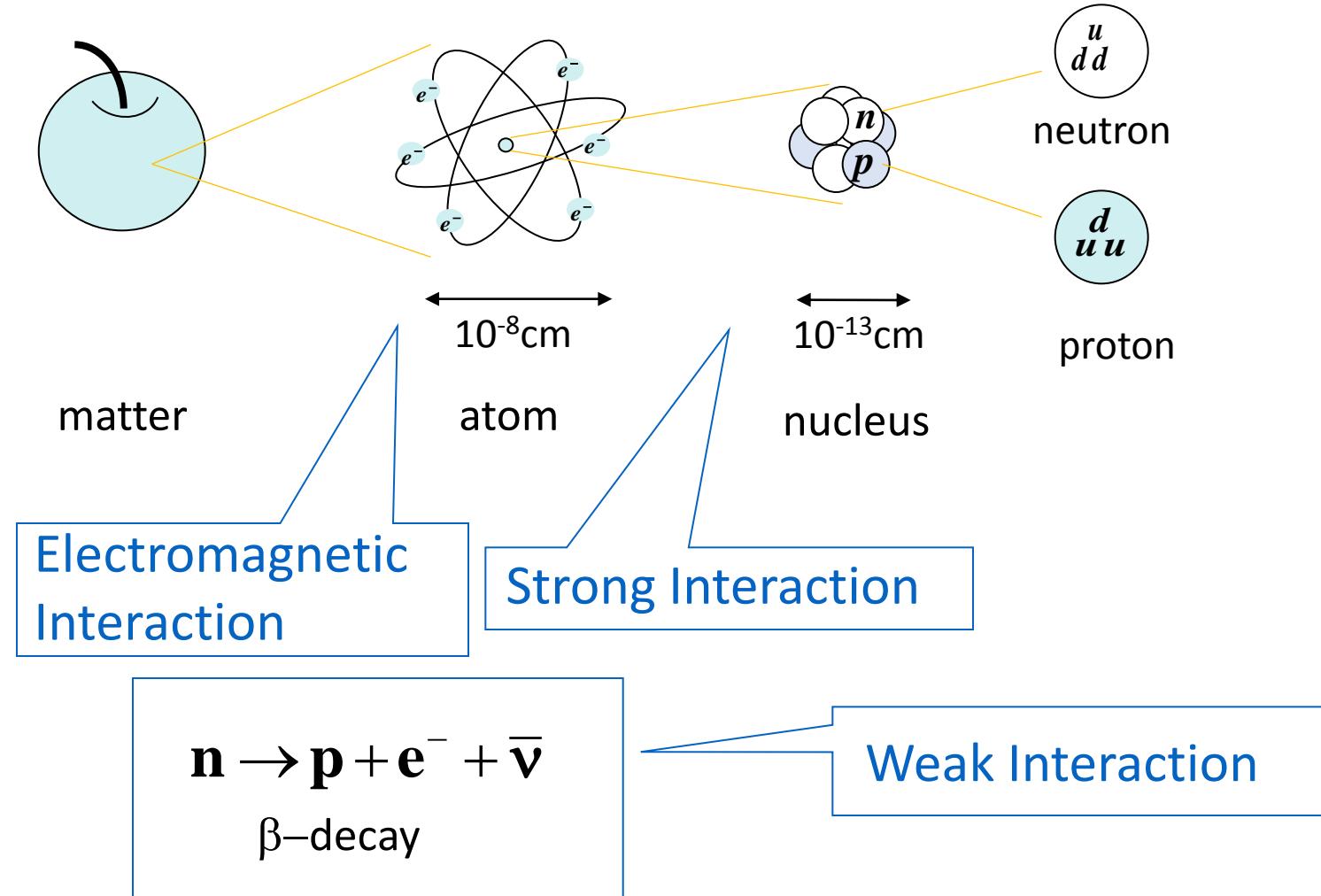


2015/10/08
Seoul Univ.

Matter and Antimatter

KEK and JSPS
M. Kobayashi

Structure of matter



Discovery of Antiparticle

Special Relativity (1905)

Quantum Mechanics (1925)

Schrodinger equation

$$i\hbar \frac{\partial}{\partial t} \psi(\mathbf{x}) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{x}^2} \psi(\mathbf{x}) \quad \longleftrightarrow \quad E = \frac{p^2}{2m}$$



Dirac equation (Dirac 1928)

$$i\hbar \frac{\partial}{\partial t} \psi(\mathbf{x}) = \left(\frac{\hbar \mathbf{c}}{\mathbf{i}} \vec{\alpha} \cdot \frac{\partial}{\partial \mathbf{x}} + \beta m c^2 \right) \psi(\mathbf{x})$$

$\vec{\alpha}, \beta$: 4×4 matrix

$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix}$$

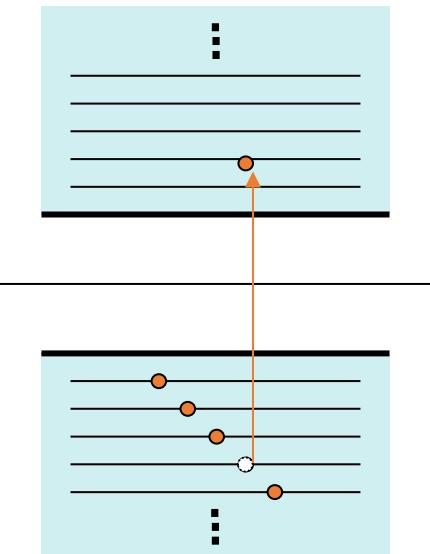
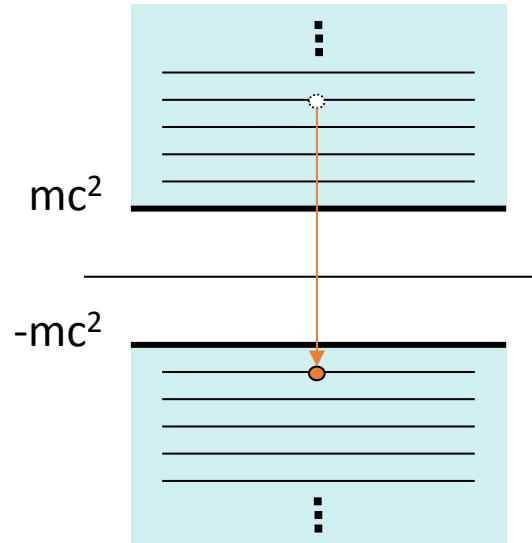
Solutions of the Dirac equation

Positive energy

$$E = \sqrt{p^2 c^2 + m^2 c^4} \geq mc^2$$

Negative energy

$$E = -\sqrt{p^2 c^2 + m^2 c^4} \leq -mc^2$$



Dirac Sea

All negative energy states are occupied

Hole in the Dirac Sea

→ Positively charged particle

1932 Anderson

Discovery of Positron

Dirac's arguments: only for fermions

How about bosons ?



