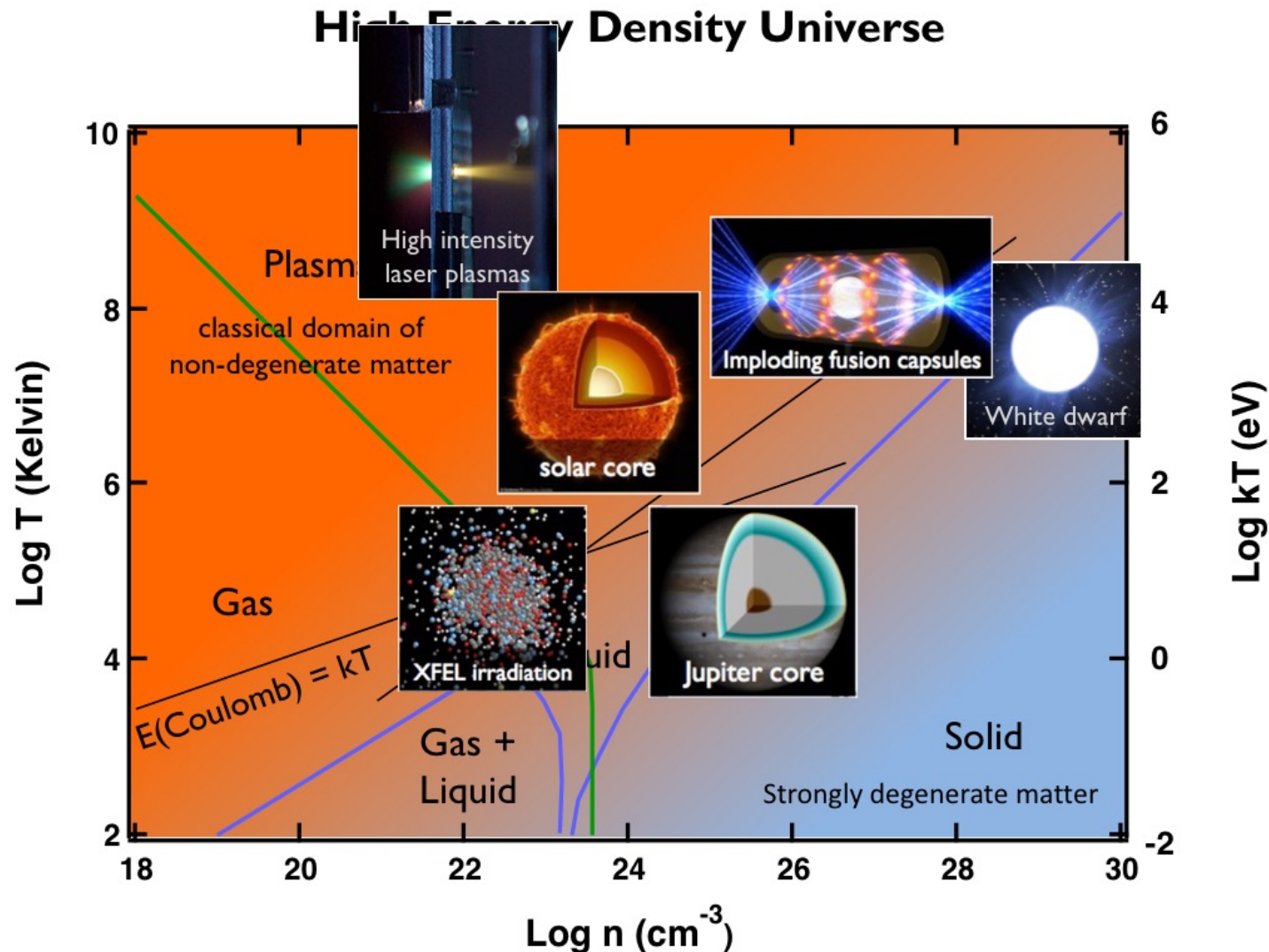
HED conditions



Energy density parameters corresponding to $\sim 10^{11} \text{ J/m}^3$	Values
Pressure	1 Mbar
Electric field strength	$1.5 \times 10^{11} \text{ V/m}$
Magnetic field strength	$5 \times 10^2 \text{ T}$
Laser intensity at 1 μm wavelength	$4 \times 10^{12} \text{ W/cm}^2$
Blackbody radiation temperature	75 eV

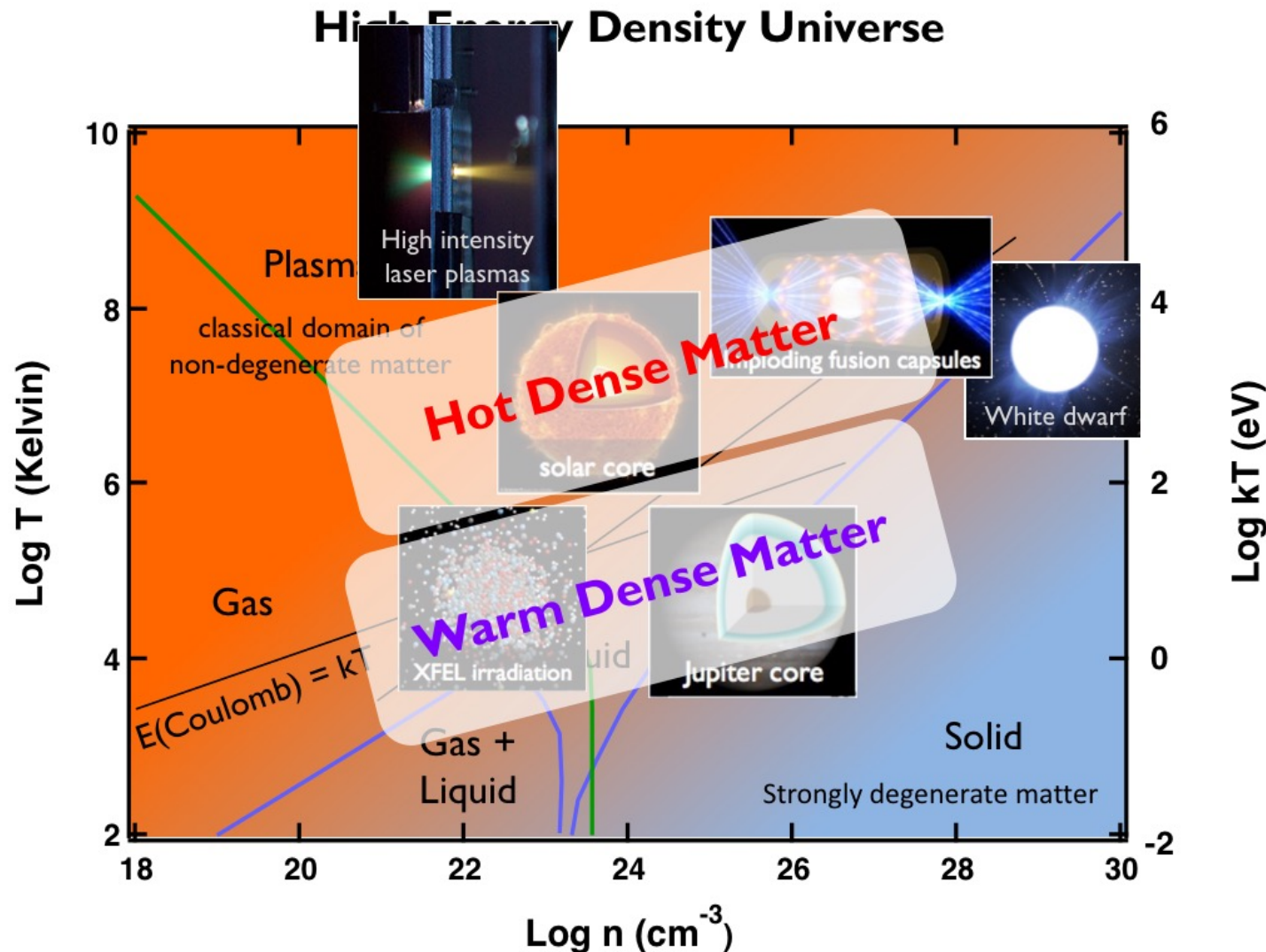
Frontiers in High Energy Density Physics - The X-Games of Contemporary Science
(National Academies Press, 2003)

HED condition is prevalent throughout the Universe



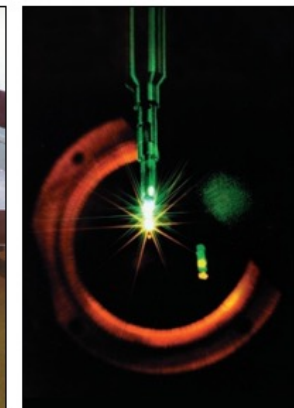
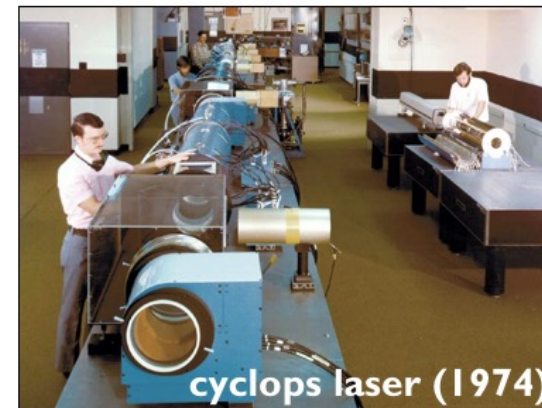
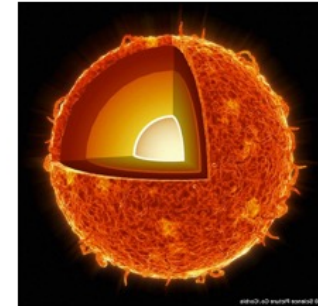
Basic Research Needs for High Energy Density Laboratory Physics, DOE Office of Science and NNSA (2010)

The complicated interplay of the physical processes creates a malfunction junction for description of states between solid and plasmas



Precursors to high energy density physics

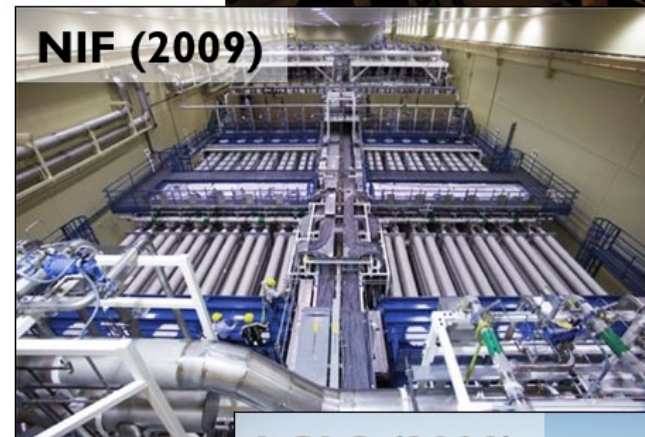
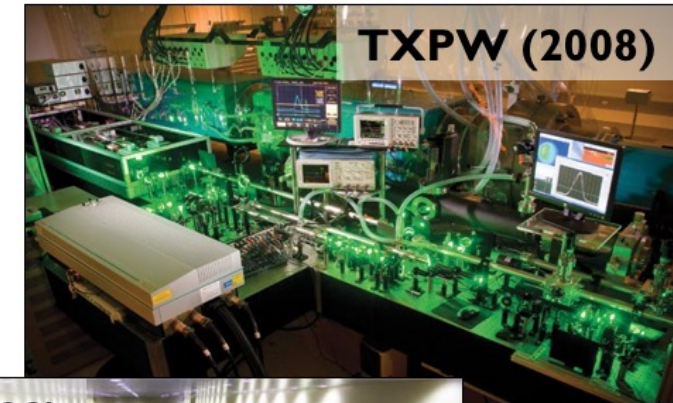
- ▶ **First half 20th Century: Stellar structure**
 - Eddington, Chandrasekhar, Schwarzschild, among others
- ▶ **Mid-20th Century: Nuclear weapons**
 - Oppenheimer, Sakharov, Teller, Bethe, Fermi, others: Compressible metals
 - Zel'dovich and Raizer: *Physics of Shock Waves and High Temperature Phenomena* (1966)
- ▶ **Post Mid-20th Century (1960-1980): Inertial confinement fusion origins**
 - Lawrence Livermore National Lab



Post 1980s: Extreme light in the laboratory

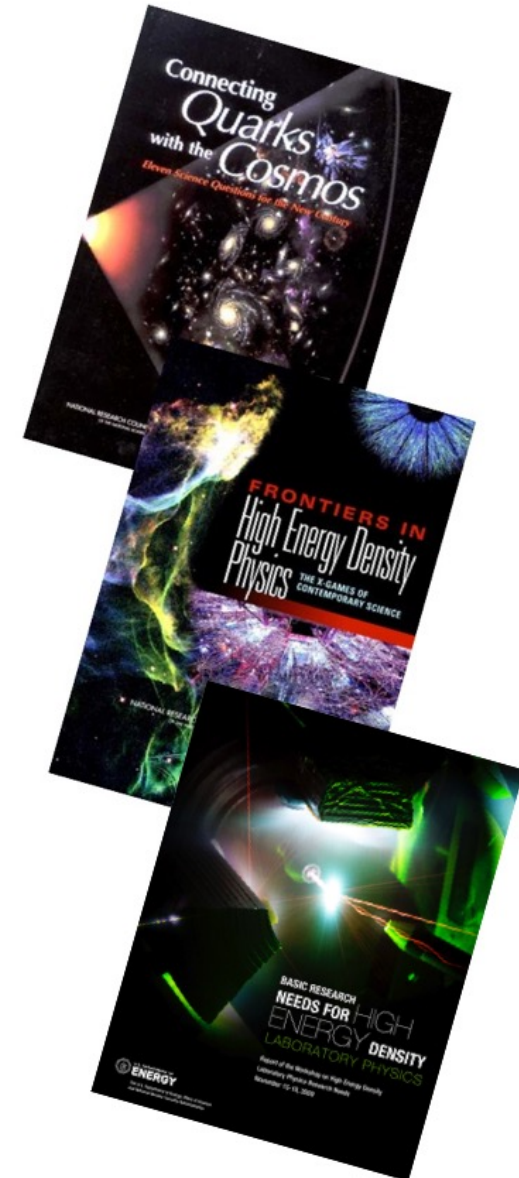
Complex quantitative *physics* experiments became feasible

- › **1984: Chirped Pulse Amplification**
- Gerard Mourou and Donna Strickland
- › **1990s: Tabletop Terawatt class lasers**
- 10^{12} Watt in laboratories, < 100 fs pulses
- › **late 2000s: Petawatt lasers**
- 10^{15} Watt, U.Texas, GIST, etc
- › **2009: National Ignition Facility (NIF)**
- LLNL
- › **2009: X-ray free electron laser**
- LCLS @ SLAC, 10^9 brighter than any other x-ray sources



21st centry: HEDP as a discipline

- **1995: 1st comprehensive report on HEDP**
(energy density $> 10^{11}$ J/m³)
 - *Science on High Energy Lasers* (Lee, Petrasso, & Falcone)
- **2003: Two NAS reports highlighted HEDP**
 - *Connecting Quarks with the Cosmos: Eleven Science Questions for the New Century*
 - *Frontiers in High Energy Density Physics: The X-Game of Contemporary Science*
- **2009, 10: DOE reports for HED experiments**
 - *Advancing the science of High Energy Density Laboratory Plasma*
 - *Basic Research Needs for High Energy Density Laboratory Physics*

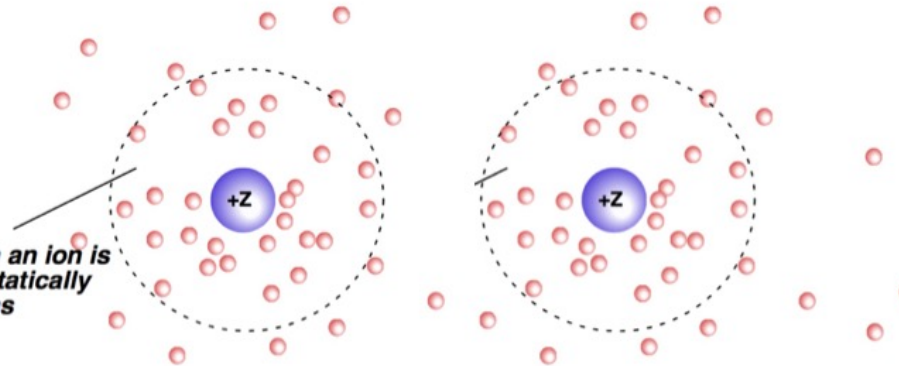


From a "plasma point-of-view", HED matter can not be described by classical, two-body interactions

Debye length

$$\lambda_D = \sqrt{\frac{kT_e + kT_i}{4\pi e^2 Z n_o}}$$

Distance over which an ion is a plasma is electrostatically shielded by electrons



Plasma parameter

$$\Lambda_p = \frac{1}{\sqrt{Z n_o}} \left(\frac{kT_e}{4\pi e^2} \right)^{3/2}$$

of e- within a sphere with radius of the Debye length

Strong coupling parameter

$$\Gamma = \frac{V_{ii}}{k_B T} = \frac{Z^2 e^2}{r_o k_B T}$$

Ratio of the interaction energy between particles V_{ii} to the kinetic energy T

Tokamak plasma

$$\begin{aligned} kT_e &= 10 \text{ keV} \\ n_e &= 10^{14} \text{ cm}^{-3} \\ \Lambda_p &= 4 \times 10^7 \\ \Gamma &\ll 1 \end{aligned}$$

