

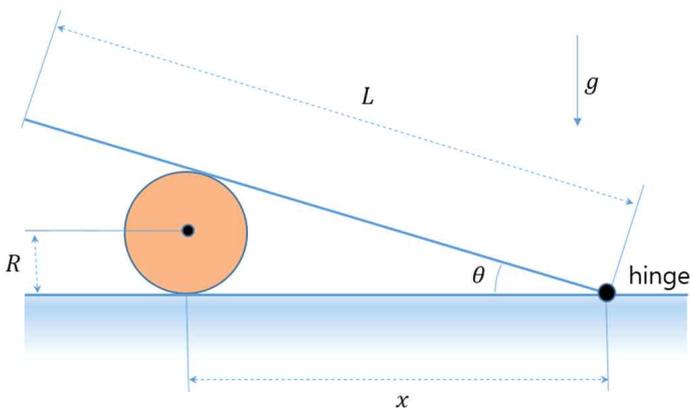
자격시험 문제지

과목명 :

고전역학

2017 . 07. 18 시행

1. [30 pts] A cylinder of radius R and uniformly distributed mass m is lying on the ground. A rod of length L and uniform mass M is hinged on the ground (rotating freely around it). The rod leans on the cylinder as shown in the Figure, forming an angle θ with the ground. The distance between the hinge and the bottom of the cylinder is x ($< L$). [Note that, $\tan(\theta/2) = R/x$.] There are no frictions in this system, and the gravitational acceleration is g . The whole system is initially at rest.



(a) [5 pts] Let us denote by N_1, N_2 the normal force on the cylinder by the ground and by the rod, respectively. Show that $N_1 - mg = N_2 \cos \theta$.

(b) [10 pts] Compute the net force F on the cylinder, in terms of N_2 and θ . Also, compute the torque τ on the rod around the hinge, in terms of N_2 , θ and other constants introduced above.

(c) [10 pts] Show that F and τ are related to each other, at the initial time, by

$$F = \frac{2mR}{ML^2 \sin^2(\theta/2)} \tau .$$

(d) [5 pts] Compute N_2 at the initial time, when $\theta = \frac{\pi}{3}$ using the results from (b) and (c).

2. [30 pts] Consider two interacting charged particles with same mass m . The particles are moving in the xy -plane, perpendicular to a uniform magnetostatic field $\vec{B} = B\hat{z}$ where $B > 0$. Neglect radiation and gravity.

(a) [15 pts] Suppose the charges are identical ($q_1 = q_2 = q$). Find the equations of motion for centre-of-mass (CM) and relative motions separately. If the charges have opposite sign, show the equations of motion do not separate into equations for CM and relative motions.

(b) [8 pts] Express the equation of motion for relative motion from (a) in terms of plane polar coordination (r, θ) , in the case of identical charges. Show the $L_w \equiv \mu r^2 \left(\dot{\theta} + \frac{1}{2}\omega \right)$ is constant ($w = \frac{qB}{m}$).

(c) [7 pts] If the equation is rewritten for the effective potential $V_e(r)$, show that the motion is always bounded in the presence of a magnetic field.

자격시험 문제지

과목명 : 양자역학

2017 . 07. 18 시행

1. [30 pts] Consider *monoenergetic* beam of spin $1/2$ particles of mass m and energy E moving in x direction. The potential energy operator is given by

$$V(x) = V_0 - \gamma B_0 S_z \text{ for } x > 0$$

$$V(x) = 0 \text{ for } x \leq 0$$

where $V_0 > 0$ is a positive constant potential, γ is gyromagnetic ratio, $B_0 > 0$ is magnitude of external magnetic field applied in z direction, and S_z is the z -direction spin operator.

(a) [4 pts] Write the Hamiltonian for the particles in the $x > 0$ region and sketch the potential energy as a function of x for particles having z -direction spins *up and down, respectively*.

(b) [4 pts] Suppose that the spin of particles coming from $-\infty$ is in *eigenstate of the x -direction spin operator S_x* with eigenvalue $+\hbar/2$ and the energy of each particle is $E > V_0 > 0$. Write down the general eigenstate of such an incoming beam considering spatial and spinor parts.