Relativistic Quantum Field Theory

Every particle has its corresponding antiparticle

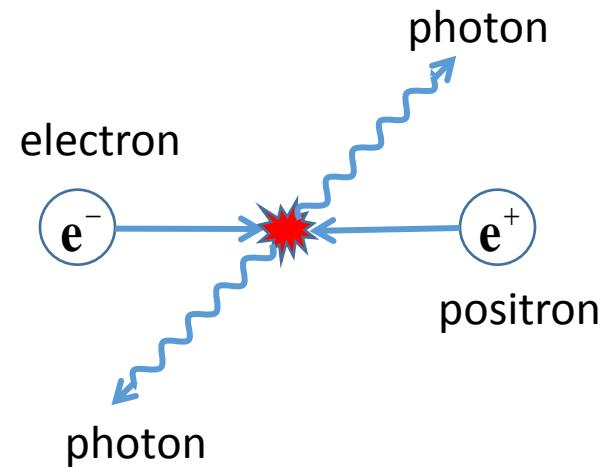
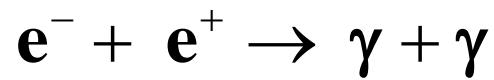
- same mass
- opposite charge

Note

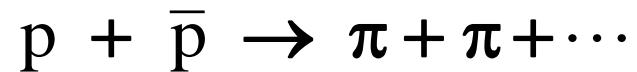
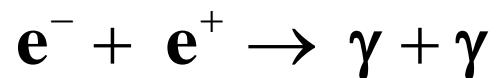
- Self conjugate particle: a particle and an antiparticle are identical
Example: **Photon**
- Neutral particles are not necessarily self conjugate
Example: **Neutron**

electron $e^- \Leftrightarrow$ positron e^+
proton $p \Leftrightarrow$ antiproton \bar{p}
neutron $n \Leftrightarrow$ antineutron \bar{n}

Pair annihilation



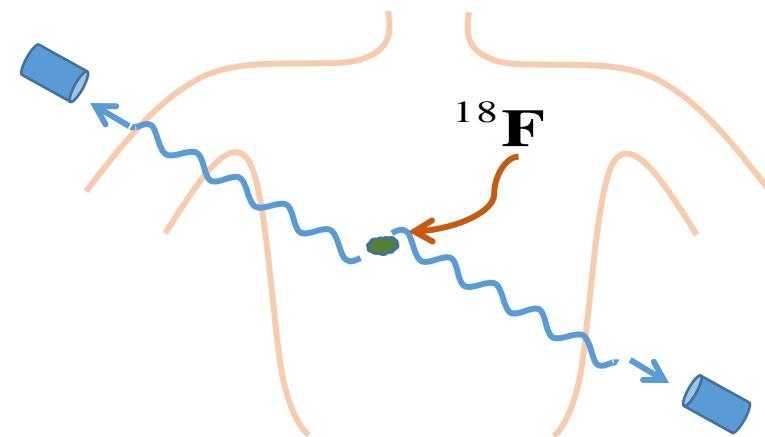
Pair annihilation



Pair creation



PET



CP symmetry

Invariance of the laws of nature under the exchange
of particle and antiparticle

C (charge conjugation) : simple exchange of particle and antiparticle

P (parity) : space inversion

CP : combined operation of C and P

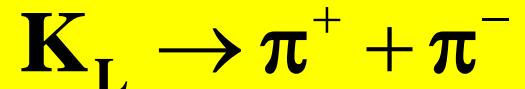
P violation was found in 1956.

→ C is also violated, but CP looked to hold true

Discovery of CP violation

1964 Cronin, Fitch et al

~0.2% of K_L decay into



$$|K_L\rangle = p |K^0\rangle + q |\bar{K}^0\rangle$$

$|p| \neq |q| \Rightarrow \text{CP violation}$

Hadron : strongly interaction particles

proton p neutron n

pion π^+ π^- π^0

kaon K^+ K^- $K^0 \bar{K}^0$

lambda Λ sigma $\Sigma^+ \Sigma^- \Sigma^0$

1956 Sakata [Sakata Model](#)
All the hadrons are composite states of

p, n, Λ

: Fundamental Triplet

$p\bar{n}$

π^+

$p\bar{\Lambda}$

K^+

$\Lambda p\bar{n}$

Σ^+

$\Lambda\Lambda\bar{n}$

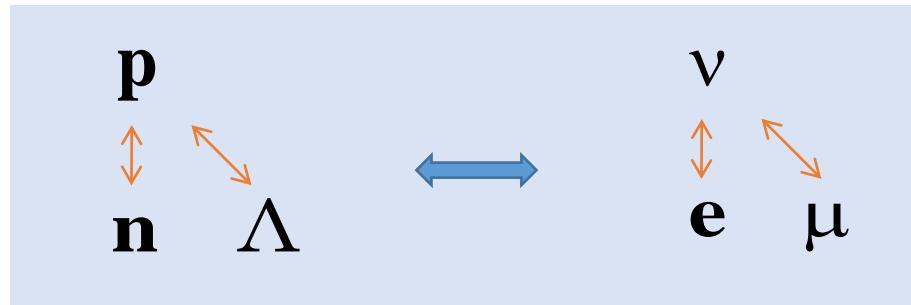
Ξ^0



Courtesy of Sakata Memorial
Archival Library

Shoichi Sakata
1911-1970

Weak Interaction in the Sakata Model



1959 Gamba, Marshak, Okubo
B-L Symmetry

1960 Maki, Nakagawa, Ohnuki, Sakata

Nagoya Model : $\mathbf{p} = \langle \mathbf{B}^+ v \rangle$, $\mathbf{n} = \langle \mathbf{B}^+ e \rangle$, $\Lambda = \langle \mathbf{B}^+ \mu \rangle$

\mathbf{B}^+ : B-matter

Discovery of Two Neutrinos

1962 Danby et al. PRL 9, 36 (1962)

$$\begin{array}{ccc} \nu_e & & \nu_\mu \\ \uparrow & & \downarrow \\ e & & \mu \end{array}$$

Modification of Nagoya Model

1962 Katayama, Matumoto, Tanaka and Yamada
Progr. Theor. Phys. 28, 675 (1962)

1962 Maki, Nakagawa and Sakata
Progr. Theor. Phys. 28, 870 (1962)

$$\mathbf{p} = \langle \mathbf{B}^+ v_1 \rangle, \quad \mathbf{n} = \langle \mathbf{B}^+ e \rangle, \quad \Lambda = \langle \mathbf{B}^+ \mu \rangle, \quad \mathbf{p}' = \langle \mathbf{B}^+ v_2 \rangle$$

$$v_1 = \cos \delta v_e + \sin \delta v_\mu$$

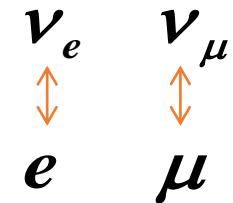
$$v_2 = -\sin \delta v_e + \cos \delta v_\mu$$

Neutrino Mixing: Maki, Nakagawa, Sakata (MNS)

$$\mathbf{p} = \langle \mathbf{B}^+ v_1 \rangle, \quad \mathbf{n} = \langle \mathbf{B}^+ e \rangle, \quad \Lambda = \langle \mathbf{B}^+ \mu \rangle, \quad \mathbf{p}' = \langle \mathbf{B}^+ v_2 \rangle$$

$$v_1 = \cos \delta v_e + \sin \delta v_\mu$$

$$v_2 = -\sin \delta v_e + \cos \delta v_\mu$$



Identify v_1, v_2 as the mass eigenstates

They formulated lepton flavor mixing precisely.

Bare field: $\psi_0 = \begin{pmatrix} \mu_0 \\ e_0 \end{pmatrix}, \quad \phi_0 = \begin{pmatrix} v_{\mu 0} \\ v_{e 0} \end{pmatrix}$

$$\mathbf{L}_{\text{int}} = [(\bar{\psi}_0 \Lambda \psi_0) + (\bar{\phi}_0 \Lambda' \phi_0)] \mathbf{X}^* \mathbf{X}$$

Λ, Λ' : 2x2 mass matrices

Difference of mass diagonalization

→ Mixing angle



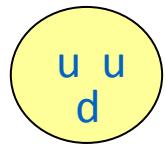
Courtesy of Sakata Memorial Archival Library

Quark Model

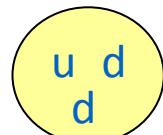
1964 Gell-Mann

quark u d s

2/3e -1/3e -1/3e



proton



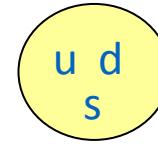
neutron



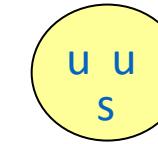
π^+



K^+



Λ



Σ^+

1963 Cabibbo

Cabibbo angle θ_C

$$u \leftrightarrow d' = \cos \theta_C \ d + \sin \theta_C \ s$$

Status at the end of 1960's

Fundamental particles

quarks	u
	d s
leptons	ν_e ν_μ
	e μ

Fundamental interactions

Electromagnetic	Renormalization theory Tomonaga-Schwinger-Feynman
Weak	No satisfactory theory
Strong	

1971 't Hooft : Renormalization of Non-Abelian gauge theory

- • Weinberg-Salam-Glashow theory for Electro-Weak Int.
• QCD for Strong Int.