HED plasma

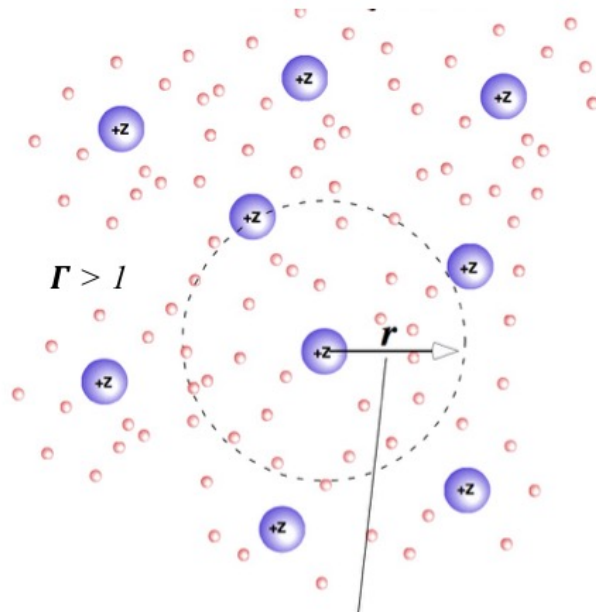
$$\begin{aligned} kT_e &= 100 \text{ eV} \\ n_e &= 10^{22} \text{ cm}^{-3} \\ \Lambda_p &= 4 \\ \Gamma &> 1 \end{aligned}$$

Some cases, $\Lambda_p < 1$

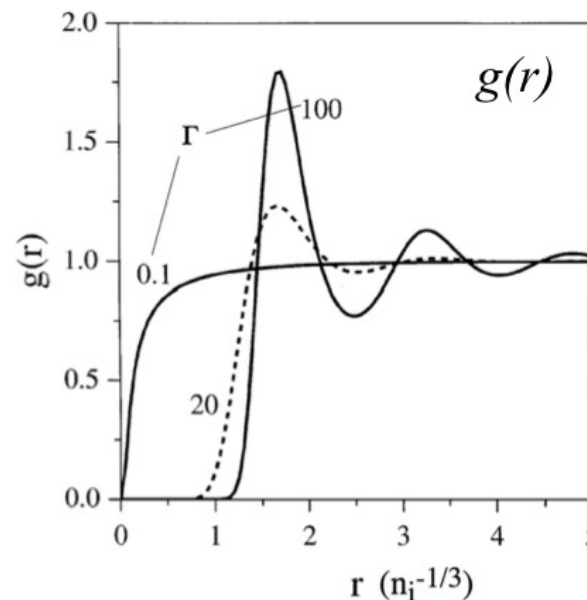
Classic Cluster expansion (BBGKY hierarchy) leading to usual kinetic equations (Vlasov; Fokker Plank) is not valid

Quantum degeneracy effects & ion correlation can become important in some HED matter

Coulomb force lead to correlations in radial position



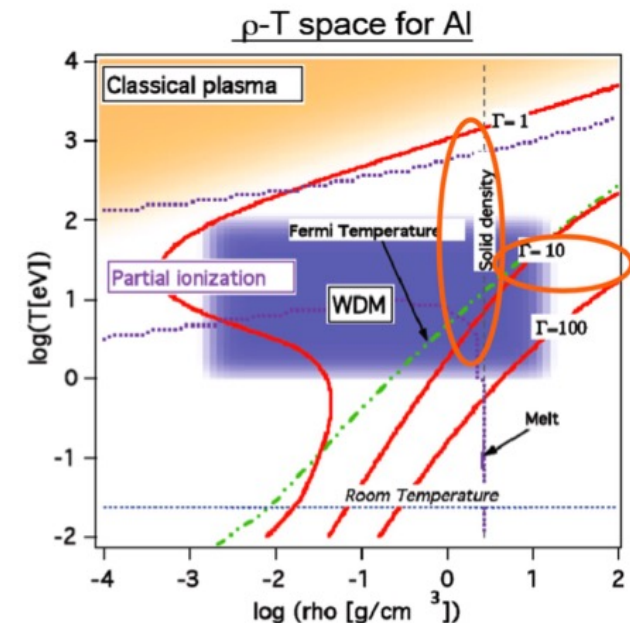
Radial distribution function of ions $g(r)$: prob of finding ion at radius r



Reproduced from Murillo and Weisheit Phys. Rep. 302, 1 (1998)

Too far from ground state

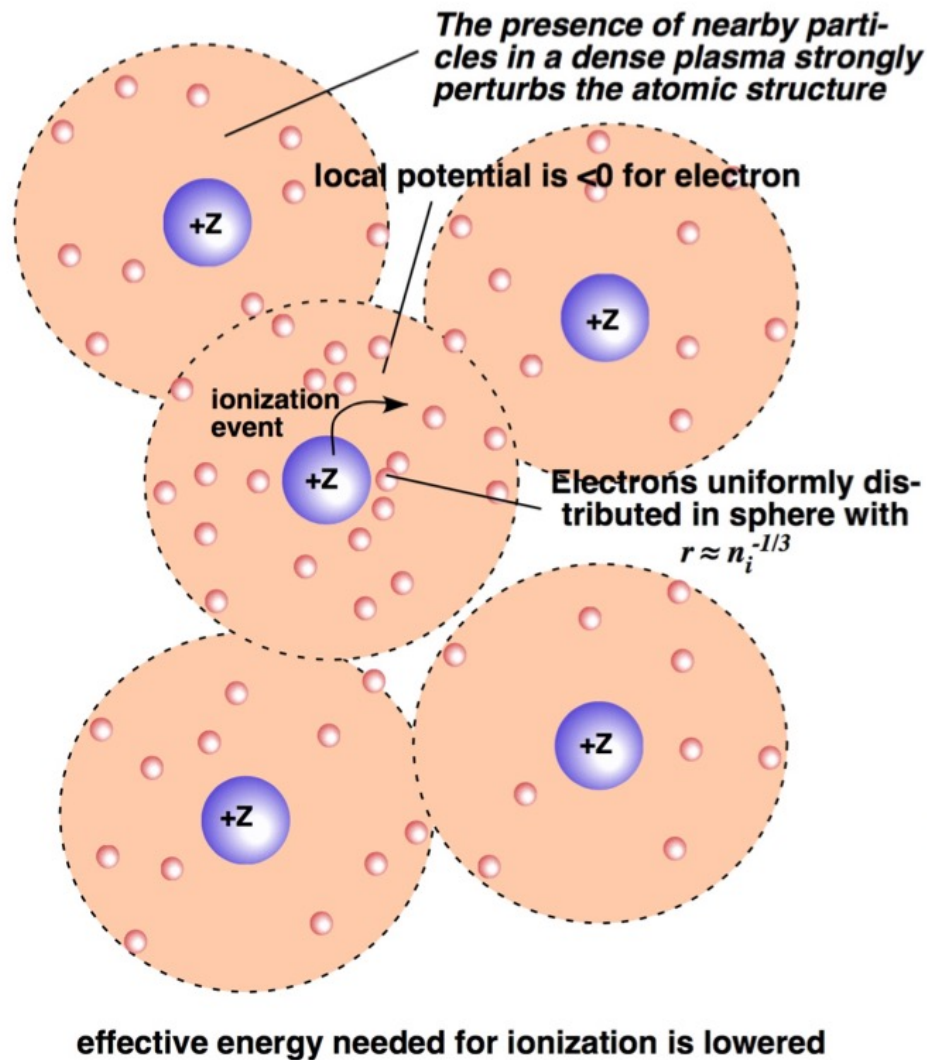
$$kT \sim E_F$$



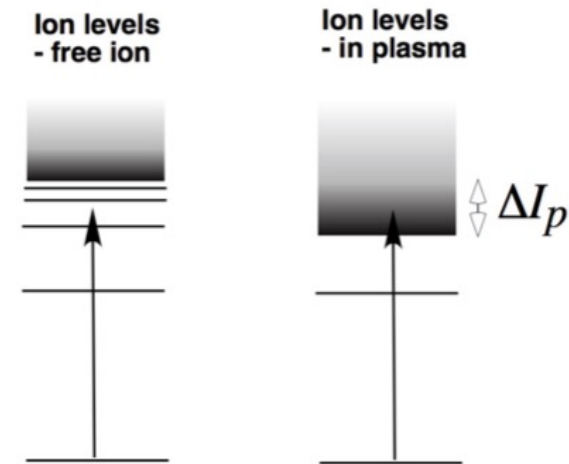
A strongly coupled plasma will exhibit ion correlations and structure akin to solid ~ liquid

At solid density, Fermi energy is ~ 10 eV. Many HED matter have temperature in this vicinity, mandating a QM treatment of the plasma

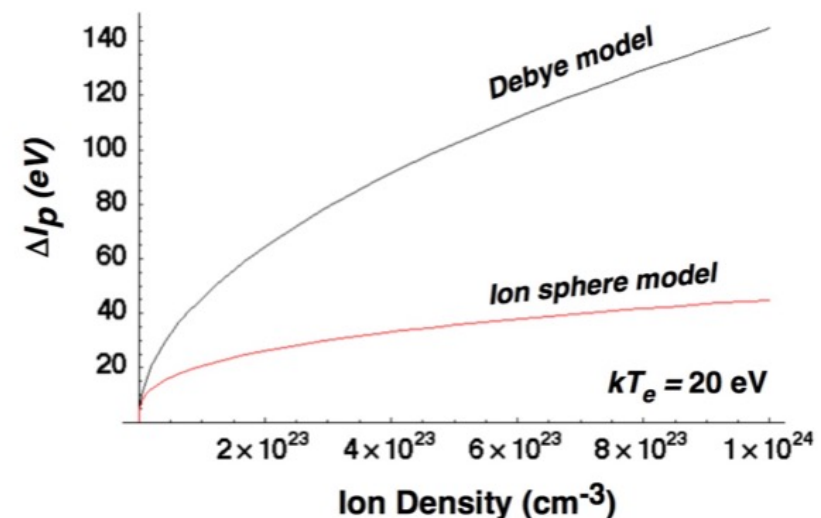
A significant question in dense plasmas is the extent to which the ionization potential of the ions is altered



One significant effect is the lowering of the ionization potential: "Continuum lowering"



Continuum lowering predictions with a Saha model for ionization



Ionization models do not converge in the HED plasma regime



Saha equilibrium equations for ion charge states

- Low to moderate density
- High temperature

$$\frac{n_e n_Z}{n_{Z-1}} = 2 \frac{W_Z(T)}{W_{Z-1}(T)} \left(\frac{m_e k T_e}{2\pi \hbar^2} \right)^{3/2} \exp \left[-\frac{I_p^{(Z-1)} - \Delta I_p^{(CL)}}{k T_e} \right] \quad n_e = \sum_{j=1}^{Z_{nuc}} n_j$$

“Average Atom” picture based on Thomas Fermi Model

- High density
- Low to moderate temperature

$$\frac{d^2 \chi}{dr^2} = \frac{4}{3\pi \hbar^3} (2m_e e^2)^{3/2} \frac{Z^{1/2}}{r^{1/2}} \chi^{3/2}$$



These models are complicated by many other factors

Co-existing constituents in a dense plasma:

- Ions
- Neutrals
- Negative ions
- Dimers
- Negatively charged dimers

Electron affinity of Au: 2.3 eV

Bond strength of Al₂: ~ 5 eV

When steep temperature gradients exist in hot dense plasmas, the heat transport becomes nonlocal

Moderate temperature plasmas/soft temperature gradients:

The heat flux is given by a local formula

$$\longrightarrow Q_{\text{Diffusive}} = \kappa_{SH} \nabla T_e$$

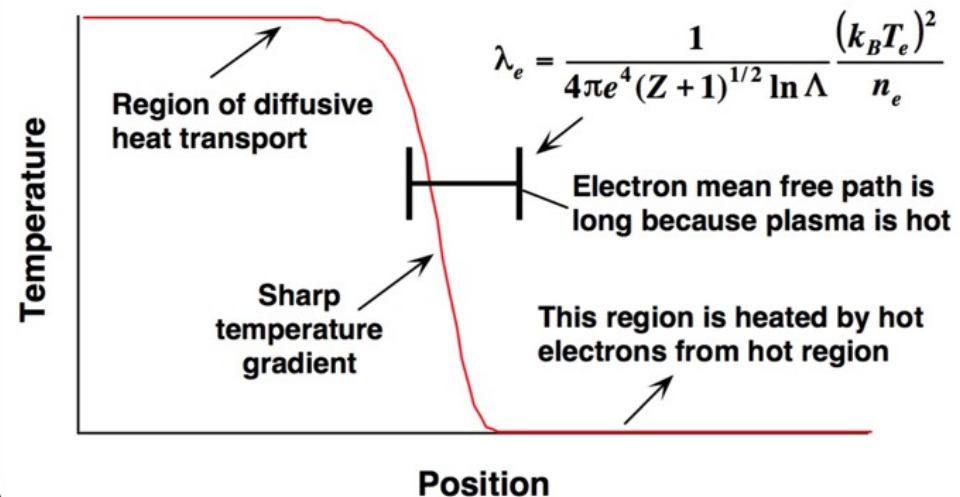
Spitzer-Härm conductivity

Sharp gradients: heat flux exceeds the maximum possible free streaming flux

$$\longrightarrow Q_{\text{max}} = \frac{3}{2} n_e k_B T_e v_{\text{therm}}$$

electron thermal velocity

When the mean free path of electrons exceeds the temperature gradient, standard diffusive heat transport breaks down



Heat transport in a sharp gradient:
Nonlocal heat flux formula

$$q_{\text{nonlocal}}(x) = \int q_{\text{Spitzer}}(x') G(x, x') dx' \longrightarrow G(x, x') = \frac{1}{2\lambda_e} \exp\left[-\frac{|x - x'|}{\lambda_e}\right]$$

Typical form of kernel $G(x, x')$ *

*Luciani, Mora, and Virmont, Phys. Rev. Lett. 51, 1664 (1983)

In many HED plasmas:

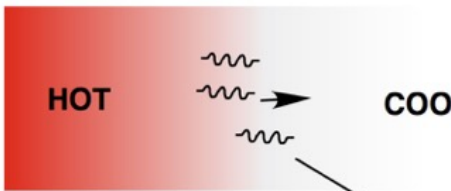
$$T_e \sim 10 - 100 \text{ eV}$$

$$n_e \sim 10^{21} - 10^{23} \text{ cm}^{-3}$$

$$\delta x \sim 1 - 100 \text{ } \mu\text{m}$$

$$\lambda_e \sim 1 - 10 \text{ } \mu\text{m}$$

Radiation energy flow in HED plasmas is an extremely complex phenomenon



Energy flow by photons

Radiative heat flux

Rossland Mean Free Path: effective photon propagation distance

$$Q_{rad} = -16 \frac{\sigma_{SB} \lambda_R T^3}{3} \nabla T$$

$$\lambda_R = \frac{15h^5}{4\pi^4 (k_B T)^5} \int_0^\infty \lambda(\nu) \frac{\nu^4 e^{-h\nu/k_B T}}{(1 - e^{-h\nu/k_B T})^3} d\nu$$

Photon frequency-dependent mean free path

Photon Absorption Mechanisms

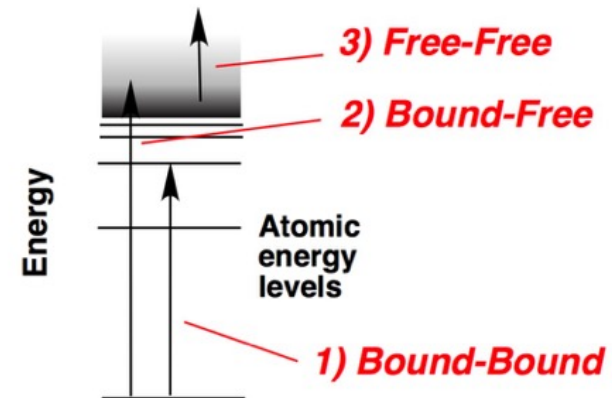
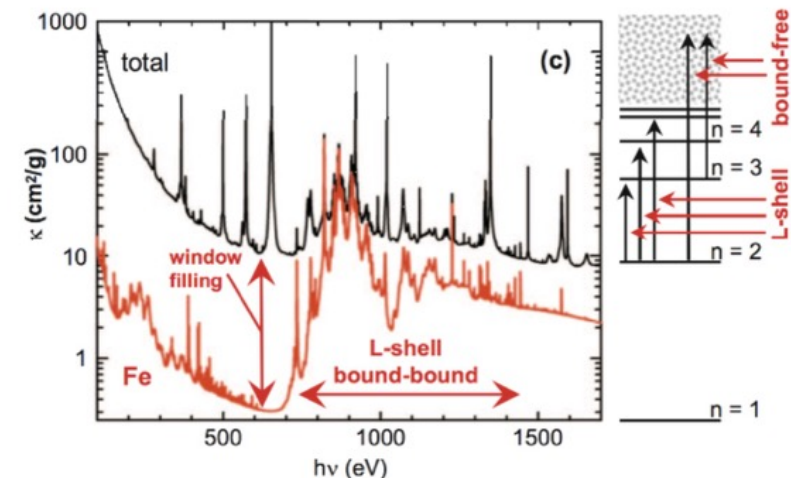