(c) [5 pts] Write down the general solution of transmitted and reflected beams ($E > V_0 > 0$).

(d) [6 pts] What are the boundary conditions at $x=0$? Using the boundary conditions, write down the equations that must be satisfied by the amplitudes appearing in parts (b) and (c).

(e) [6 pts] If $E = V_0 > 0$, what is the probability of measuring z direction spin angular momentum $+\hbar/2$ for the transmitted beam?

(f) [5 pts] Now, instead of the incoming beam with the x -direction spin polarization (*i.e.*, eigenstate of S_x), if we start with *unpolarized* incoming beam (but still $E = V_0 > 0$), what is the probability of measuring the z -direction spin angular momentum $+\hbar/2$ for the transmitted beam?

2. [30 pts] Consider that we have a hydrogen atom. The orbital angular momentum of the electron is $l=1$ (p -state). The total angular momentum operator is defined as the sum of the orbital and spin angular momentum operators: $\mathbf{J} = \mathbf{L} + \mathbf{S}$ where $\mathbf{L} = \mathbf{r} \times \mathbf{p}$. The Pauli spin matrices are given by $\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. Also, it is known that the time-reversal operator for an electron is given by $T = -i\sigma_y K$ where K is the complex-conjugation operator.

(a) [8 pts] What are the possible eigenstates of the total angular momentum operators J^2 and J_z ? In other words, determine all the possible eigenstates of the total (orbital and spin) angular momentum operators of the form $|j, m_j\rangle$. For example, $\left|j = \frac{1}{2}, m_j = +\frac{1}{2}\right\rangle$ is one of such eigenstates.

(b) [8 pts] Express the eigenstate $\left|j = \frac{1}{2}, m_j = +\frac{1}{2}\right\rangle$ in terms of the eigenstates of S_z and L_z operators, *i.e.*, $|m_s, m_l\rangle$ states. Hints: Start from $\left|j = \frac{3}{2}, m_j = +\frac{3}{2}\right\rangle = \left|m_l = +1, m_s = +\frac{1}{2}\right\rangle$ and apply J_- on both sides. By doing this, you will obtain the results for $j = \frac{3}{2}$ states. Then, use the orthonormality condition to obtain the expressions for $j = \frac{1}{2}$ states. You may also use the relations $J_{\pm}|j, m_j\rangle = \sqrt{j(j+1) - m_j(m_j \pm 1)} \hbar |j, m_j \pm 1\rangle$ and $J_z|j, m_j\rangle = m_j \hbar |j, m_j\rangle$ without proof.

(c) [5 pts] Prove that $T^{-1} = i\sigma_y K$.

(d) [9 pts] Obtain $T L_i T^{-1}$ and $T \sigma_i T^{-1}$ ($i = x, y, z$). Note that you need to obtain 6 answers. Each correct answer counts 1.5 pts.

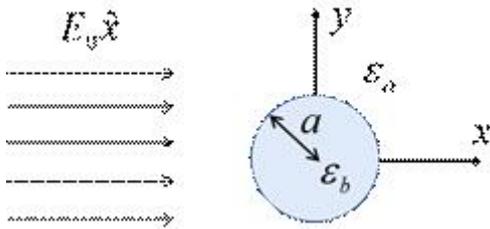
자격시험 문제지

과목명 :

전기역학

2017 . 07. 18 시행

1. [30 pts] A uniform electric field $E_0 \hat{x}$, perhaps produced by means of a parallel plate capacitor, exists in a dielectric having permittivity ϵ_a . With its axis perpendicular to this field, a sufficiently long circular cylindrical dielectric rod [extended along the z (out-of-paper) direction] having permittivity ϵ_b and radius a is introduced.