1973 Kobayashi, Maskawa

How to accommodate CP violation

What we found

- Not possible with three or four quarks
- Existence of unknown particles
- A possible candidate is six-quark model

u c t
d s b

Flavor Mixing of 4-Quark Model

$$\begin{pmatrix} u \\ d' \end{pmatrix}_L \quad \begin{pmatrix} c \\ s' \end{pmatrix}_L$$

$$\begin{pmatrix} d' \\ s' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} \\ V_{cd} & V_{cs} \end{pmatrix} \begin{pmatrix} d \\ s \end{pmatrix}$$

Redefinition of quark fields

$$u \rightarrow e^{i\delta_u} u, \quad d \rightarrow e^{i\delta_d} d, \dots$$

$$\begin{pmatrix} e^{-i\delta_u} & 0 \\ 0 & e^{-i\delta_c} \end{pmatrix} \begin{pmatrix} V_{ud} & V_{us} \\ V_{cd} & V_{cs} \end{pmatrix} \begin{pmatrix} e^{i\delta_d} & 0 \\ 0 & e^{i\delta_s} \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix}$$



No CP Violation

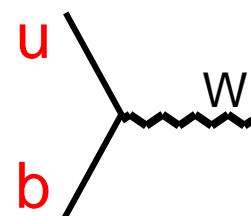
6-Quark Model for CP Violation

$$\begin{pmatrix} \mathbf{u} \\ d' \end{pmatrix} \quad \begin{pmatrix} \mathbf{c} \\ s' \end{pmatrix} \quad \begin{pmatrix} \mathbf{t} \\ b' \end{pmatrix} \quad \begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \approx \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

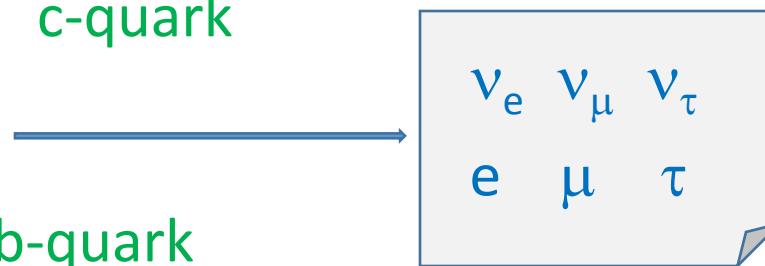


Complex Coupling Constant



Development of Experiments

- 1974 Discovery of J/ ψ particle c-quark
- 1975 Discovery of τ -lepton
- 1977 Discovery of Y particle b-quark
- 1995 Discovery of t-quark



Theory of Strong Interactions

- 1973 QCD Asymptotic freedom
- Quark confinement



Standard Model

Status at the end of 1960's

Fundamental particles

	u	c	t
quarks	d	s	b
	ν_e	ν_μ	ν_τ
leptons	e	μ	τ

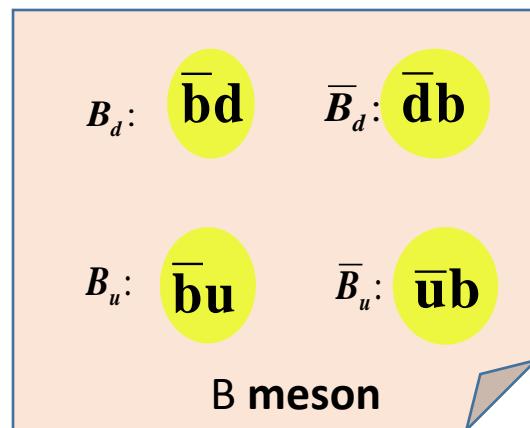
Fundamental interactions

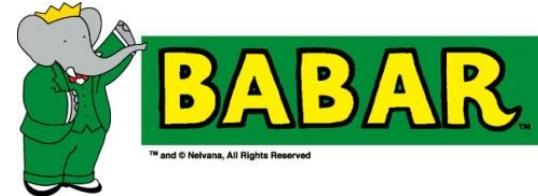
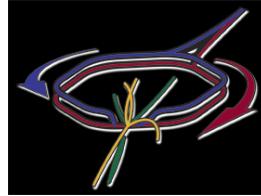
Electromagnetic	}	Renormalization theory
Weak		Tammega-Schwinger-Feynman theory
Strong		No satisfactory theory
		QCD

CP violation experiments at B-factories

B-factories at KEK and SLAC

- e^+e^- collider
- optimized for B-meson production
- asymmetric energy





KEKB/Belle (Japan)