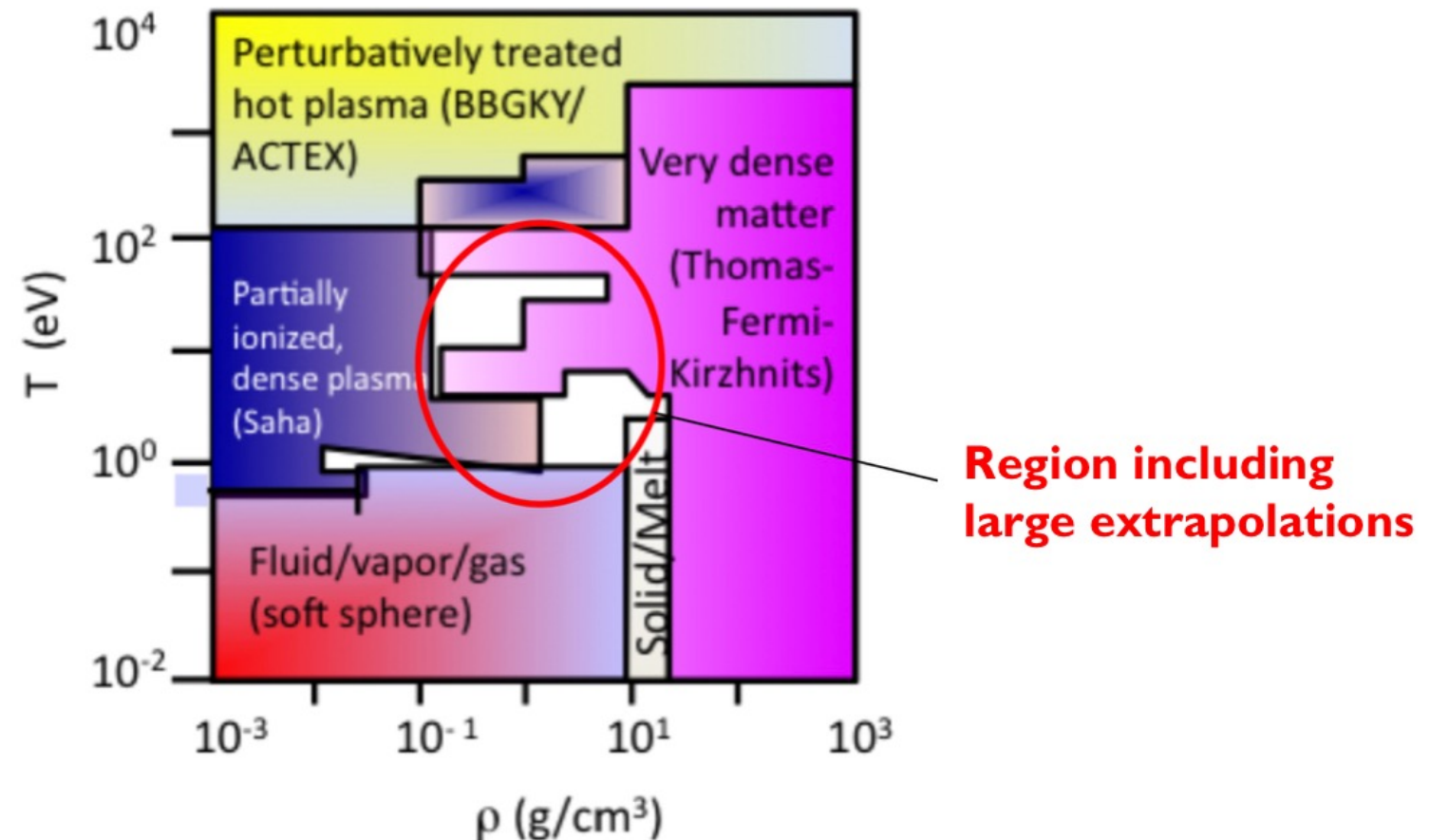
Photo-absorption spectrum of dense Fe



J. E. Bailey et al. Phys. Plas. **16**, 058101 (2009)

The common theoretical practice is to piece together many different models in the Warm, Hot Dense Matters

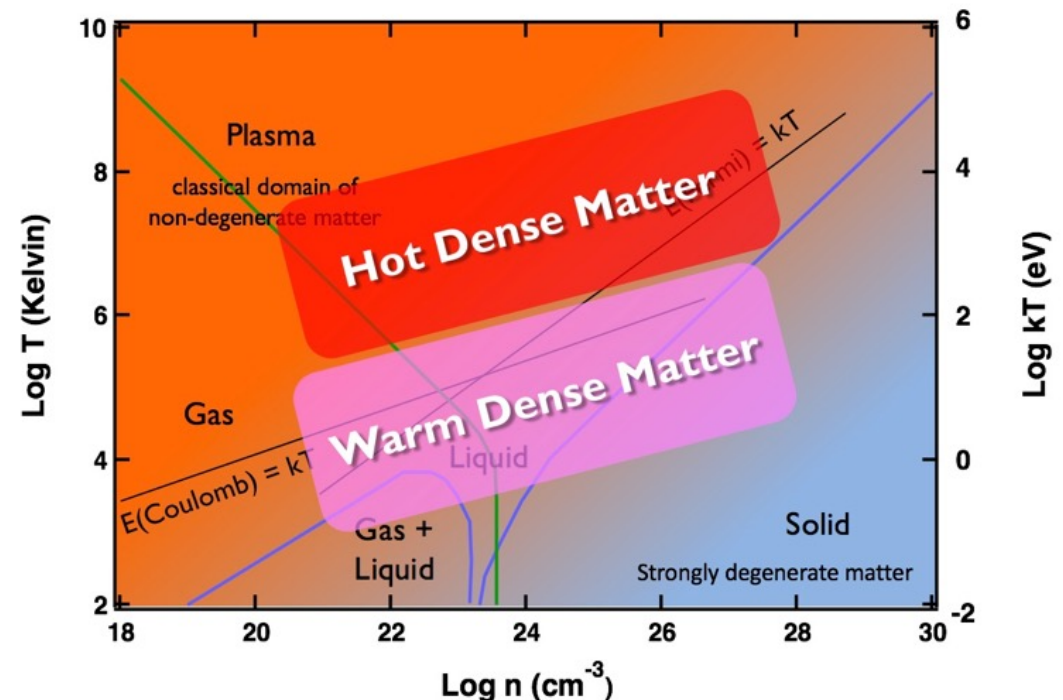
ρ -T diagram (copper) :
The range of application for the best theories in each region



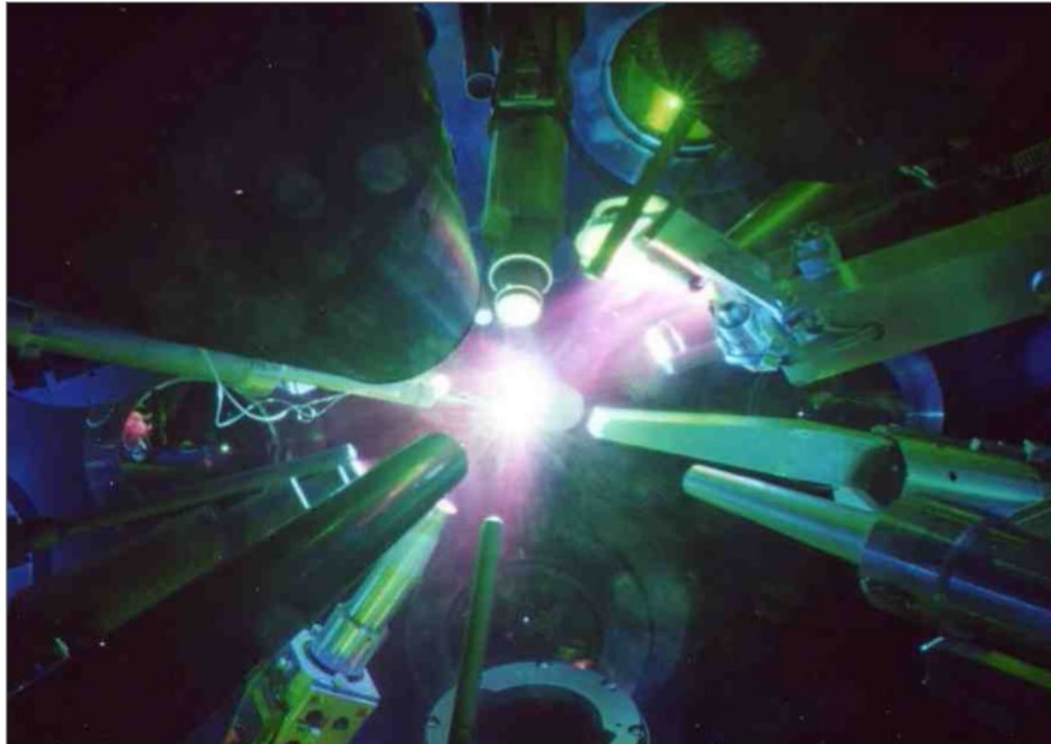
The challenge of HED science centers on difficulties in describing the basic properties of an HED substance

Research Needs

- ▶ **Equation of state:** Measurement of temperature is essential but has proven to be extremely difficult
- ▶ **Material structure:** Measurements of ionic and electronic structures provide significant insight and are necessary to understand the phase of the material and its transport and mechanical properties
- ▶ **Transport properties:** A critical need is to develop measurements of transport properties, such as electrical and thermal conductivity.
- ▶ **Approach to equilibrium:** Definitive time-dependent information is lacking, and successful collection of relevant data would represent a major breakthrough.
- ▶ and more...



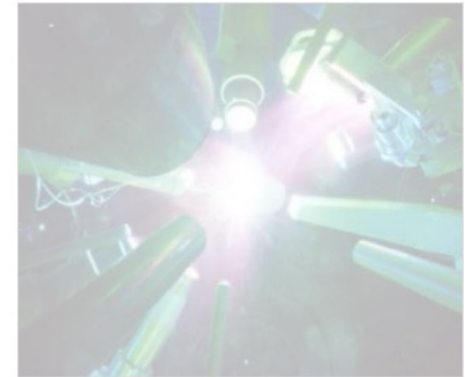
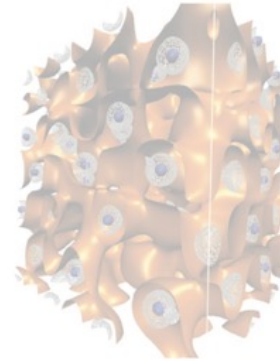
A major hurdle in understanding HED matter is the lack of single-state data



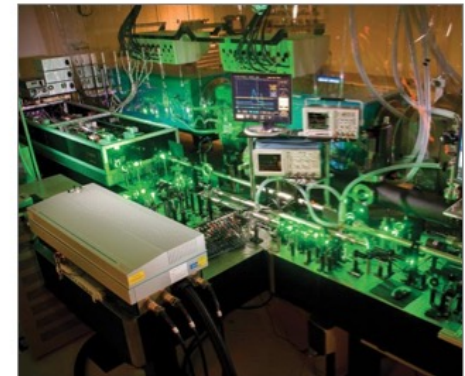
- **Short-living, Fast-varying:** Laboratory HED matter tends to be *non-uniform* due to hydrodynamic expansion at extreme pressure. Properties measured on non-uniform or *multi-state* systems can only be compared with theory through code simulations that includes modeling of *gradient effects*

Experimental requirement :
Ultrafast and Intense Light Sources

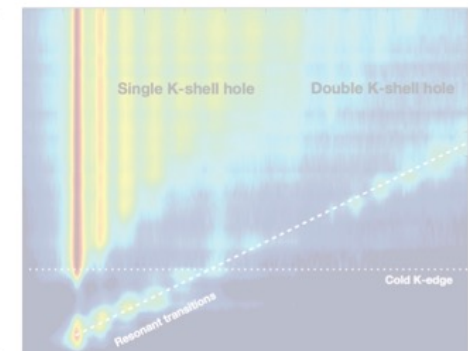
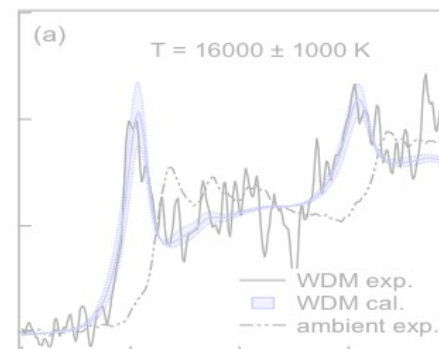
- What is HEDP?



- Extreme Light :**
X-ray Free electron laser & High power laser



- Ultrafast X-ray spectroscopies**
for IPD, x-ray transmission,
el-ph coupling of HED matters, etc



Unprecedentedly extreme power & brightness at optical and x-ray regime in femtosecond time scales



X-ray Free Electron Laser (XFEL)



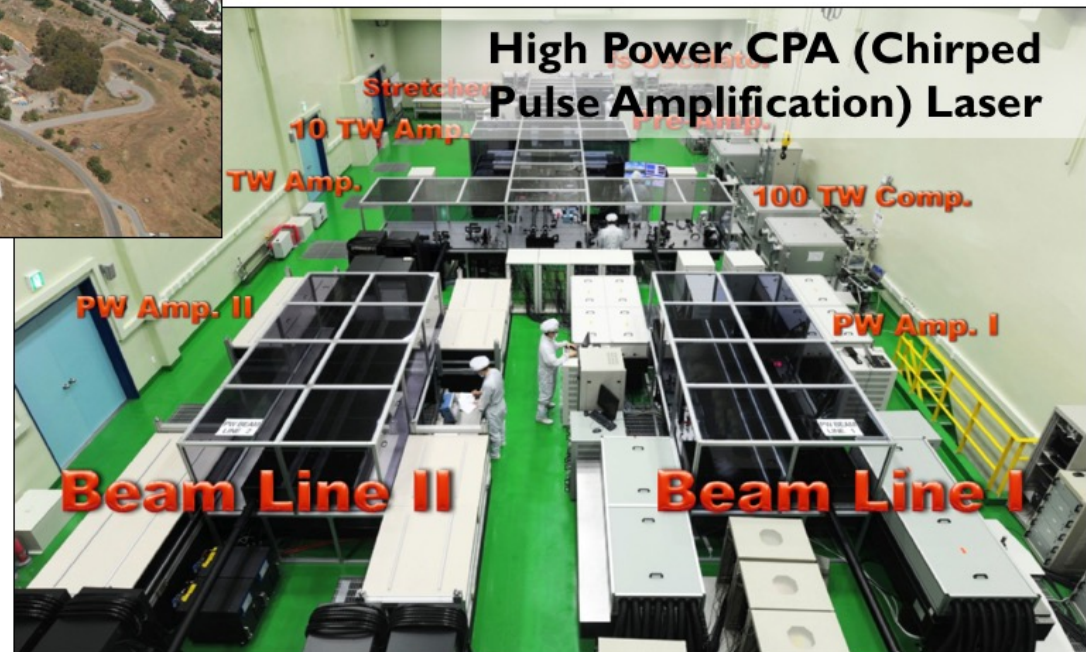
PW laser @ GIST

- λ : 0.8 μm
- pulse length: 35 fs
- pulse energy : > 40 J
- focusing to achieve IR intensities $> 10^{21} \text{ Wcm}^{-2}$

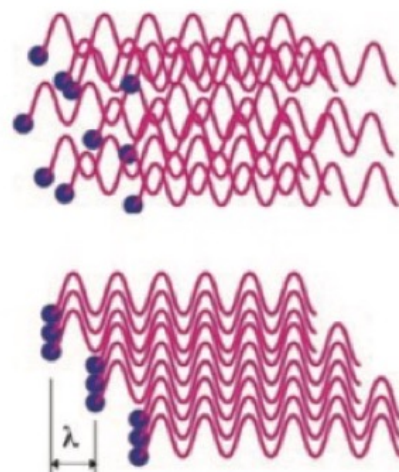
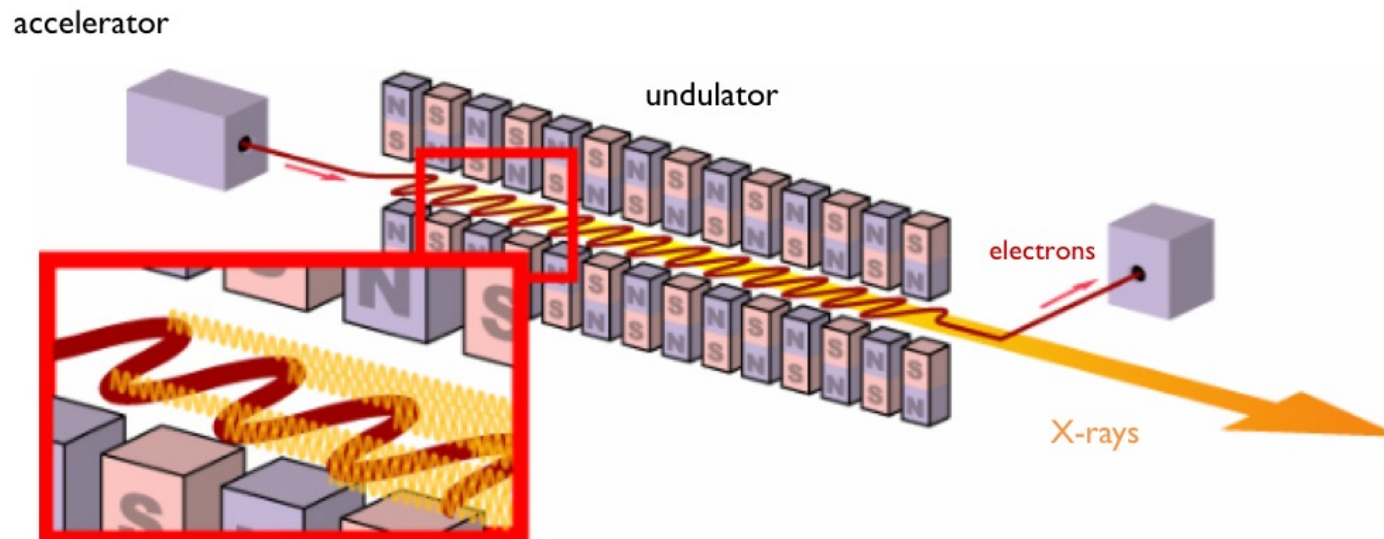
LCLS @ SLAC

- $h\nu$: 250 eV – 10.5 keV
- pulse length: <5 fs - 500 fs
- pulse energy : < 4.7 mJ
- focusing to achieve X-ray intensities $> 10^{18} \text{ Wcm}^{-2}$

High Power CPA (Chirped Pulse Amplification) Laser



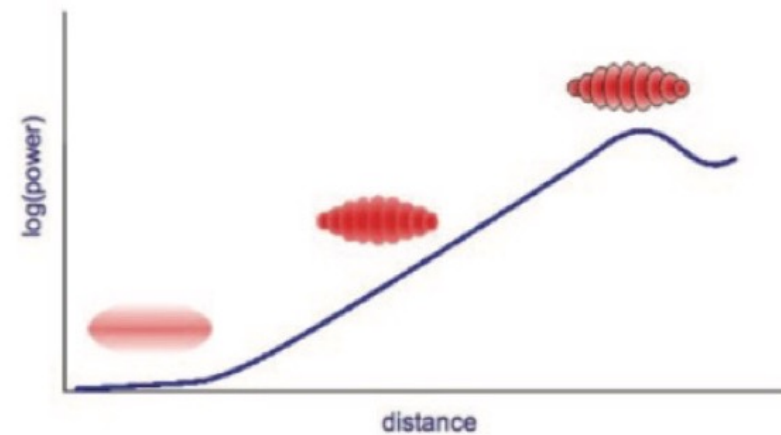
What is X-ray Free Electron Laser (XFEL) ?