Because the rod is long and the original field is uniform, the problem reduces down to a two-dimensional polar coordinate problem. Let the axis of cylinder be on the z -axis. (ρ, θ, z) is the cylindrical coordinate. We will look for solutions of Laplace equation inside ($\phi_b(\rho, \theta)$) and outside ($\phi_a(\rho, \theta)$) of rod subjecting to the boundary conditions

(i) $\phi_a = \phi_b$ at $\rho = a$,

(ii) $\epsilon_b \frac{\partial \phi_b}{\partial \rho} = \epsilon_a \frac{\partial \phi_a}{\partial \rho}$ at $\rho = a$, and

(iii) $\phi_a \rightarrow -E_0 x = -E_0 \rho \cos \theta$ for $\rho \gg a$.

(a) [7 pts] Justify boundary conditions (ii) and (iii). You may use the three boundary conditions in solving the next problems even if you haven't solved (a).

(b) [15 pts] The general form of the potential outside and inside the rod satisfying the boundary conditions is given, respectively, by

$$\phi_a = -E_0 \rho \cos \theta + \sum_{k=1}^{\infty} \rho^{-k} (C_k \cos k\theta + D_k \sin k\theta)$$

and $\phi_b = \sum_{k=1}^{\infty} \rho^k (A_k \cos k\theta + B_k \sin k\theta)$.

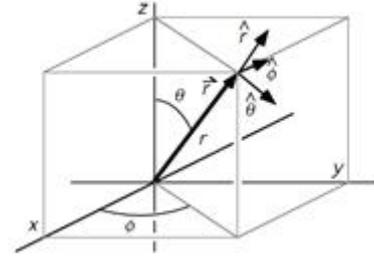
Determine the potential and the electric field at points outside and inside the rod, neglecting end effects.

(c) [8 pts] Sketch the electric field lines inside and outside the rod when $\epsilon_a > \epsilon_b$.

2. [30 pts] The simplest possible spherical electromagnetic wave can be written as

$$\mathbf{E}(r, \theta, \phi, t) = A \frac{\sin \theta}{r} \left[\cos(kr - \omega t) - \frac{1}{kr} \sin(kr - \omega t) \right] \hat{\phi},$$

where $\omega/k = c$.



(a) [10 pts] Obtain the associated magnetic field \mathbf{B} using Faraday's law.

(b) [10 pts] Calculate the Poynting vector $\mathbf{S} = \frac{1}{\mu_0} \mathbf{E} \times \mathbf{B}$.

(c) [10 pts] Take the time average of the Poynting vector to get the intensity vector $\mathbf{I} = \langle \mathbf{S} \rangle$. Discuss about the results, including the r -dependence.

※ You may use the following formula:

$$\nabla_s = \hat{r} \frac{\partial s}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial s}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial s}{\partial \phi}$$

$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$$

$$\begin{aligned} \nabla \times \mathbf{v} = & \hat{r} \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \\ & + \hat{\theta} \left[\frac{1}{r \sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r v_\phi) \right] \\ & + \hat{\phi} \frac{1}{r} \left[\frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \end{aligned}$$

※ You may also use the following vacuum Maxwell equation if necessary:

$$\nabla \cdot \mathbf{E} = 0, \quad \nabla \cdot \mathbf{B} = 0, \quad \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E},$$

$$\nabla \times \mathbf{B} = \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

자격시험 문제지

과목명 :

통계역학

2017 . 07. 18 시행

1. [60 pts] In a gas of free and independent electrons in three-dimensions, the one-electron levels are specified by the wave vector \mathbf{k} and spin quantum number s with the energy given by $E(k) = \frac{\hbar^2 k^2}{2m}$. At zero temperature, the Fermi energy E_F can be defined in a way that the levels with $E(k) \leq E_F$ ($E(k) > E_F$) are occupied (unoccupied).