$E(e^-)=8\text{GeV}$,
 $E(e^+)=3.5\text{GeV}$

Finite angle beam crossing

Feature

$E(e^-)=9\text{GeV}$,
 $E(e^+)=3.1\text{GeV}$

Zero angle beam crossing

1994

Governmental
Approval

1993

Experiment

May 1999 –
Jun. 2010

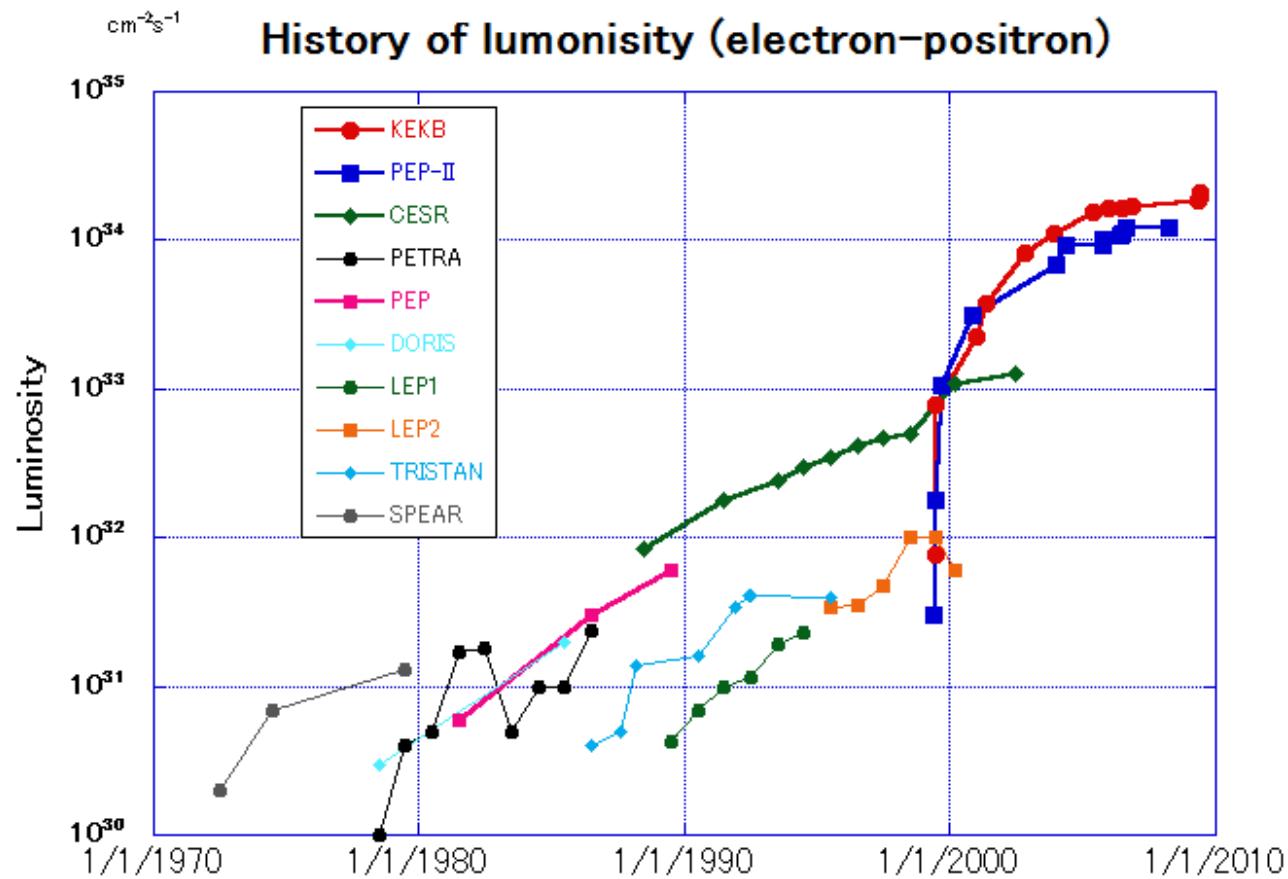
May 1999 –
Apr. 2008

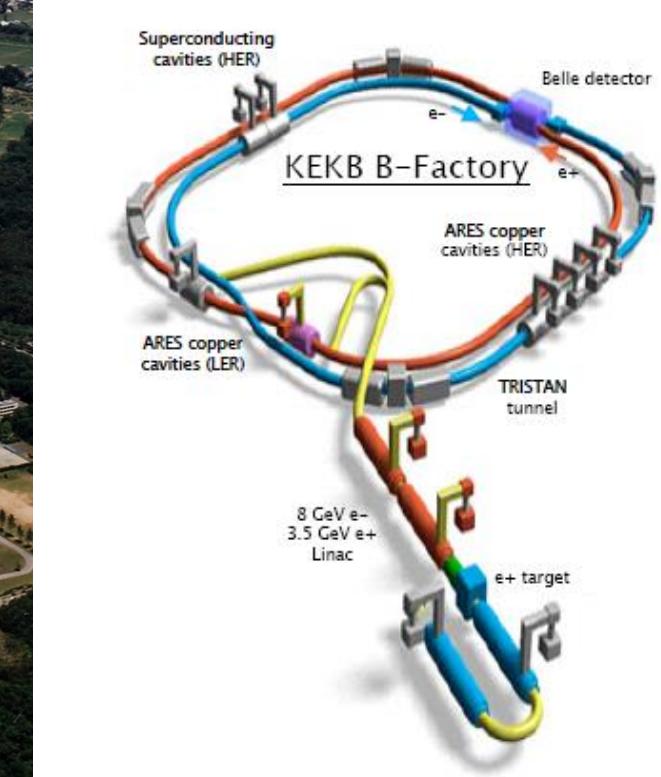


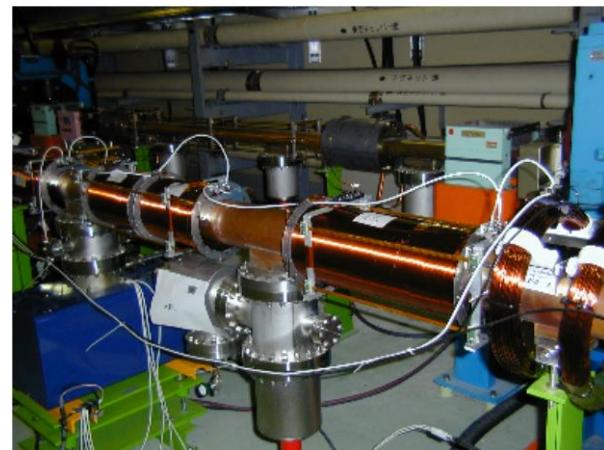
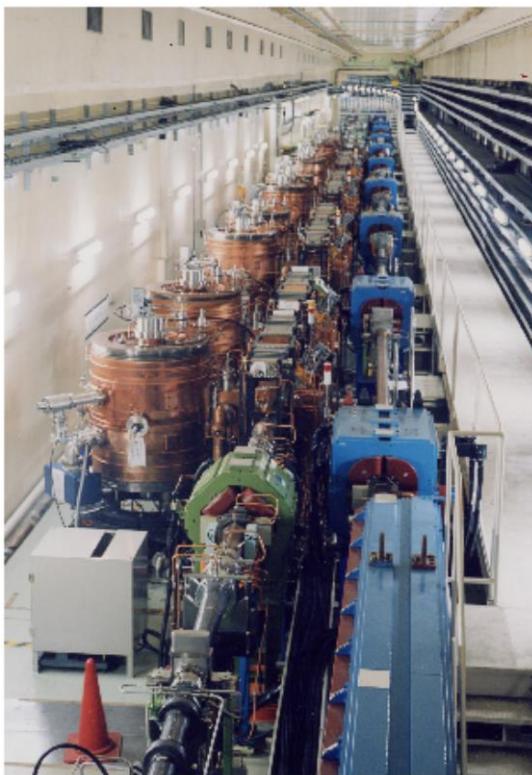
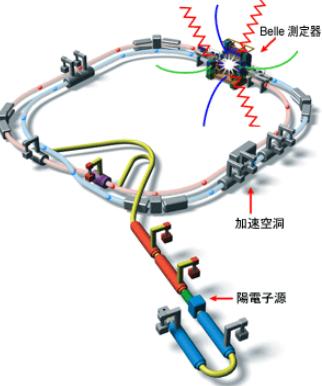
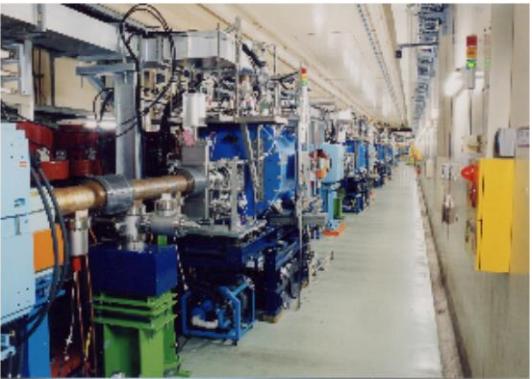
Super KEKB

$$\text{KEKB } L = 2.1 \times 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$$

$$\text{PEP-II } L = 1.2 \times 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$$