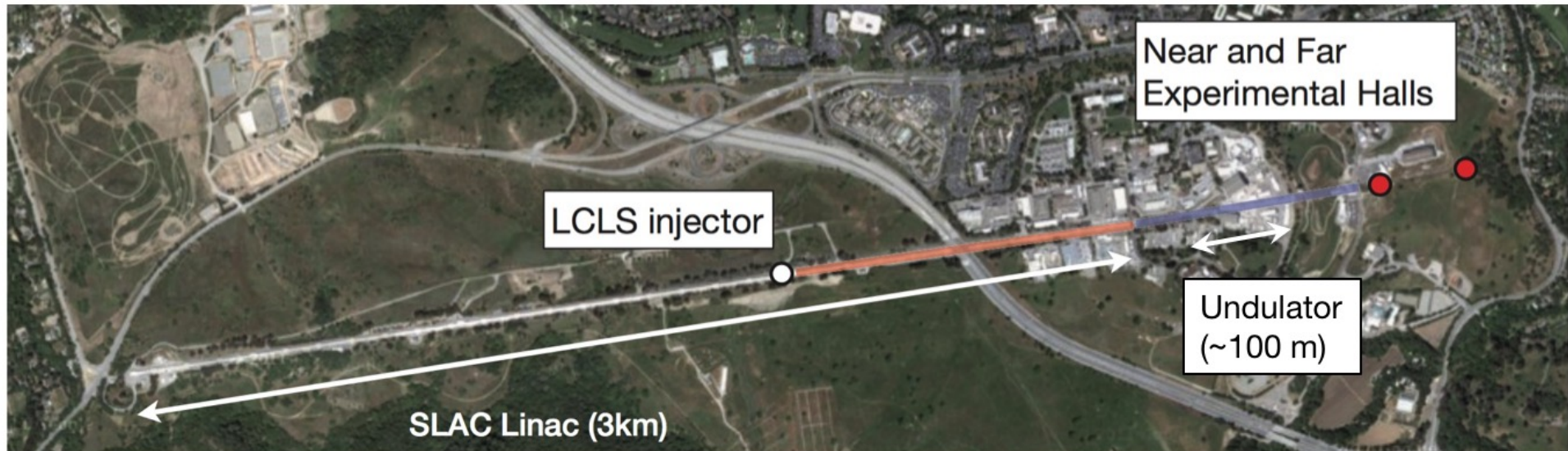
$$P = N P_1$$

$$E = N E_1$$

$$P = N^2 P_1$$

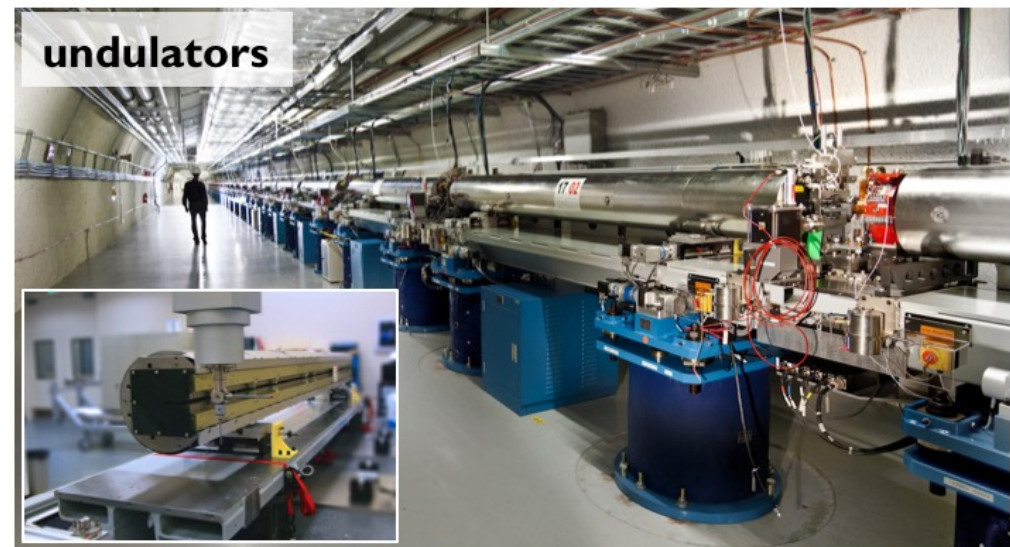


Linac Coherent Light Source (LCLS) @ Stanford Linear Accelerator Center (SLAC)



World's first XFEL

- ▶ tunable photon energy range:
250 eV – 10.5 keV (2.4 – 15.4 GeV e-beam)
- ▶ pulse repetition rate: 120 Hz
- ▶ pulse energy: < 4.7 mJ
- ▶ bandwidth: 0.3%
- ▶ variable pulse length: <5 fs - 500 fs
- ▶ focusing to achieve X-ray intensities
> 10^{18} Wcm^{-2}



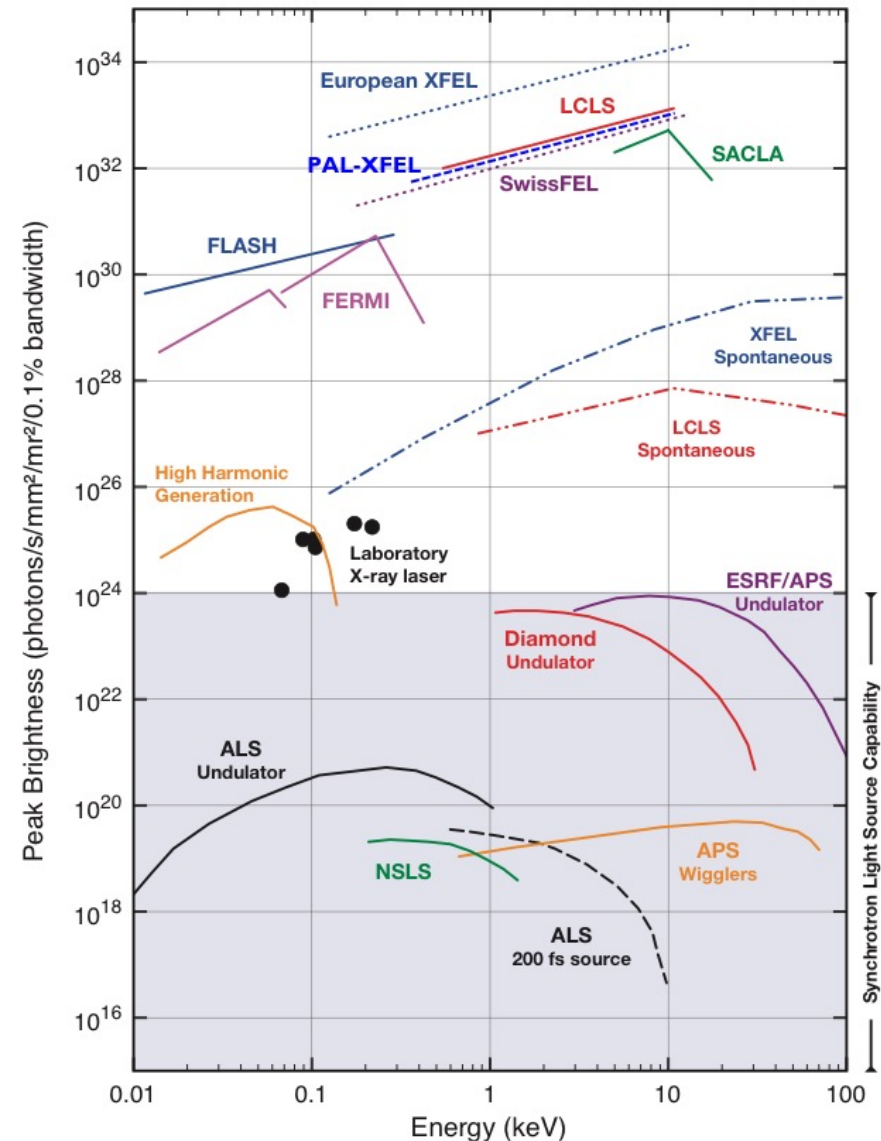
XFEL provides tremendous possibilities for HEDP

Peak brightness 9 orders of magnitude higher than 3rd gen. sources

- ▶ Short pulse duration (1–100 fs)
- ▶ Full transverse coherence
- ▶ High repetition rate
- ▶ Tunable
- ▶ High # of photons / bunch ($\geq 10^{12}$)

X-ray FELs worldwide

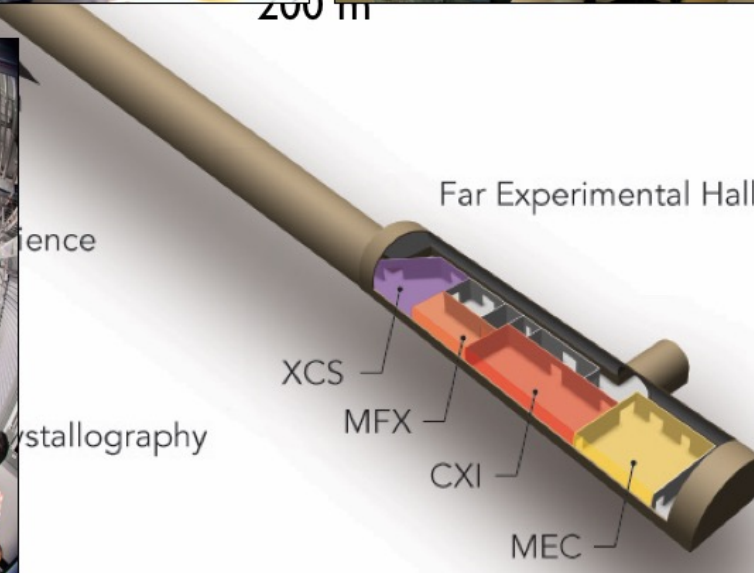
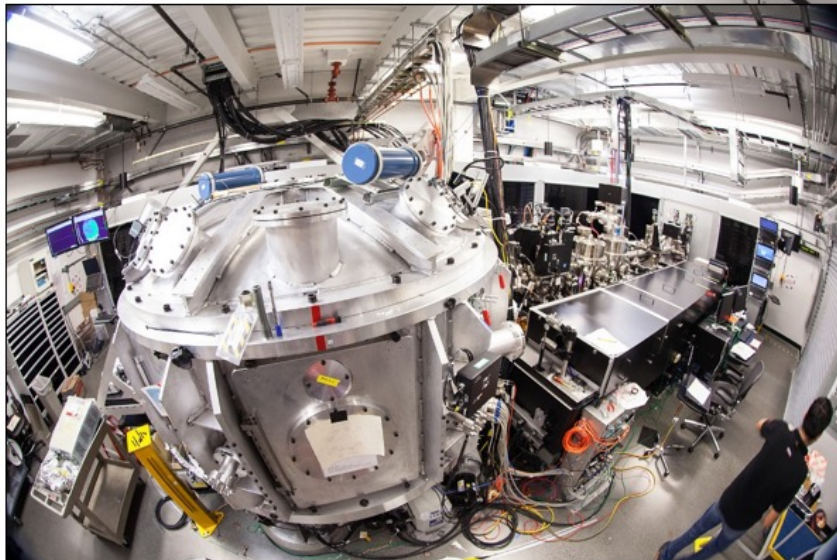
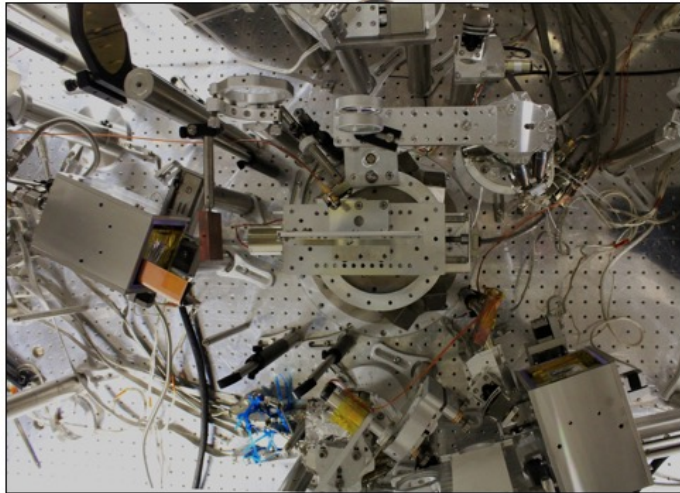
- ▶ LCLS, USA : 0.12 ~ 2.5 nm
- ▶ SACLA, Japan : 0.1 ~ 3.6 nm
- ▶ PAL-XFEL, Korea : 0.06 ~ 4.5 nm
- ▶ Euro X-FEL, Germany : 0.1 ~ 6 nm
- ▶ Swiss-FEL, Swiss : 0.1 ~ 7 nm



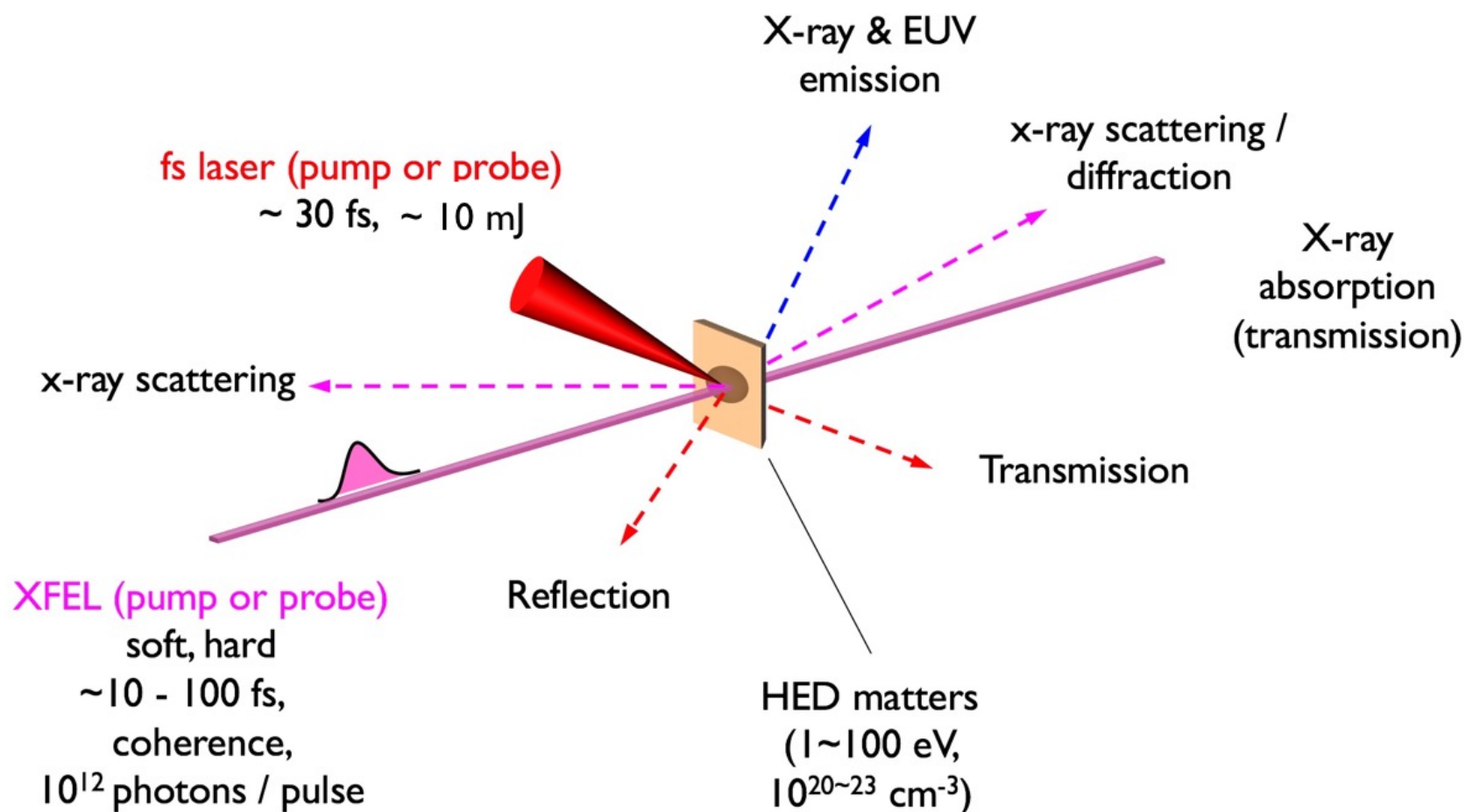
XFEL worldwide



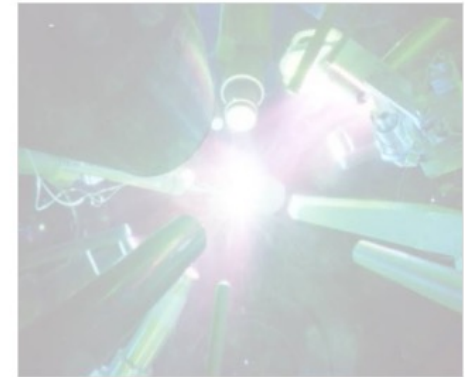
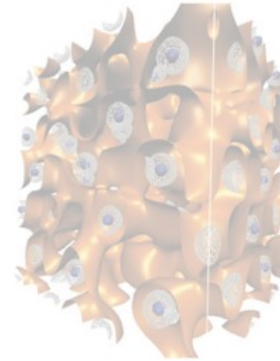
LCLS experimental hutches & control rooms



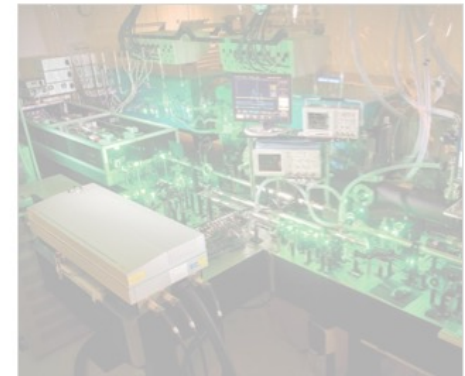
HED experiment on XFEL: Concept



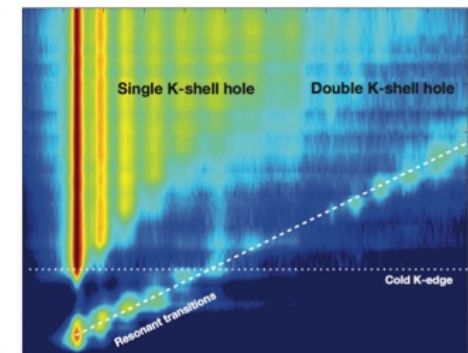
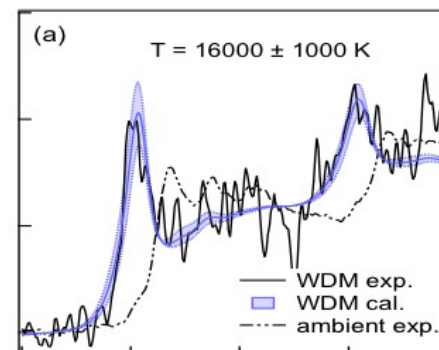
- What is HEDP?