(a) [5 pts] At the temperature T , the energy density ϵ and the particle number density n can be written as

$$\epsilon = \frac{2}{V} \sum_{\mathbf{k}} E(k) f(E(k)) \quad \text{and} \quad n = \frac{2}{V} \sum_{\mathbf{k}} f(E(k))$$

where $f(E) = \frac{1}{e^{(E-\mu)/k_B T} + 1}$ is the Fermi-Dirac distribution function with the Boltzmann constant k_B and the chemical potential μ . Using the definition of the density of states $g(E) = \frac{2}{V} \sum_{\mathbf{k}} \delta(E - E(k))$ show that ϵ and n can be written as

$$\epsilon = \int_{-\infty}^{\infty} dE g(E) E f(E) \quad \text{and} \quad n = \int_{-\infty}^{\infty} dE g(E) f(E).$$

(b) [10 pts] Compute $g(E)$ explicitly and show

$$g(E) = \frac{m}{\hbar^2 \pi^2} \sqrt{\frac{2mE}{\hbar^2}} = \frac{3}{2} \frac{n}{E_F} \left(\frac{E}{E_F}\right)^{1/2}.$$

(c) [10 pts] In metals, when the temperature T is much smaller than the chemical potential μ , one can perform the above energy integral by using the Sommerfeld expansion, which leads to

$$\epsilon = \int_0^{\mu} dE g(E) E + \frac{\pi^2}{6} (k_B T)^2 [\mu g'(\mu) + g(\mu)] + O(T^4)$$

$$n = \int_0^{\mu} dE g(E) + \frac{\pi^2}{6} (k_B T)^2 g'(\mu) + O(T^4).$$

Assuming that the chemical potential μ differs from its $T=0$ value E_F by terms of order T^2 , one can write

$$\int_0^{\mu} dE h(E) = \int_0^{E_F} dE h(E) + (\mu - E_F) h(E_F)$$

where $h(E)$ is an arbitrary function. Applying this formula show that $\mu = E_F - \frac{\pi^2}{6} (k_B T)^2 \frac{g'(E_F)}{g(E_F)}$ and

$\epsilon = \epsilon_0 + \frac{\pi^2}{6} (k_B T)^2 g(E_F)$ where ϵ_0 is the energy density in the ground state.

(d) [5 pts] Show that the specific heat of the free electron gas is $c_v = \frac{\pi^2}{2} \left(\frac{k_B T}{E_F}\right) n k_B$.

(e) [10 pts] To compute the heat capacity $C = \gamma T$, the specific heat should be multiplied by the volume of the system. Suppose that a mole of free electron metal contains $Z N_A$ conduction electrons (Z is the valence and N_A is Avogadro's number). Calculate γ using $R = k_B N_A = 8.314$ joules/(mole K) = 1.99 calories/(mole K) and $E_F/k_B = 10000$ K and $Z = 3$. Compute γ in units of J/(mol K) and round up to two significant figures.

(f) [5 pts] In the presence of external magnetic field H , suppose that the energy of an electron with the momentum \mathbf{k} whose spin is parallel (antiparallel) to H is given by $E_+(k) = E(k) - \mu_B H$ [$E_-(k) = E(k) + \mu_B H$]. What is the relevant g -factor?

(g) [5 pt] Since the magnetic field induces just a constant shift of single particle energies, one can define the density of states for electrons with the spin parallel (antiparallel) to H as $g_+(E) = \frac{1}{2} g(E - \mu_B H)$ ($g_-(E) = \frac{1}{2} g(E + \mu_B H)$), and the particle density of each species is given by $n_{\pm} = \int dE g_{\pm}(E) f(E)$ satisfying $n = n_+ + n_-$. When the Zeeman energy is much smaller than the Fermi energy, one can assume that $g_{\pm}(E) = \frac{1}{2} g(E \pm \mu_B H) = \frac{1}{2} g(E) \pm \frac{1}{2} \mu_B H g'(E)$. Using this approximation, find the chemical potential of the system. (Note the the chemical potential is μ when $H=0$.)

(h) [10 pts] The magnetization density is given by $M = -\mu_B (n_+ - n_-)$. Here one can again assume

$$g_{\pm}(E) = \frac{1}{2} g(E \pm \mu_B H) = \frac{1}{2} g(E) \pm \frac{1}{2} \mu_B H g'(E)$$

Show that the zero temperature magnetic susceptibility is $\chi = \frac{M}{H} = \mu_B^2 g(E_F)$.