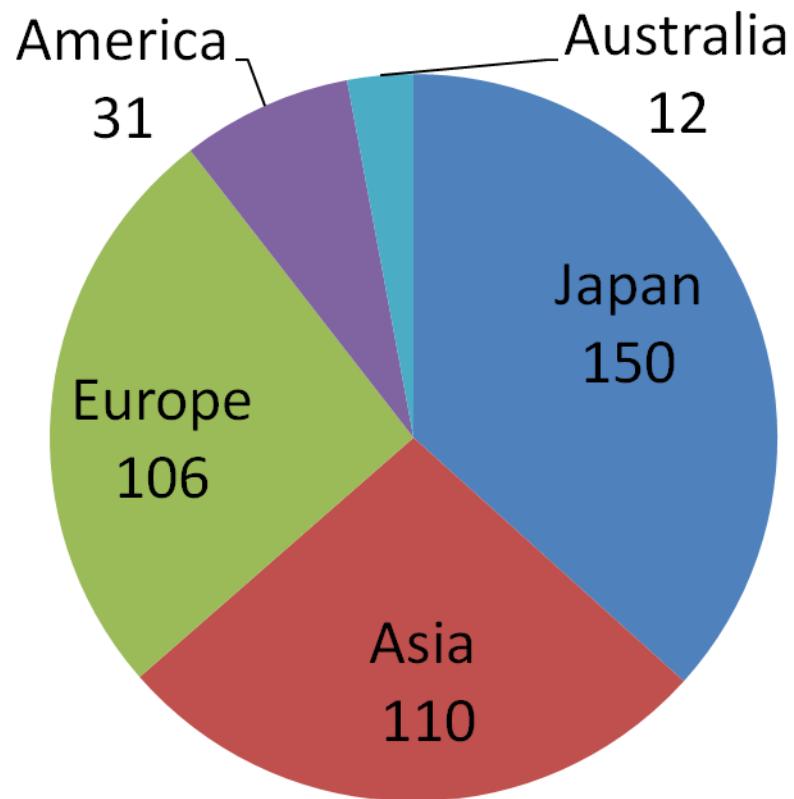


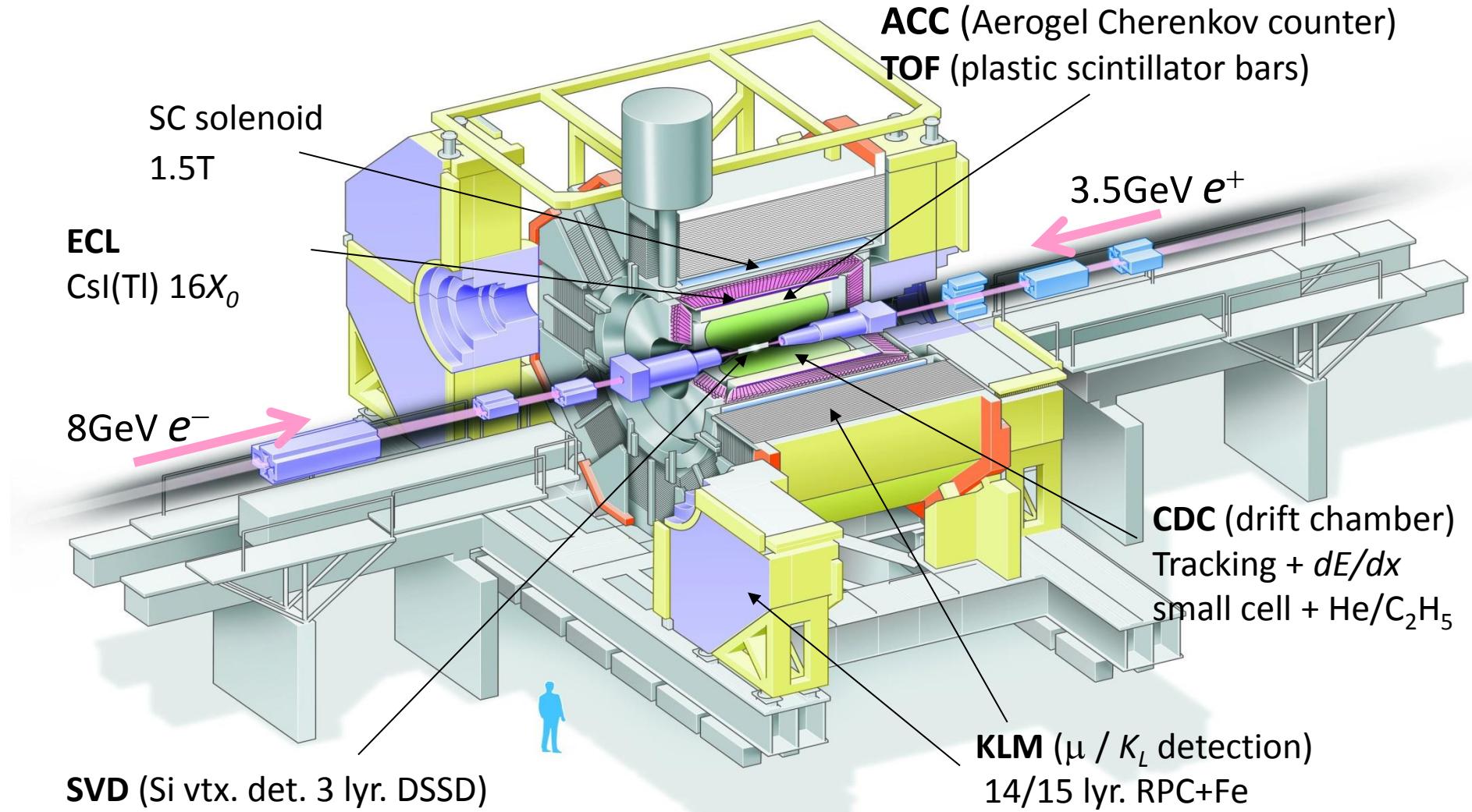


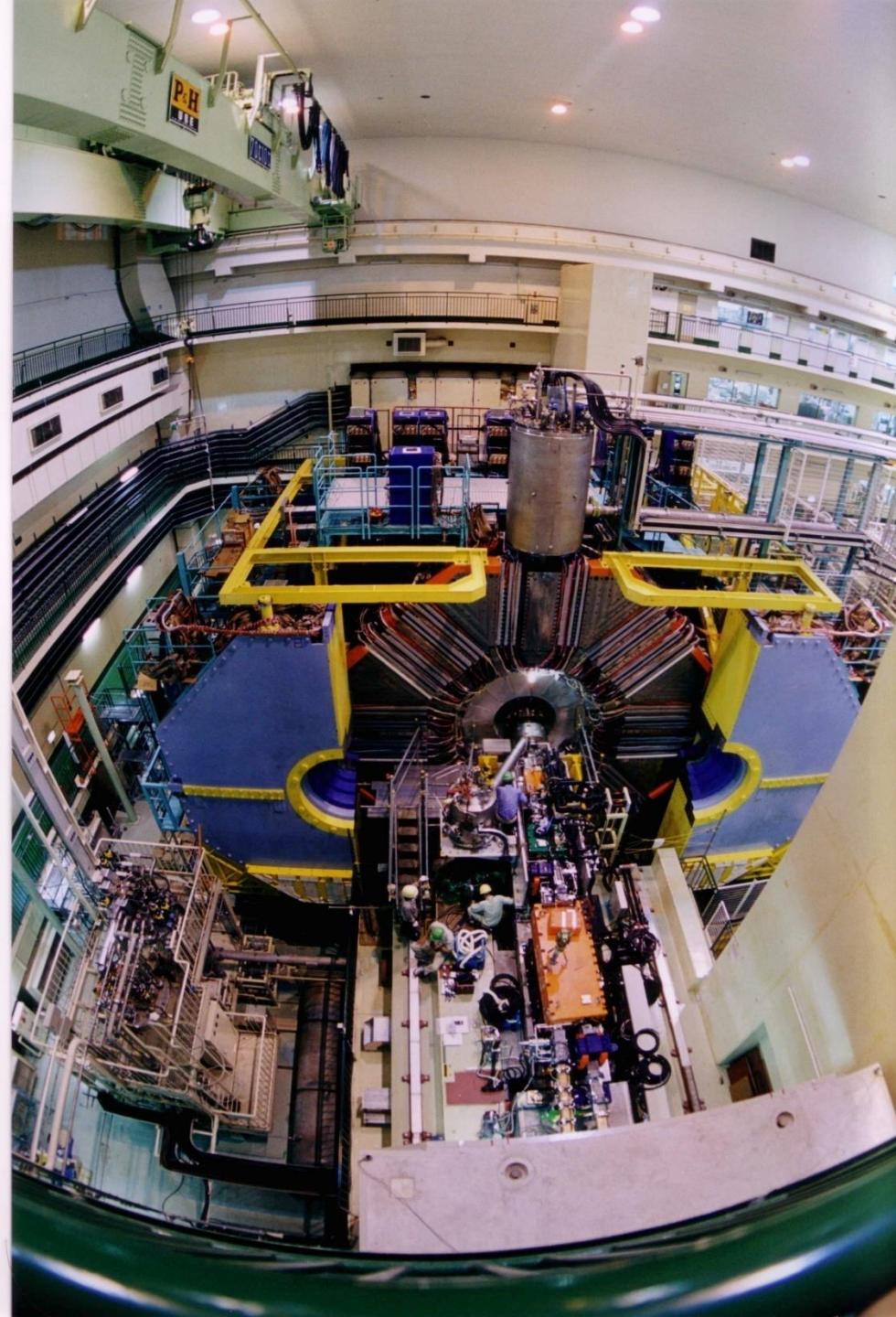