- Extreme Light :**
X-ray Free electron laser & High power laser



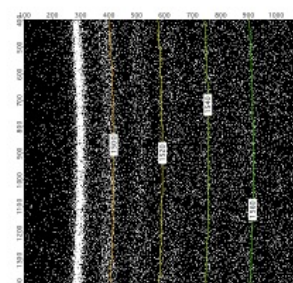
- Ultrafast X-ray spectroscopies**
for IPD, x-ray transmission,
el-ph coupling of HED matters, etc



K-shell spectroscopy of hot dense aluminum

Isochoric heating of solid sample with an intense X-ray pulse

X-ray spectrometer : K-alpha emission Al around 1200~1900 eV



x-ray CCD

PD

Bragg crystal

LCLS pulse

Photon energy : 1480 ~ 1850 eV

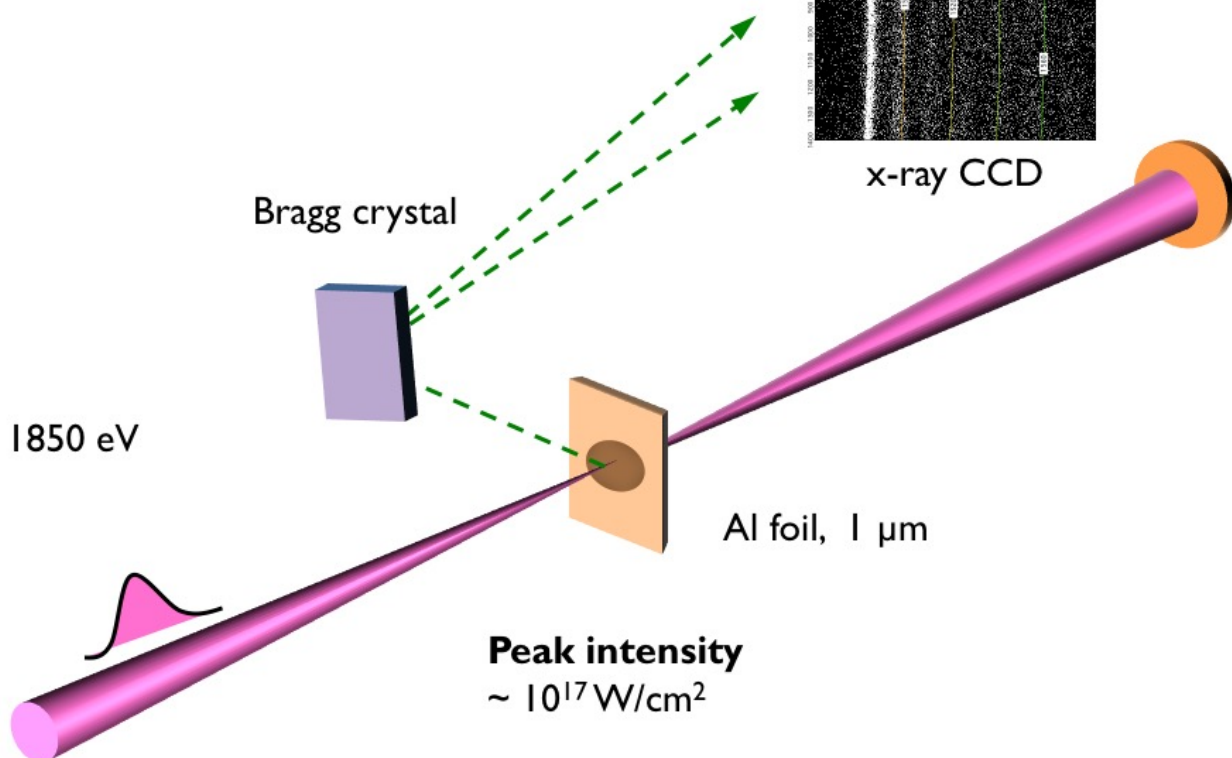
Pulse length < 100 fs

Pulse energy ~ 2 mJ

Bandwidth ~ 0.3 %

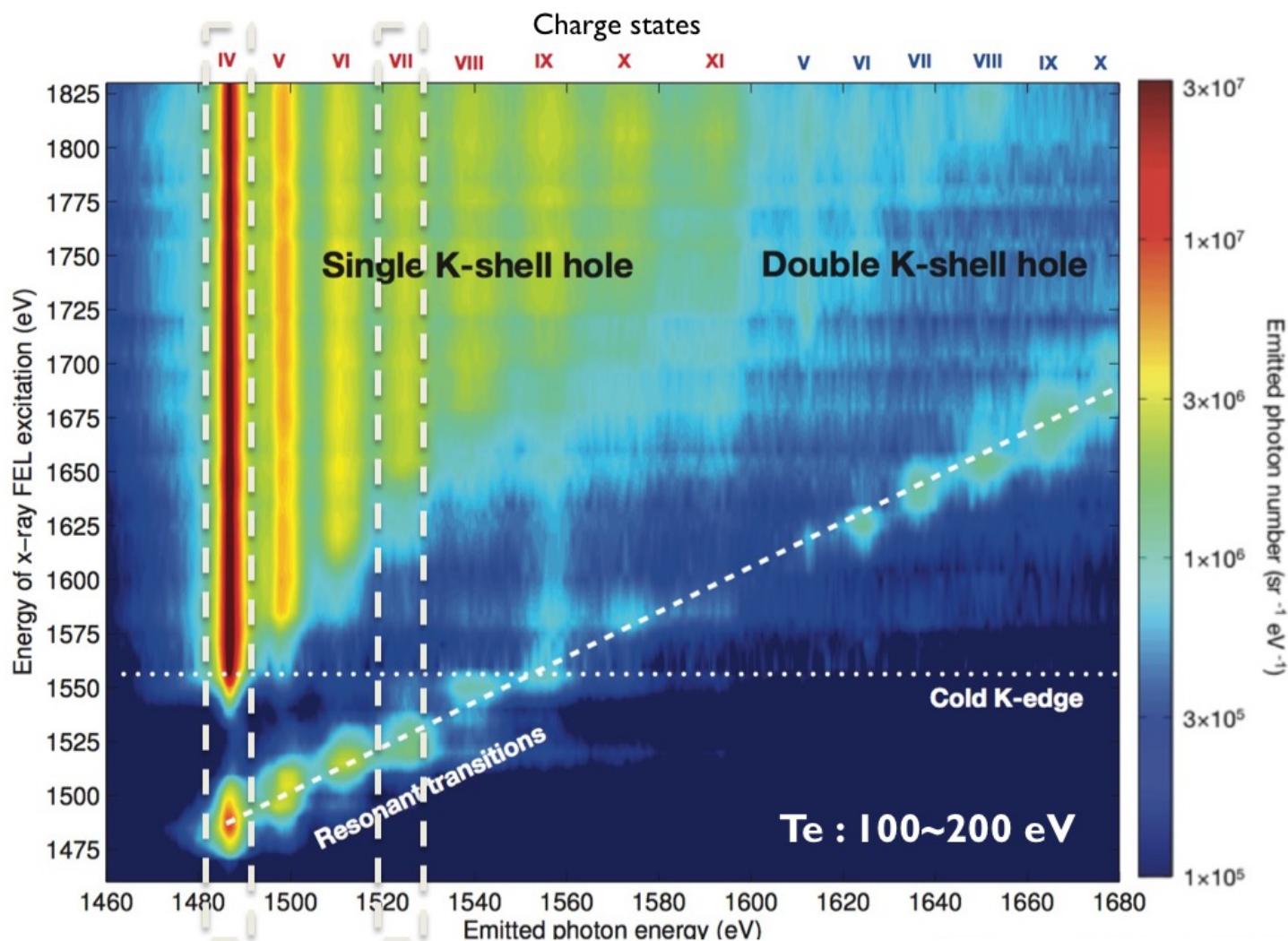
Al foil, 1 μm

Peak intensity
 $\sim 10^{17} \text{ W/cm}^2$



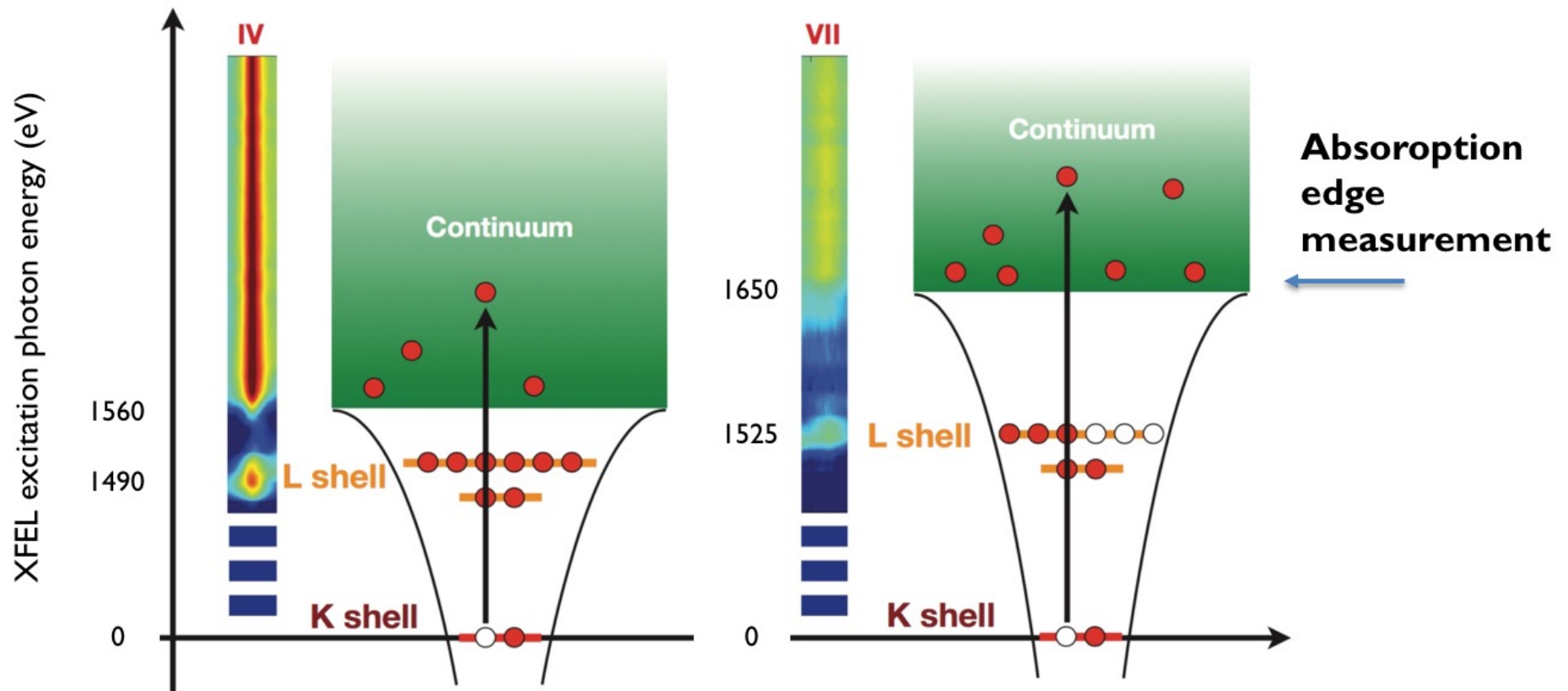
K-shell spectroscopy of hot dense aluminum

XFEL pumping : 10^{12} photons, 10^{17} W/cm²



K-shell emission spectra represent energy levels of ionized aluminum in hot dense conditions

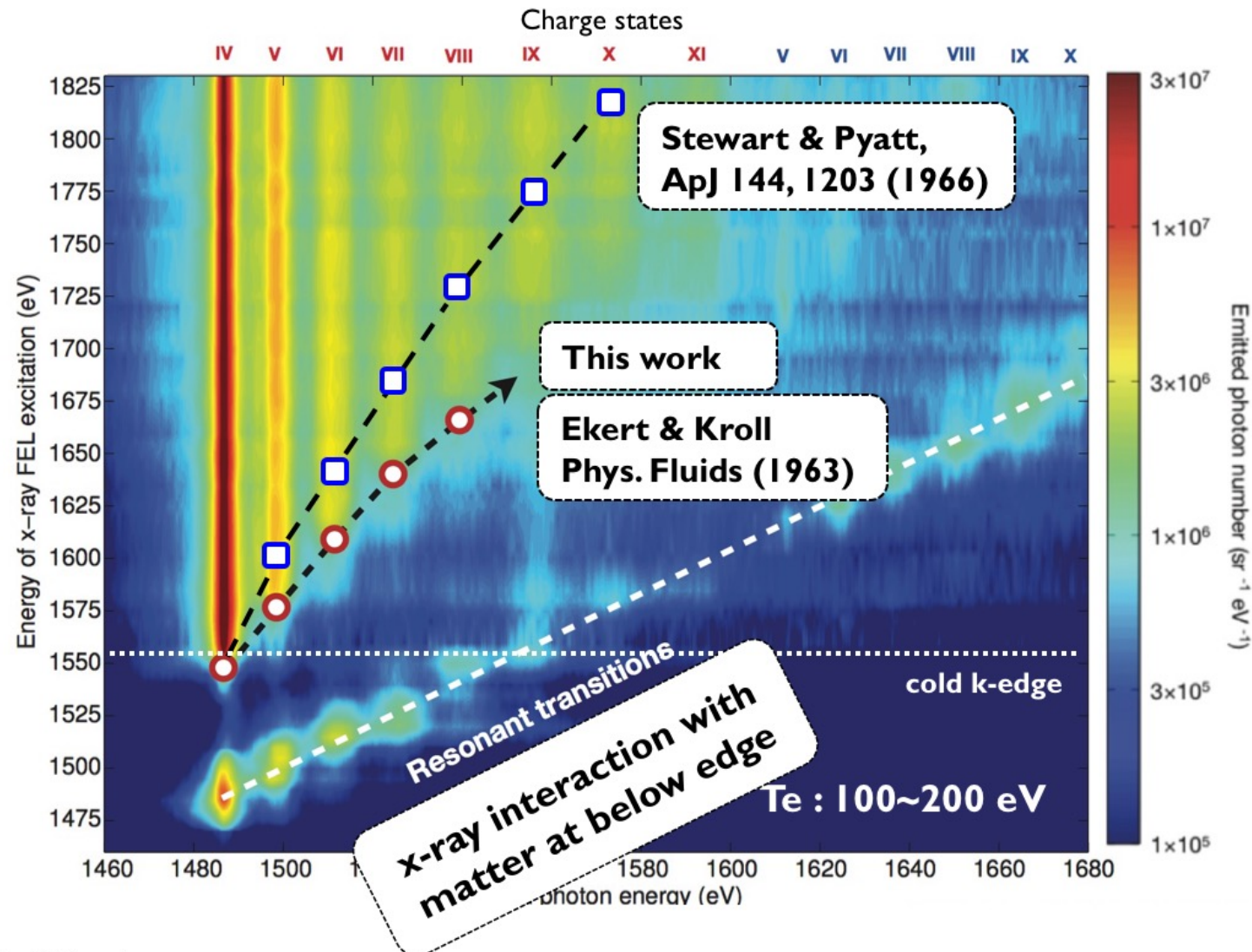
Absorption thresholds is observed



Experimental IPD = Atomic edge - Measured edge

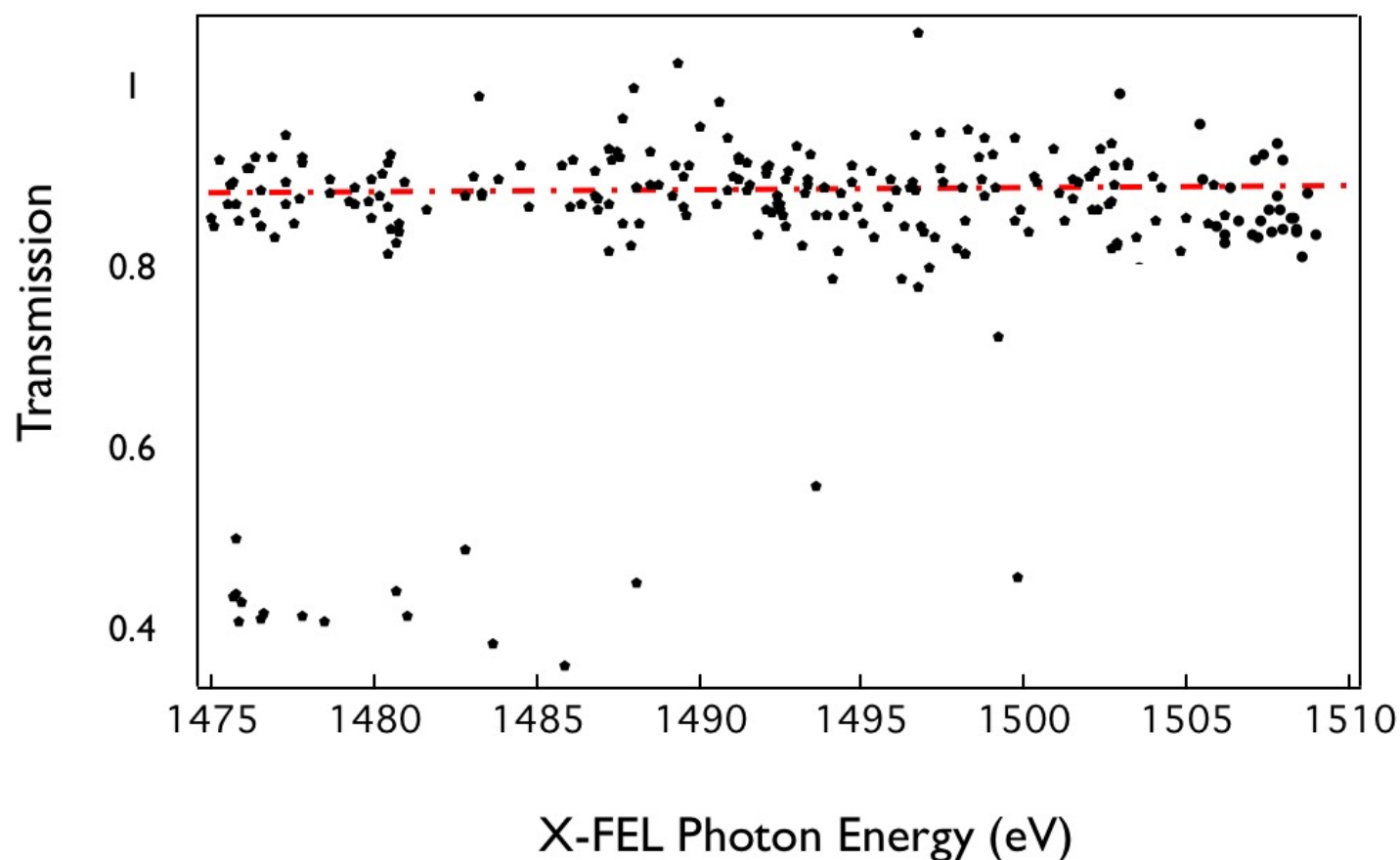
Experimental benchmarks for IPD theories

Ionization Potential Measurement in Hot Dense Matter



X-FEL transmission *below* K-edge: Experiment

Peak intensity $< 10^{16} \text{ W/cm}^2$ Pulse length $< 100 \text{ fs}$
Photon energy : 1470 ~ 1510 eV Aluminum foil : 1.4 μm



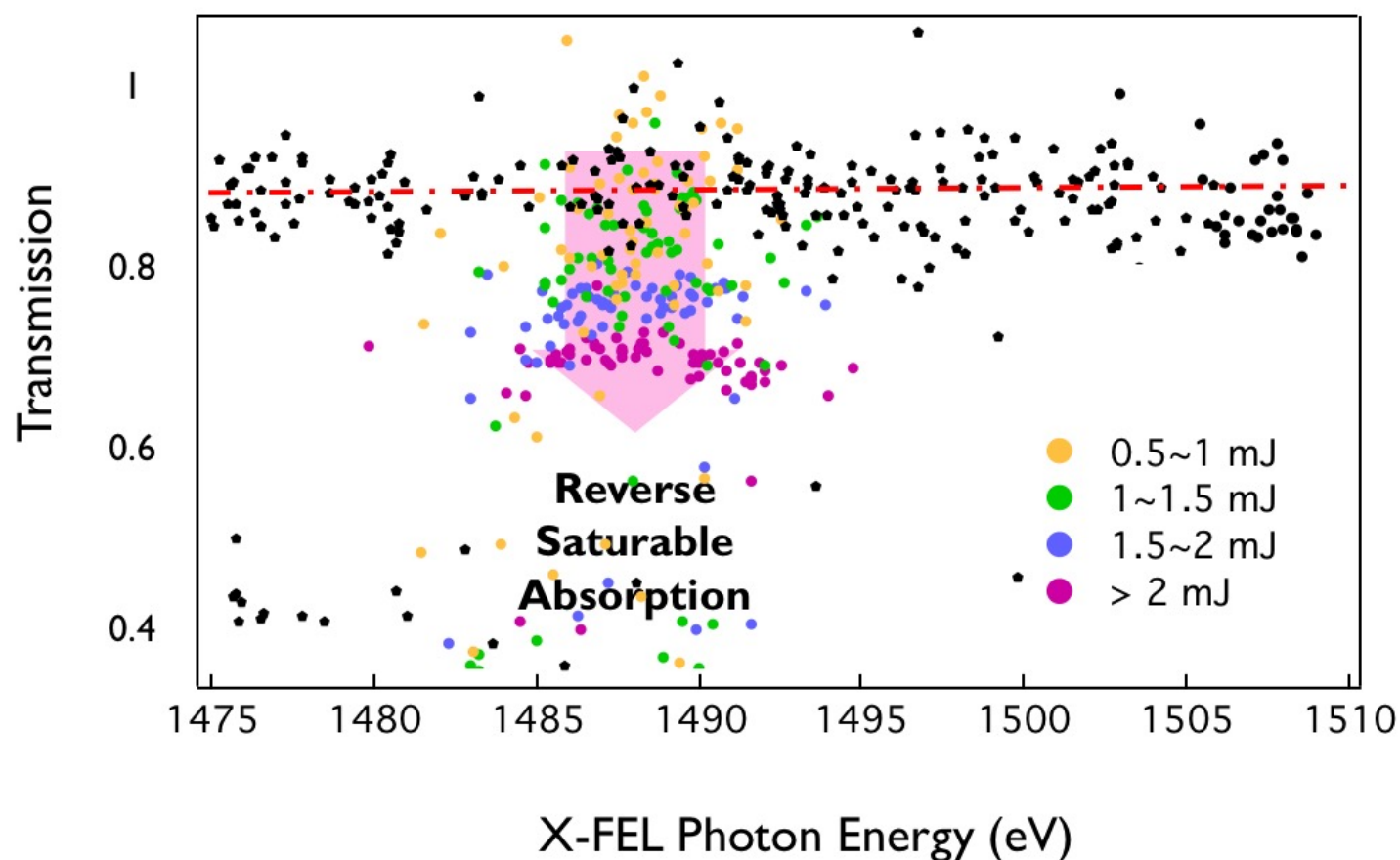
X-FEL transmission *below* K-edge: Experiment

Peak intensity $\sim 10^{16-17}$ W/cm²

Pulse length < 100 fs

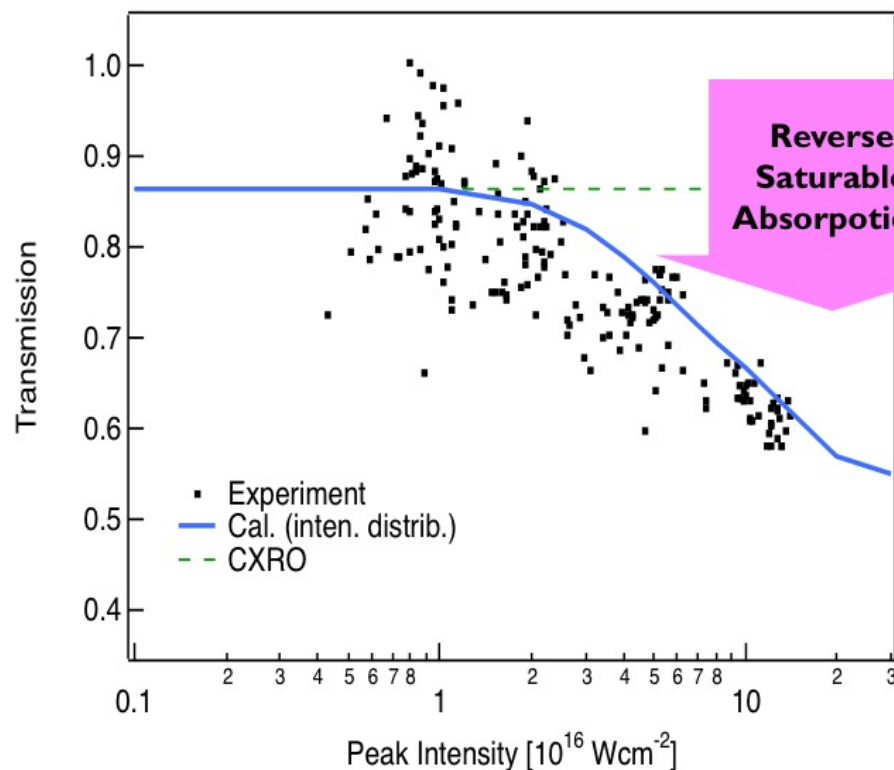
Photon energy : ~ 1487 eV

Aluminum foil : 1.4 μ m



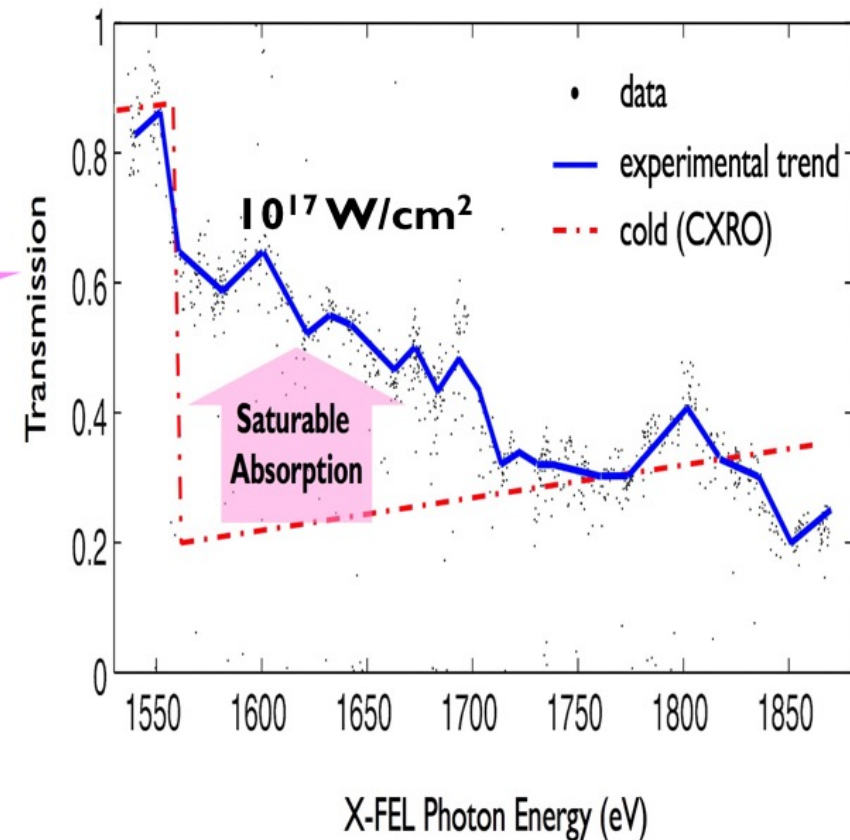
Intensity dependent x-ray transmission of hot dense aluminum

Transmission vs Intensity



Cho et al, - submitted

Transmission vs photon energy

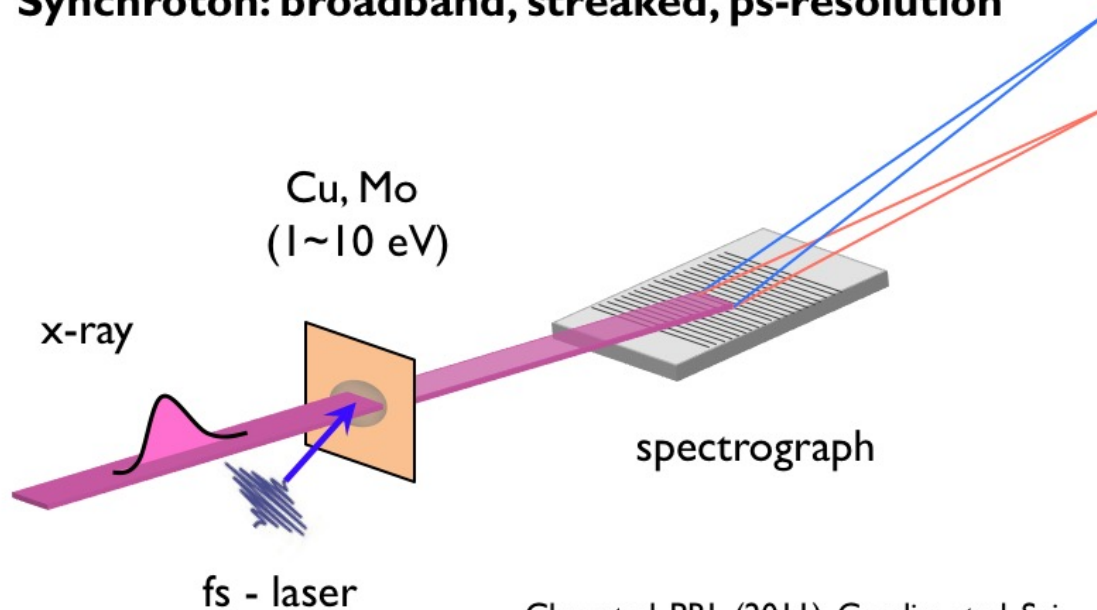


Rackstraw et al, PRL (2015)

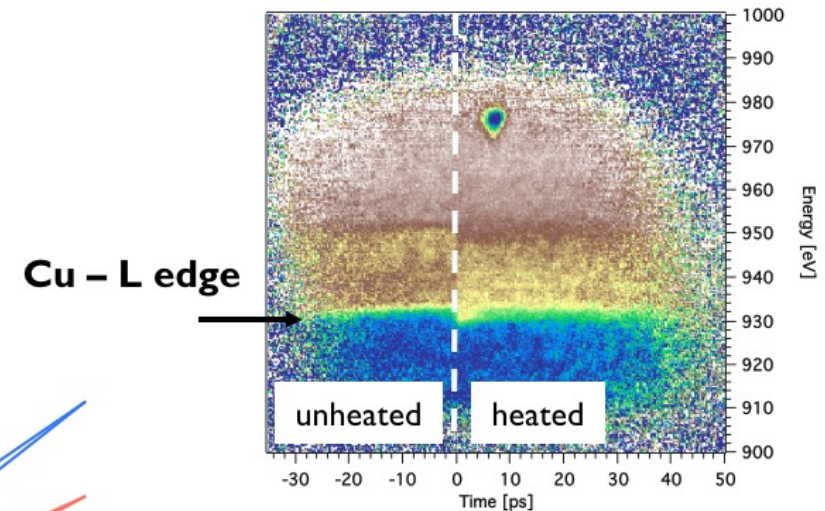
Opacity & energy transport via radiation in HED matters

Time resolved x-ray absorption spectroscopy for super heated metals

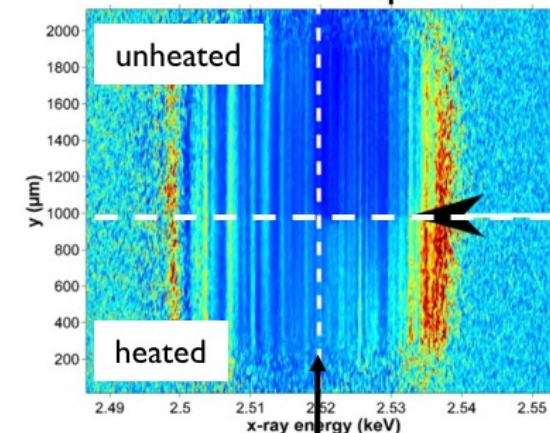
- Femtosecond laser pulse
~ 10 mJ / pulse
heating sample above the damage threshold
Te: 1~10 eV
- X-rays
FEL: SASE, fs-resolution
Synchrotron: broadband, streaked, ps-resolution