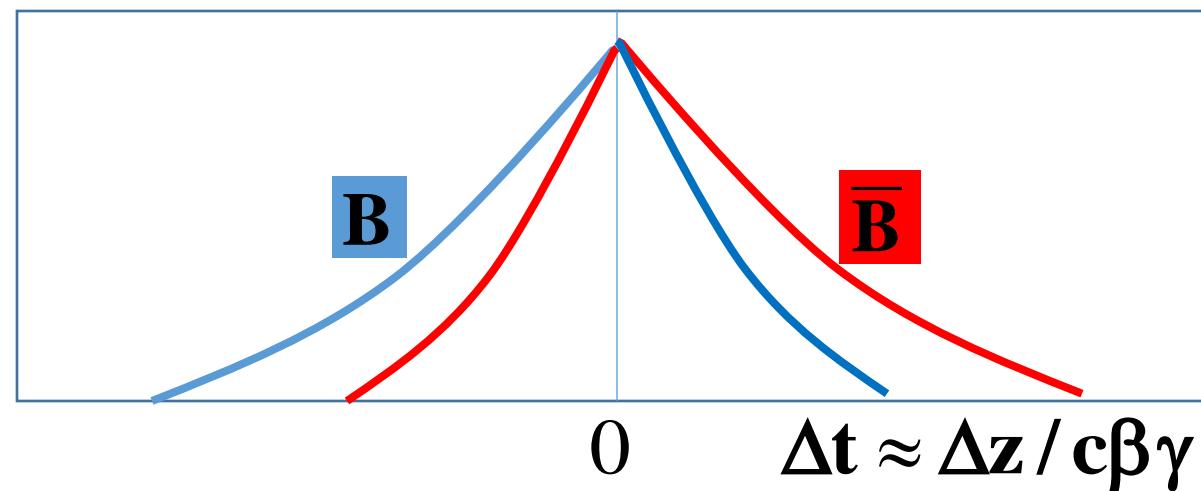
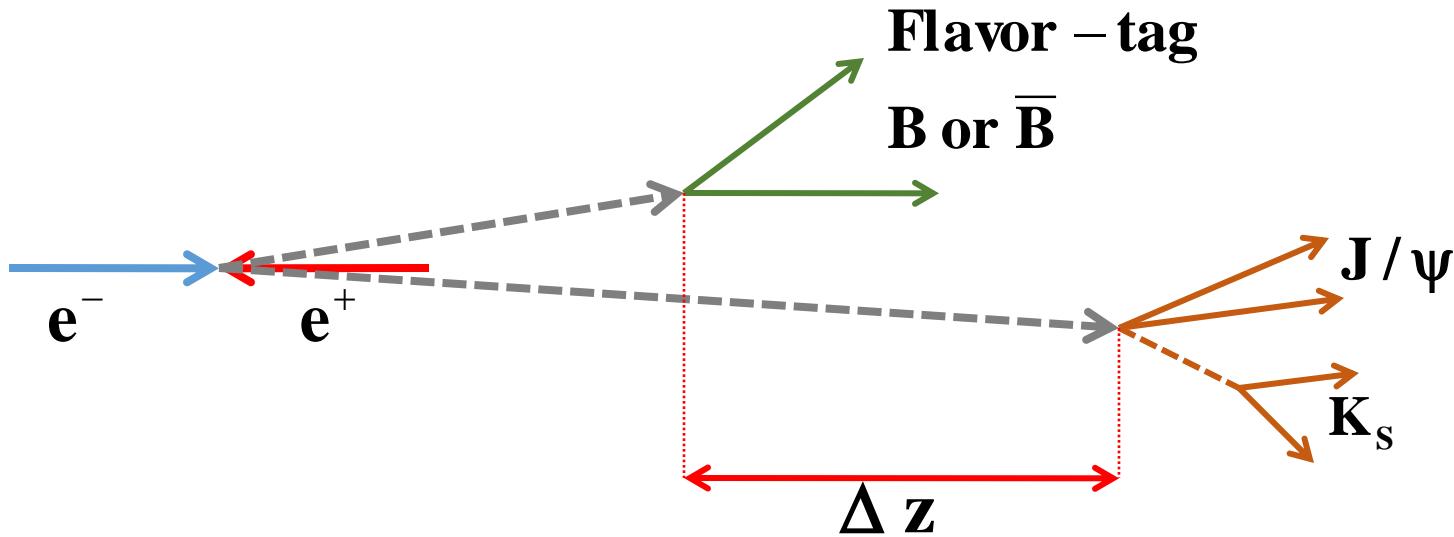
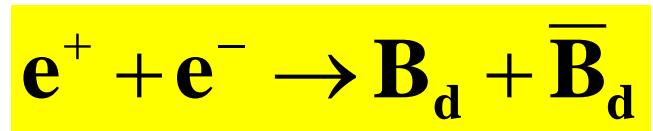
Belle Members