streaked spectrum of
70 ps x-ray pulse



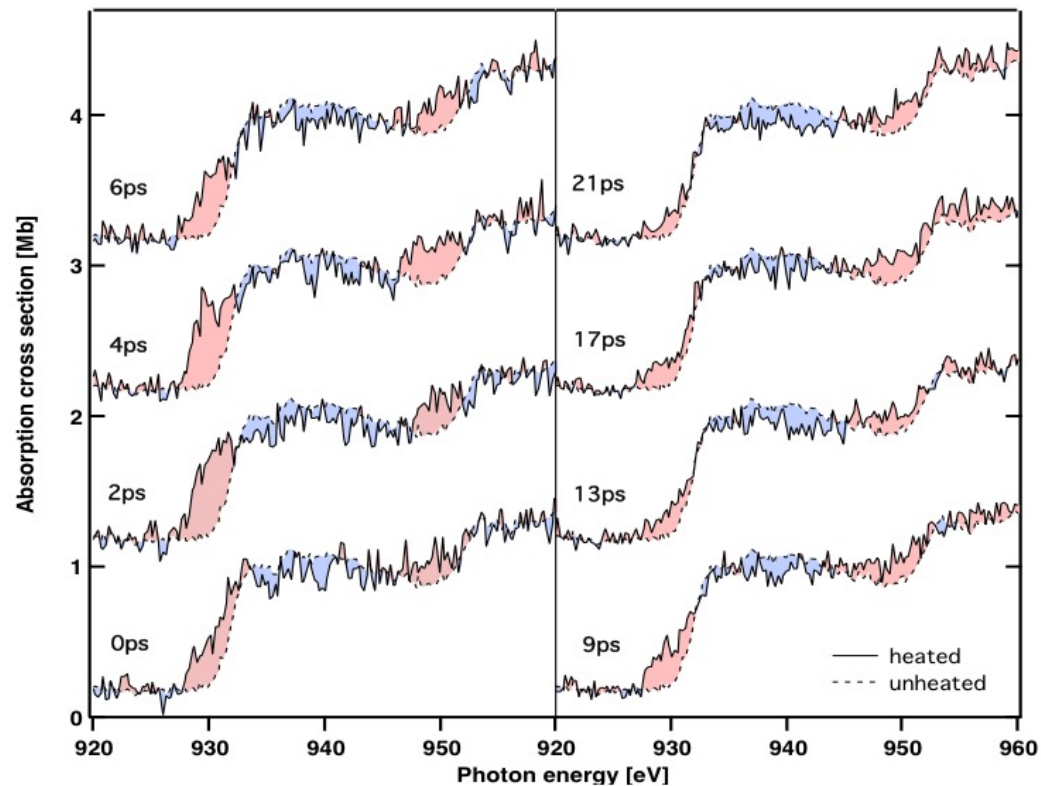
80 fs XFEL spectra



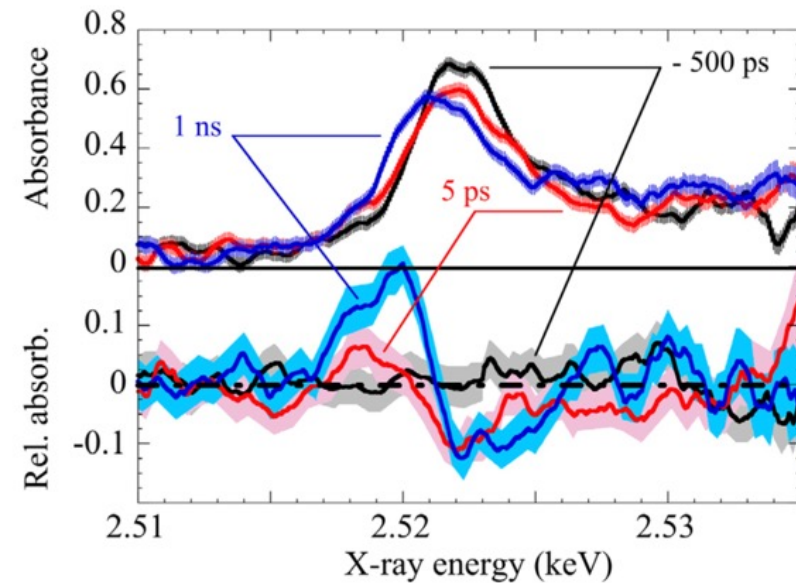
Mo - L edge

Time resolved x-ray absorption spectroscopy

near L-edge of Cu



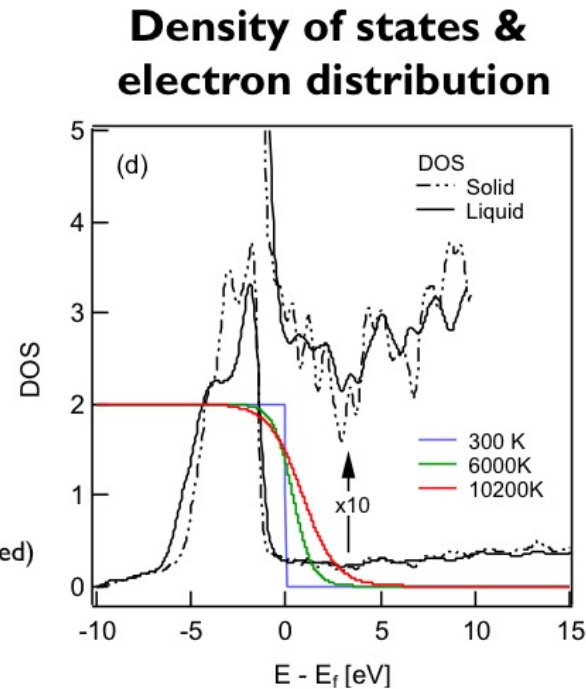
near L-edge of Mo



The first principle DFT-MD simulations to produce electronic DOS of HED matters

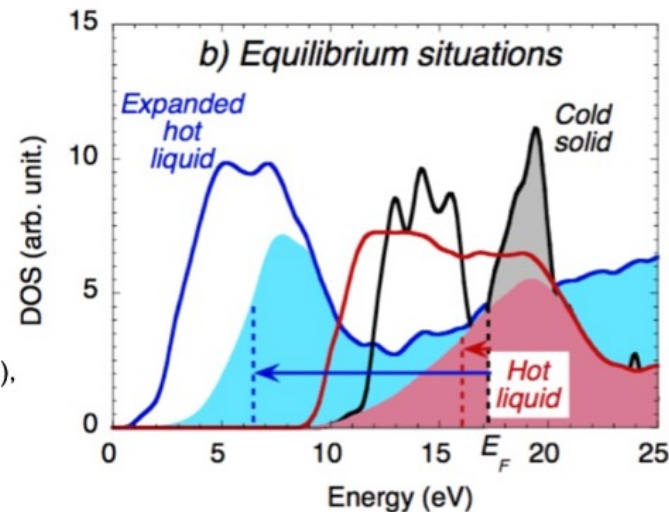
Cu

Cho, et al, PRL (2011),
Cho et al, Sci Rep (to be published)

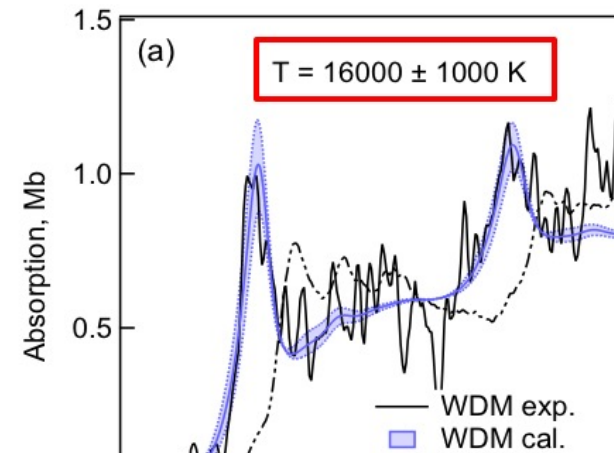


Mo

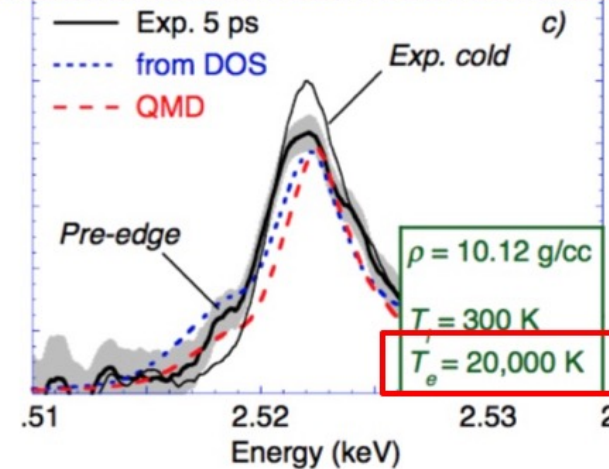
Gaudin, et al, Sci Rep. (2014),
Dorchies et al, PRB (2015)



XANES

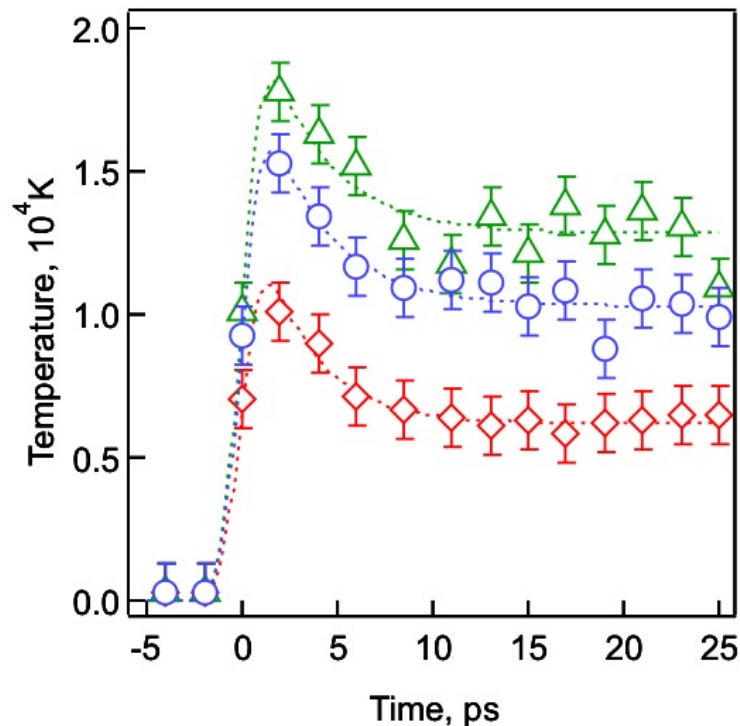


XANES serves as an electron temperature diagnostics of HED matter

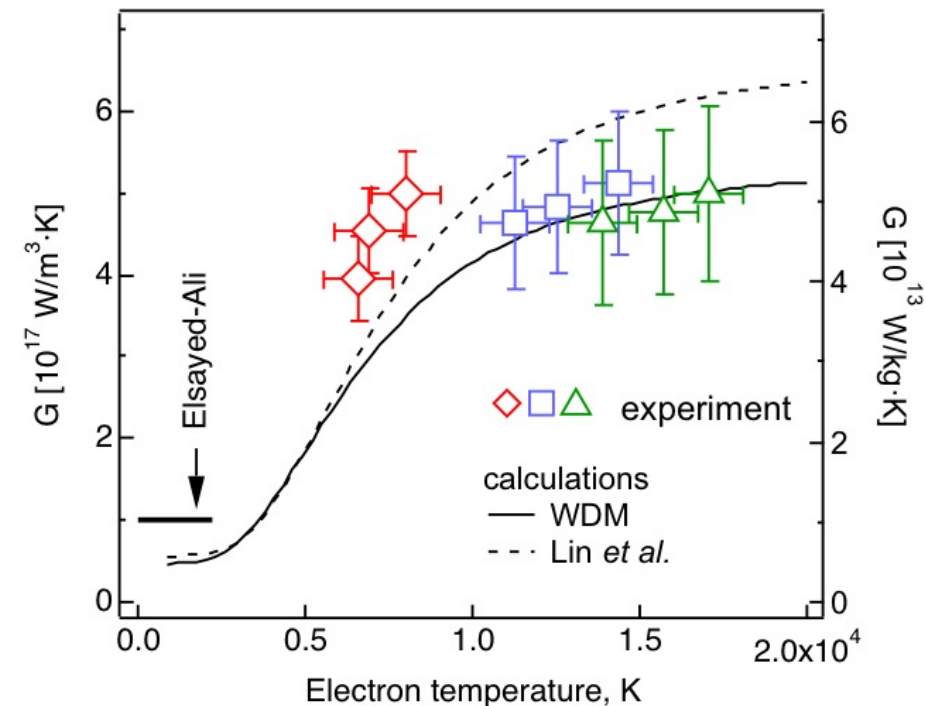


Electron - ion energy relaxation (electron-phonon coupling) of HED matters are tested

Electron temperature



Temperature dependent el-ph couplings



TTM eq.

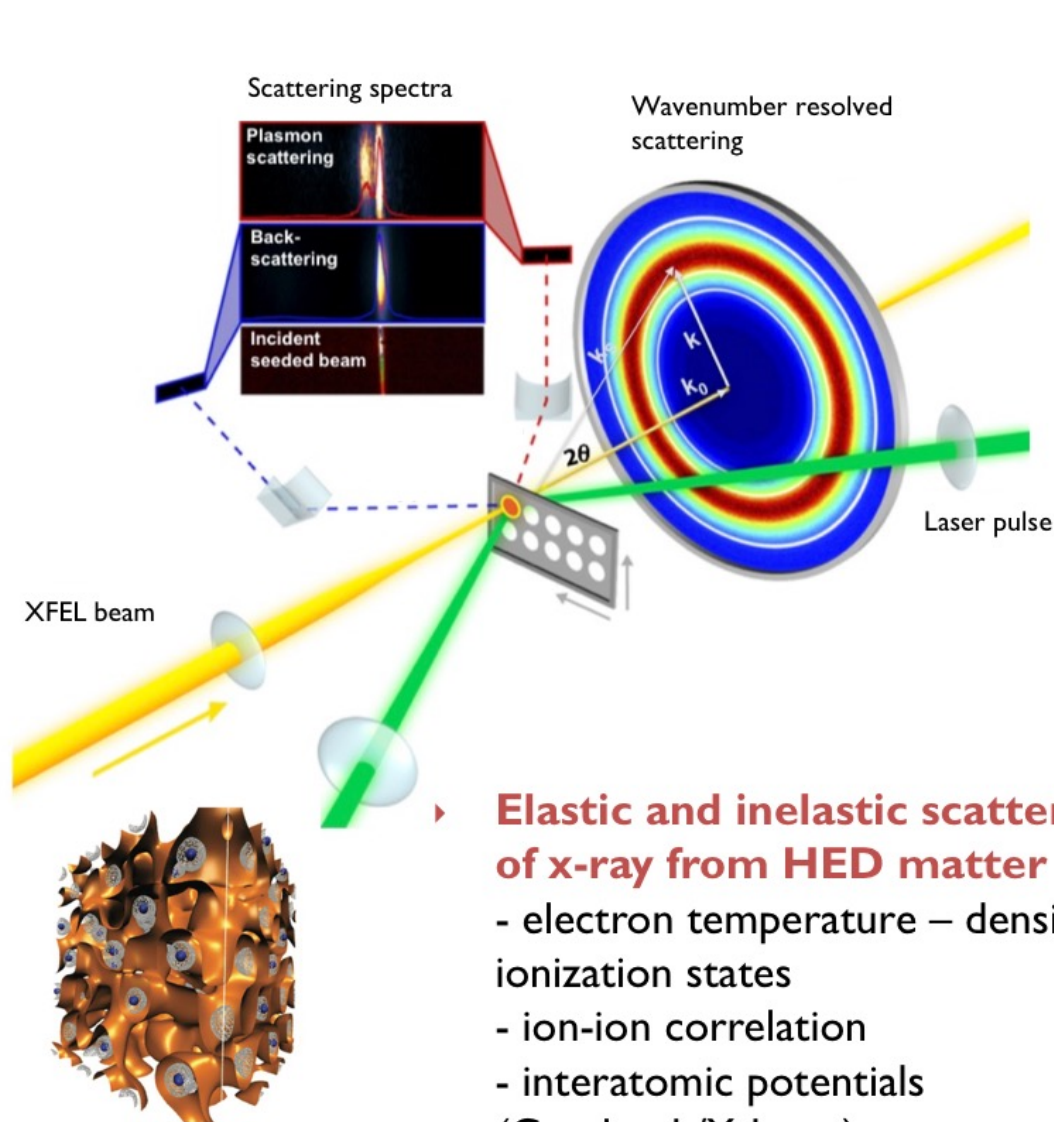
$$C_e(T_e)T_e'(t) = -G [T_e(t) - T_i(t)] + S_L(t)$$

$$C_i T_i'(t) = G [T_e(t) - T_i(t)]$$

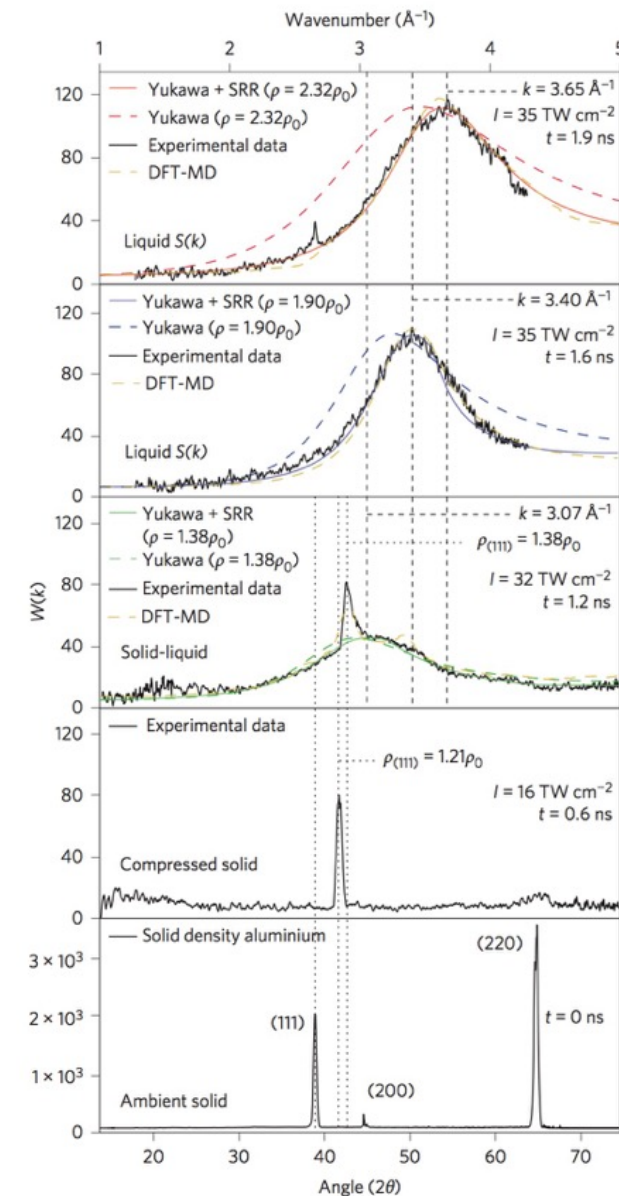
el-ph coupling

$$G(T_e) = \frac{\pi \hbar k_B \lambda \langle \omega^2 \rangle}{g(\epsilon_F)} \int_{-\infty}^{\infty} g^2(\epsilon) \left(-\frac{\partial f}{\partial \epsilon} \right) d\epsilon$$