15 Countries and Regions
62 Institution
409 Members



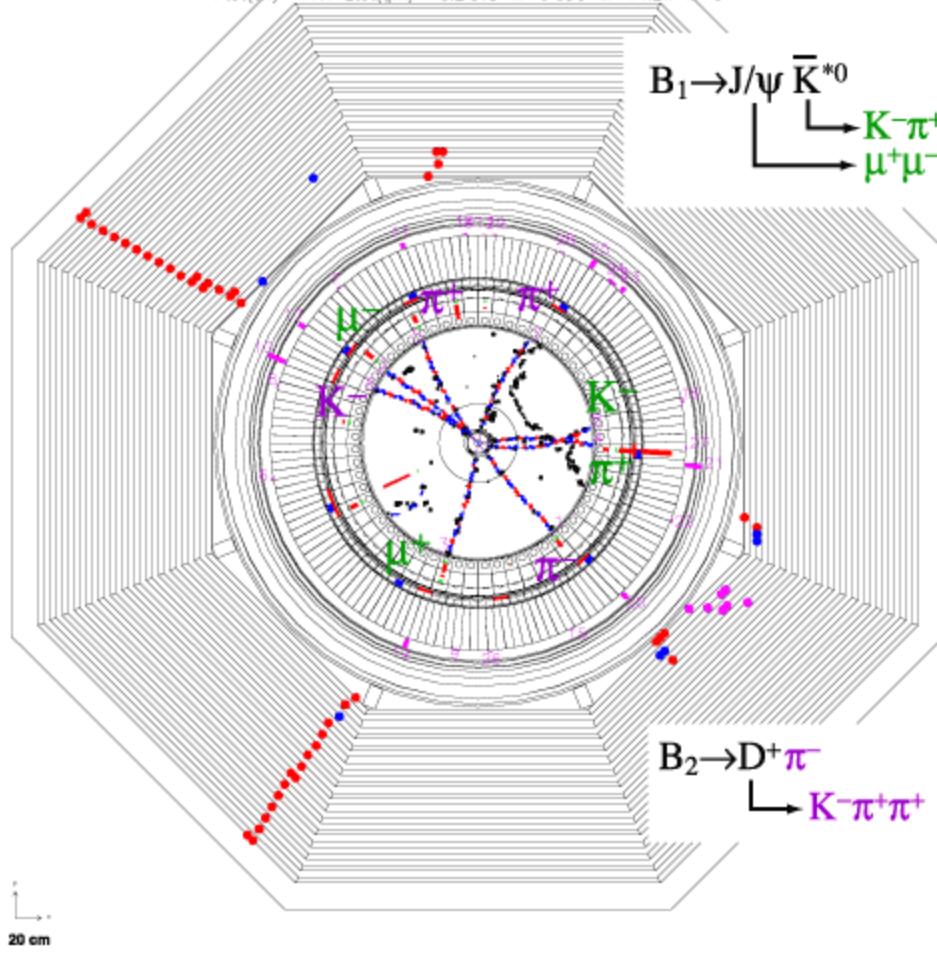


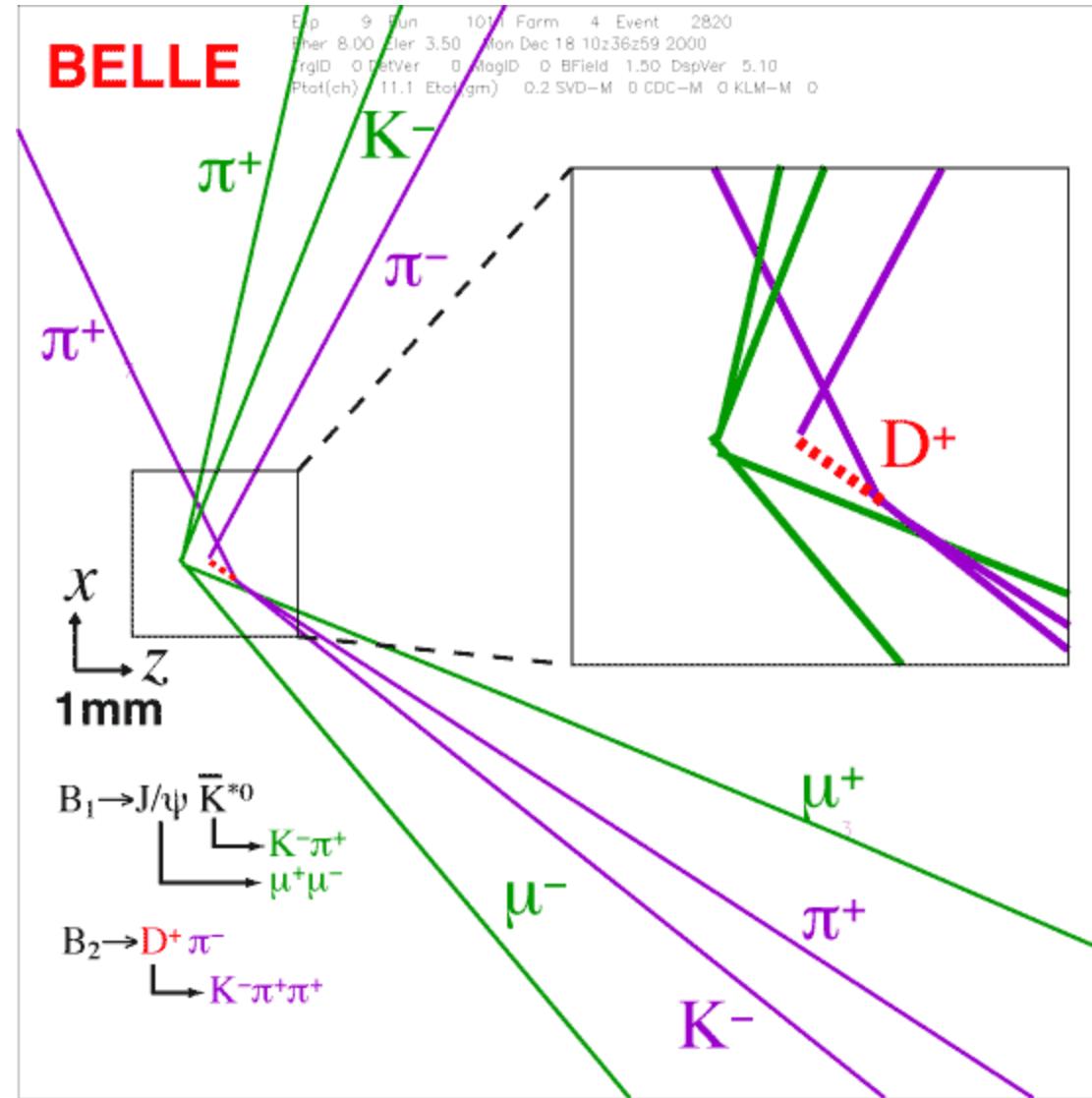


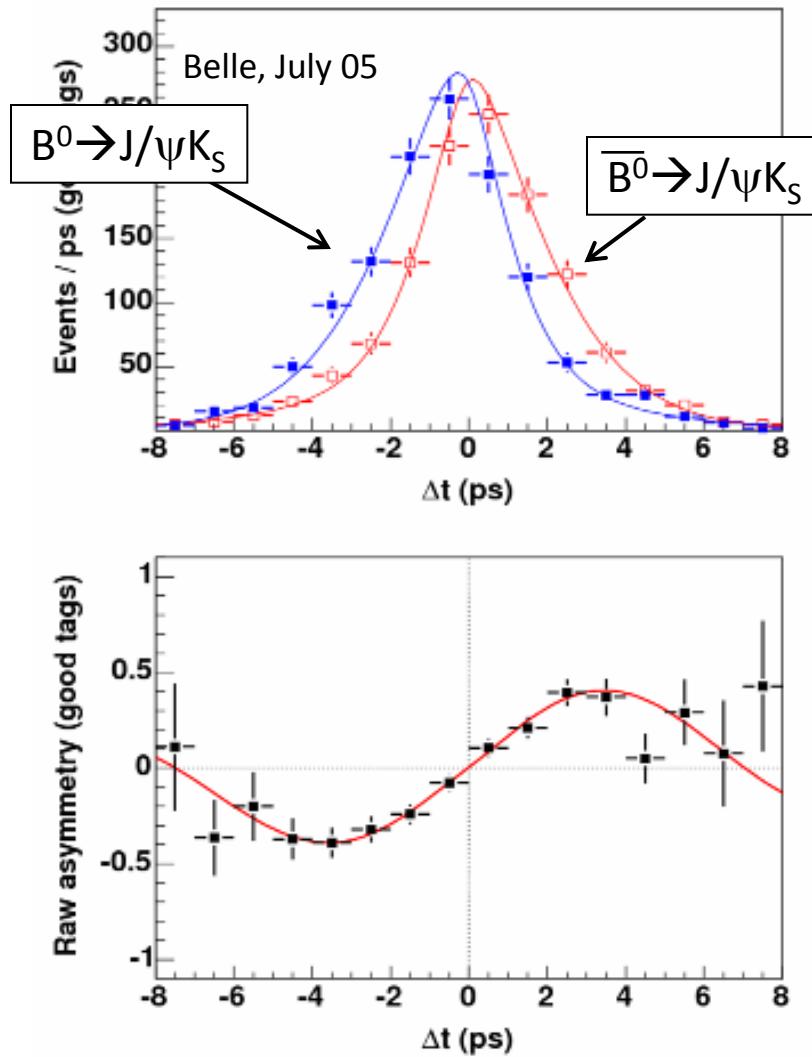


BELLE

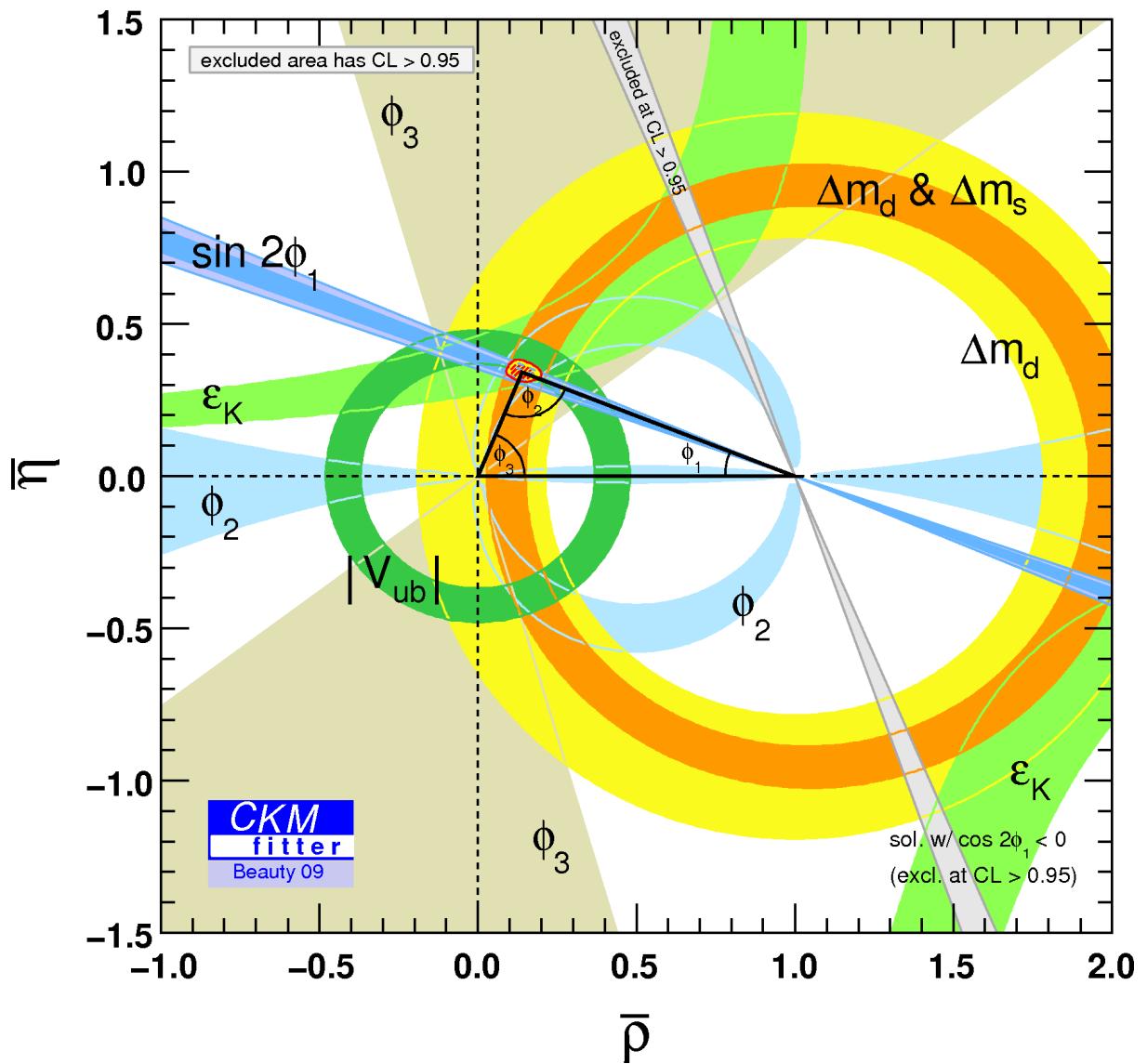
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TrgID 0-DetVer 0-MagID 0-BField 1.50-DspVer 5-10
Ptot(ch) 11.1 Etot(gm) 0.2 SVD-M 0 CDC-M 1 KLM-M 0







Belleグループ提供

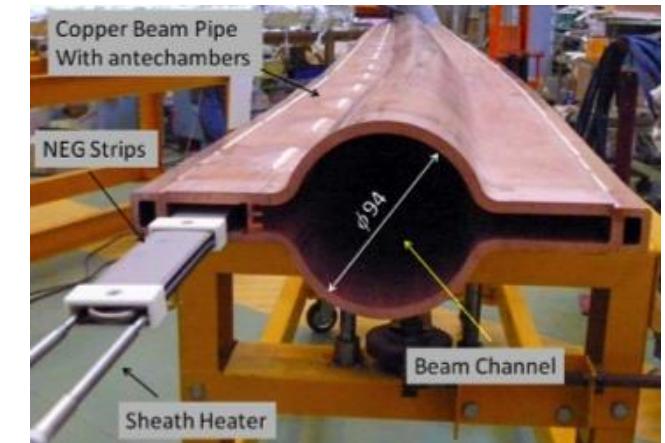
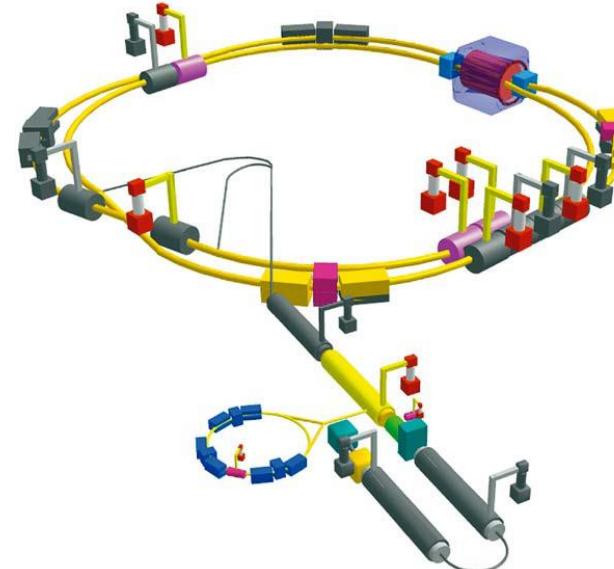


[CKMfitter Group \(J. Charles *et al.*\)](#),
[Eur. Phys. J. C41, 1–131 \(2005\) \[hep-ph/0406184\]](#),
[updated results and plots available at: <http://ckmfitter.in2p3.fr>](#)

Super KEKB

40 times higher luminosity

Commissioning starts in January 2016

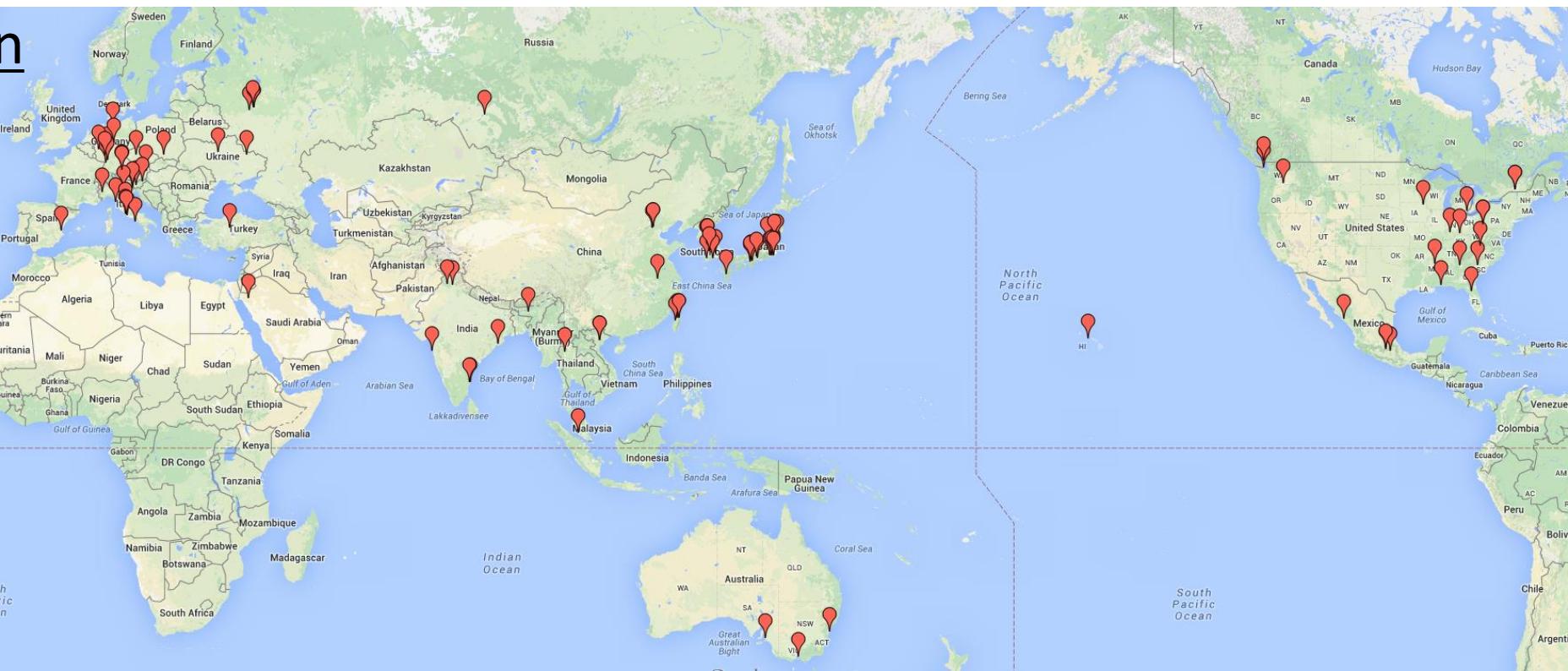


Belle-II Collaboration

23 countries/regions

99 institutions

>600 scientists



Big-Bang

The Universe was in a high-temperature high-density state

Particles and antiparticles coexisted

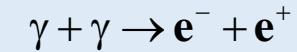


Pair annihilation



Matter was composed from remnant particles

$$\langle kT \rangle \gg mc^2$$



How did number difference occur?

Sakharov's three conditions

- Baryon number non-conservation
- C and **CP violation**
- Imequilibrium

Present status of CP violation

- The standard model explains CP violation in both K-mesons and B-mesons
- The standard model, however, cannot explain the matter dominance of the Universe.
There will be yet unknown particles and unknown CP violating interactions.