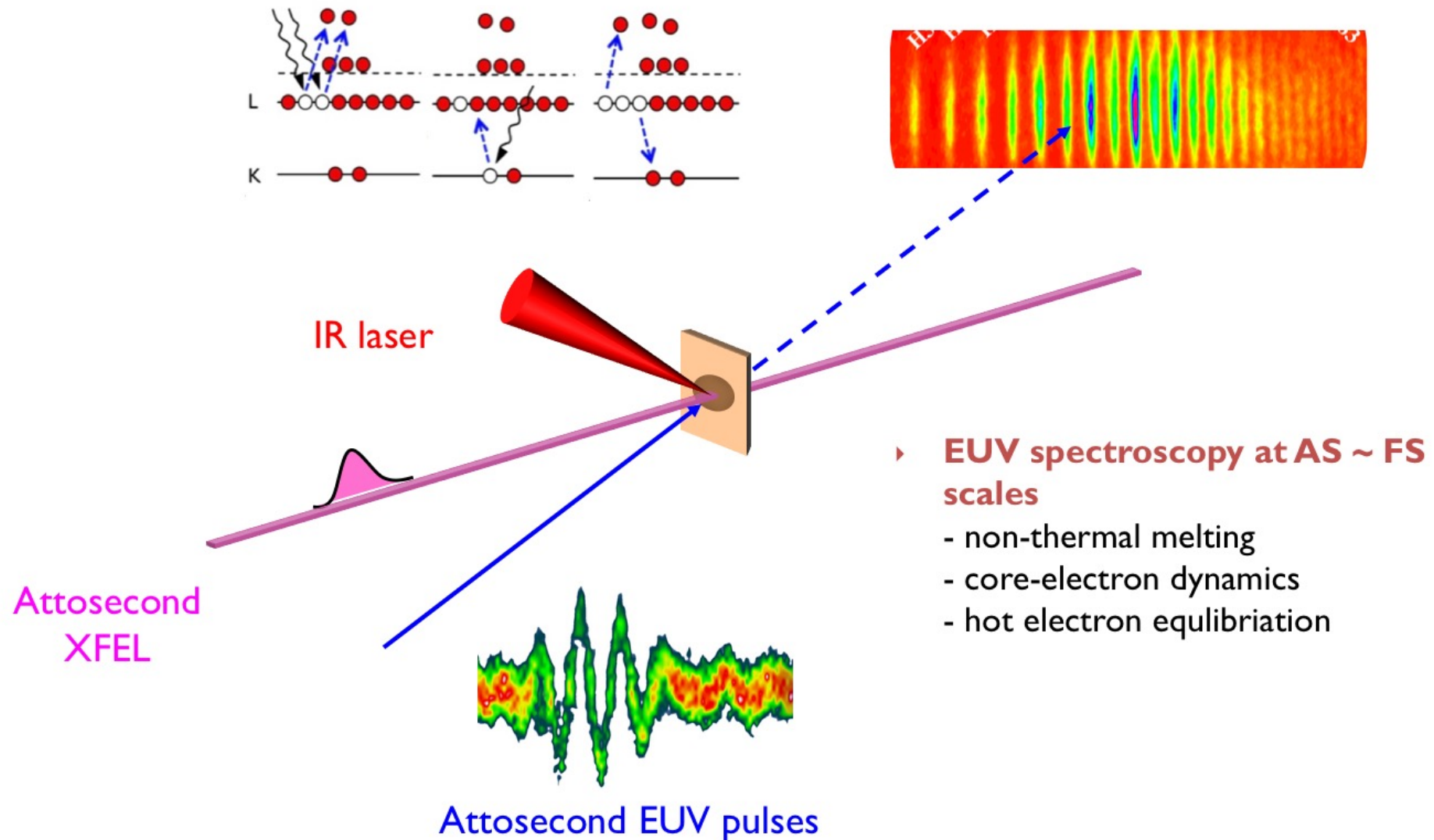
X-ray scattering : lattice compression and melting in HED conditions



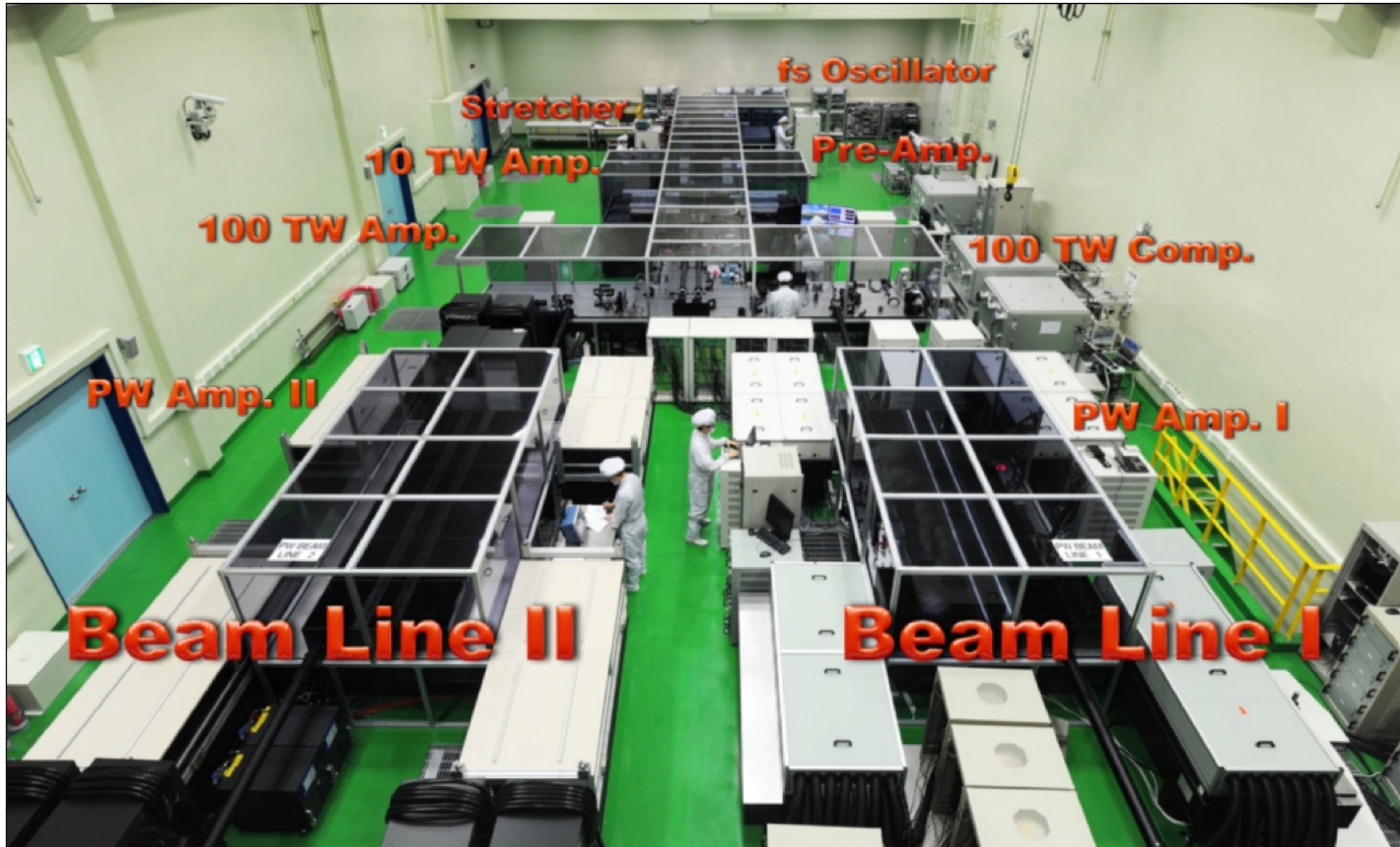
- Elastic and inelastic scattering of x-ray from HED matter
 - electron temperature – density & ionization states
 - ion-ion correlation
 - interatomic potentials (Coulomb/Yukawa)



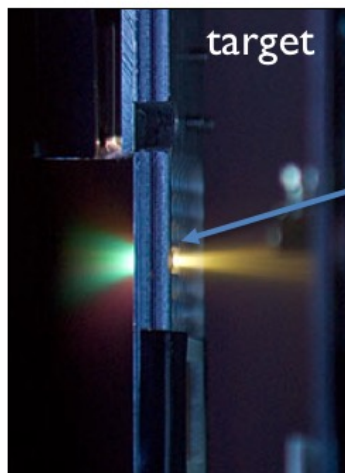
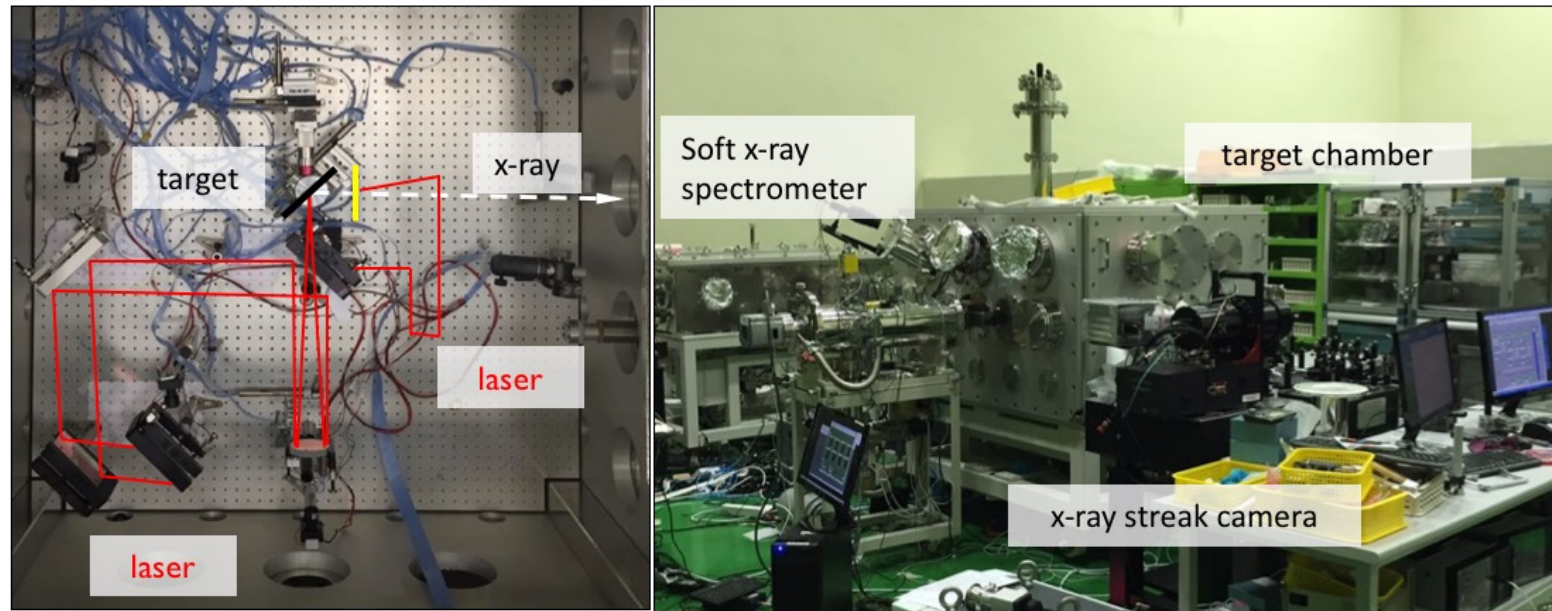
Attosecond EUV pulses: Core electron dynamics in HED conditions



HED researches at GIST PW laser facilities

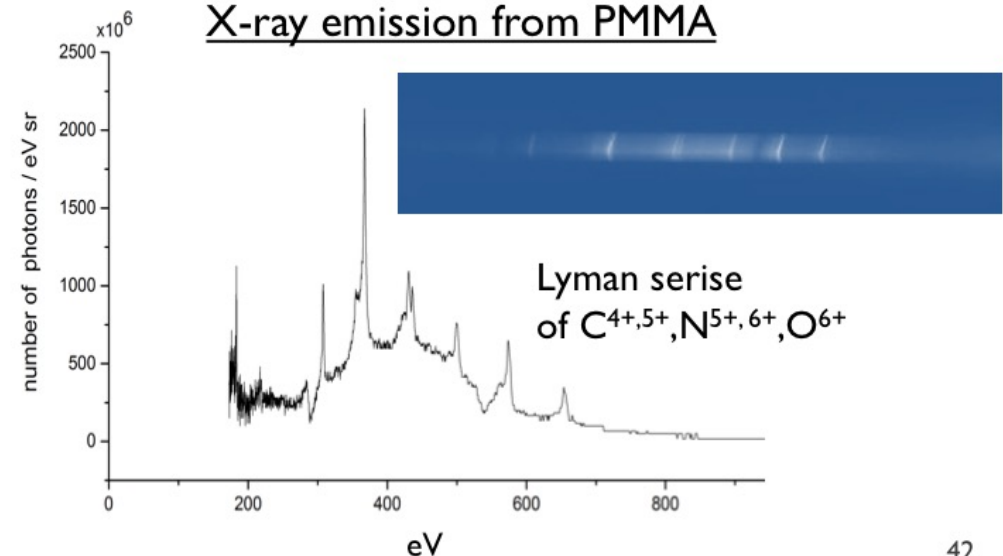


X-ray emission & absorption spectroscopy for laser plasmas

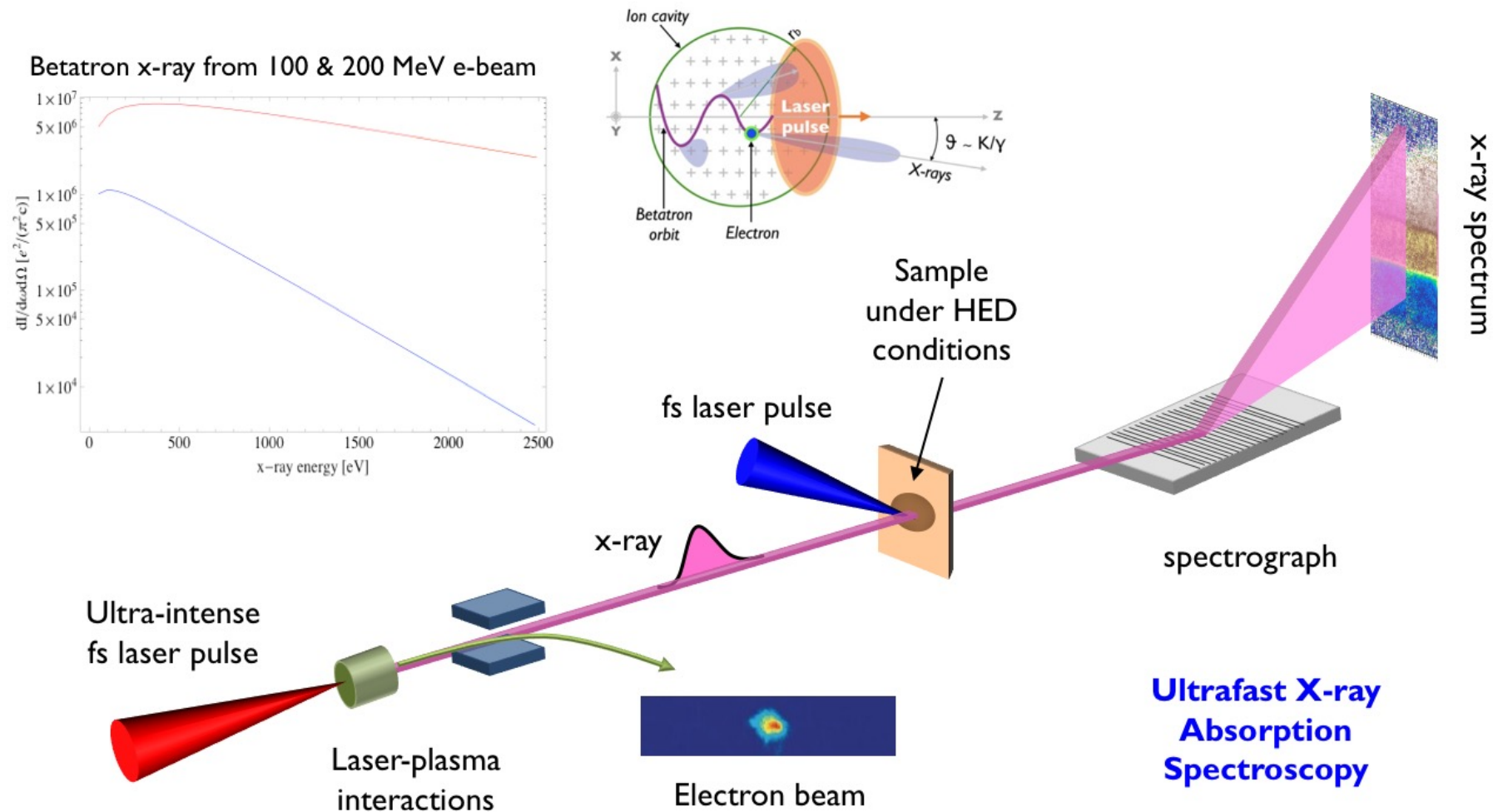


$T_e \sim 10 \text{ keV}$
 $N_e \sim 10^{20} \text{ cm}^{-3}$

X-ray emission from PMMA



Betatron radiation from the wakefield accelerated electrons is broadband – femtosecond x-ray source



Ultrafast structural changes - non-thermal melting
Ultrafast electron dynamics – hot electron thermalization

- High Energy Density Physics is the study of matter at extremely high pressures and temperatures. The challenge of HEDP centers on difficulties in describing and measurement of the basic properties.
- With recent development of extreme light sources, which delivers large amount of energy in very short time scale (XFEL and PW lasers), qualitative physics experiment became feasible.
- Some fundamental properties of a few HED conditions (temperature, ionization potential, transmission, energy relaxation, etc) were experimentally measured recently, starting to provide benchmark data for theories.
- HEDP is closely related to plasma physics, condensed matter physics, optics, x-ray science, astrophysics, etc. We want to bring new experimental techniques and physical understanding from various disciplines and apply them for this new science.

High Energy Density Physics group at GIST

